

Bounding the set of classical correlations of a many-body system

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Outline

- *Non-negative polynomials*
- *Sum-of-squares representations*
- *Characterizing many-body correlations*



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Non-negative polynomials



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 - *Bounding the set of quantum correlations*



Non-negative polynomials

- *A bit of history*



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Who among us would not be happy to lift the veil behind which is hidden the future; to gaze at the coming developments of our science and at the secrets of its development in the centuries to come? What will be the ends toward which the spirit of future generations of mathematicians will tend? What methods, what new facts will the new century reveal in the vast and rich field of mathematical thought?

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Riemann Hypothesis and other number theory conjectures



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Given a multivariate non-negative polynomial, does it admit a sum-of-squares representation?



Non-negative polynomials

- *Hilbert's 17th problem*



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Given $p(\vec{x}) \in \mathbb{R}[\vec{x}]$ satisfying $p(\vec{x}) \geq 0 \forall \vec{x} \in \mathbb{R}^n$,
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- Converse is trivial. When does equivalence hold?
- How powerful is this representation?
- Can one extend it to subsets of \mathbb{R}^n ?

- Semialgebraic sets
$$\begin{cases} f_i(\vec{x}) = 0 \\ g_j(\vec{x}) \geq 0 \end{cases}$$



Non-negative polynomials

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n	1	2	3	4	5
2d					
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8					
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Non-constructive proof



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Example: Cauchy-Schwarz Inequality

$$\|\vec{x}\|^2 \cdot \|\vec{y}\|^2 - \langle \vec{x}, \vec{y} \rangle^2 = \sum_{i < j} (x_i y_j - x_j y_i)^2$$



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Semidefinite Programming

- *Primal-dual formulation*



Semidefinite Programming

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$$\begin{aligned} \min_X \quad & \langle C, X \rangle \\ \text{s.t.} \quad & \langle A_i, X \rangle = b_i \\ & X \succeq 0 \end{aligned}$$



Semidefinite Programming

- Primal-dual formulation

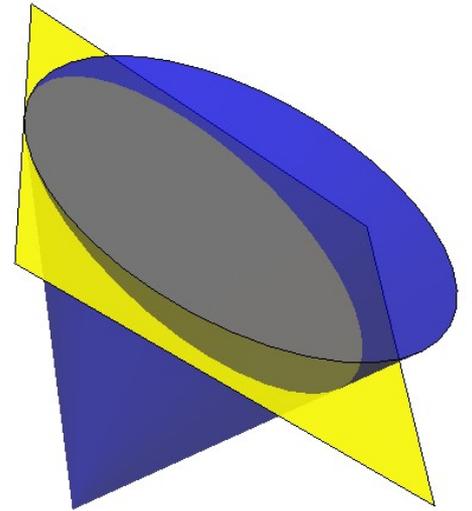
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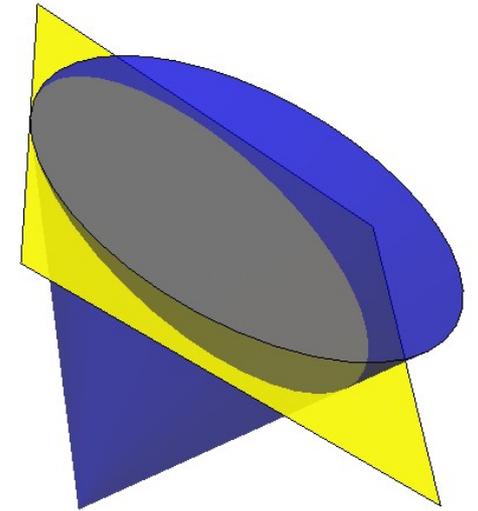
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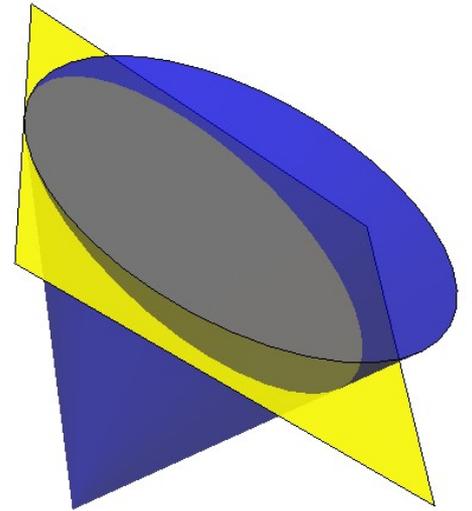
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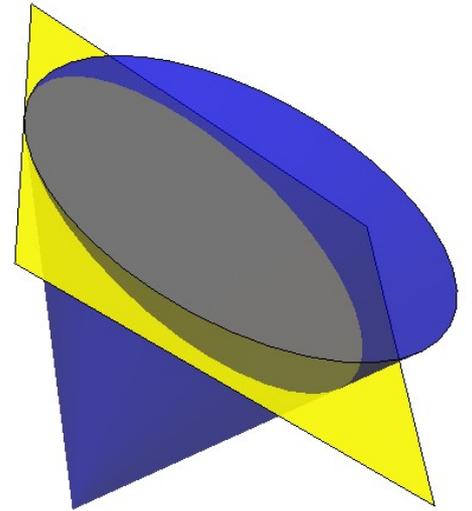
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Deciding if $p(\vec{x})$ admits a sos representation is simply an SdP in disguise

An easy one, since one can assume $\deg(q_i) \leq d$
 $(\deg(p) = 2d)$

$$p(\vec{x}) = \sum_i q_i^2(\vec{x})$$



Semidefinite Programming

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$$Q = \begin{pmatrix} 2 & -3 & 1 \\ -3 & 5 & 0 \\ 1 & 0 & 5 \end{pmatrix} = L^T L \quad L = \frac{1}{\sqrt{2}} \begin{pmatrix} 2 & -3 & 1 \\ 0 & 1 & 3 \end{pmatrix}$$



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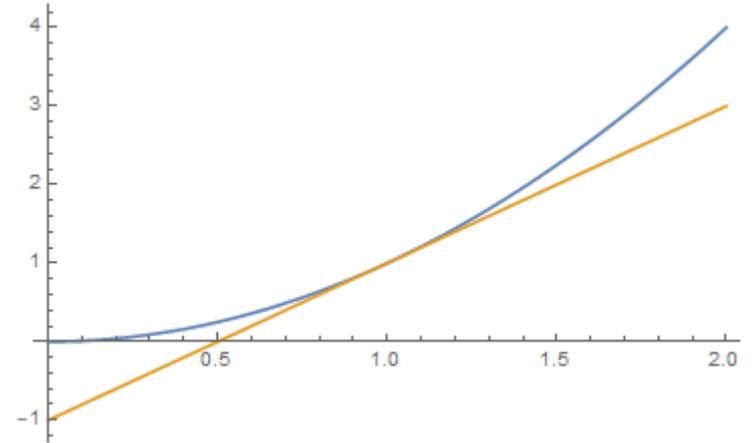
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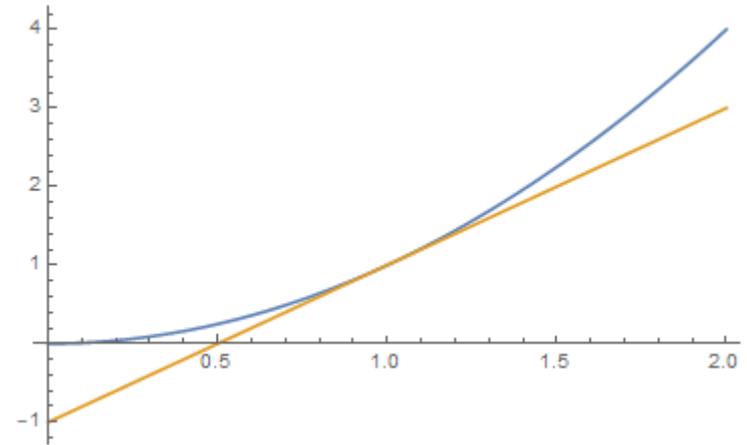


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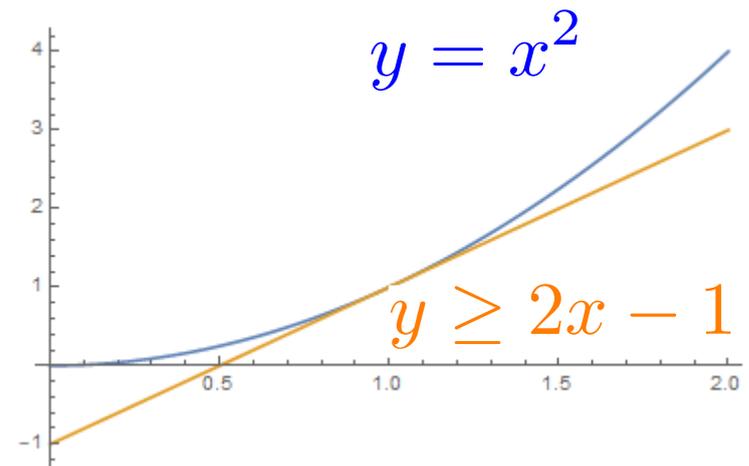


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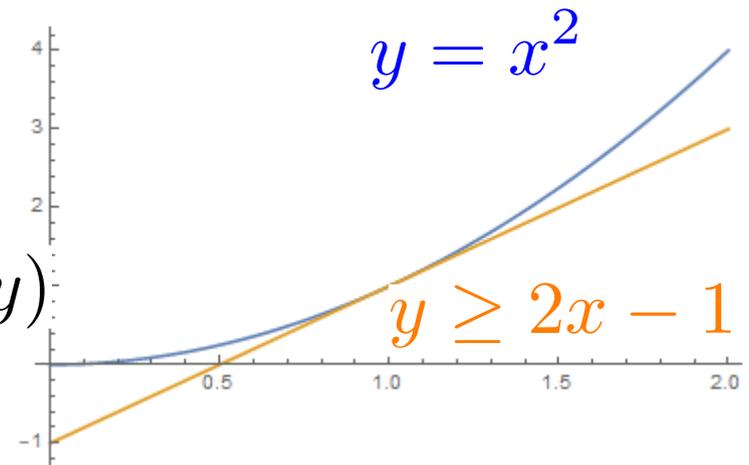
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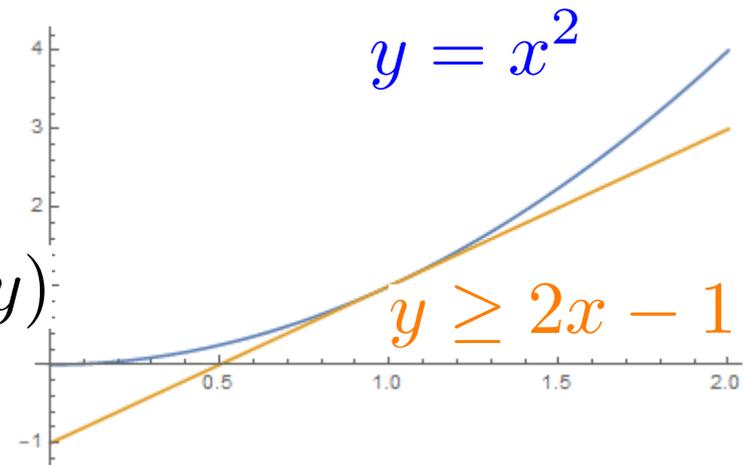
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Sums-of-squares modulo ideals are powerful!



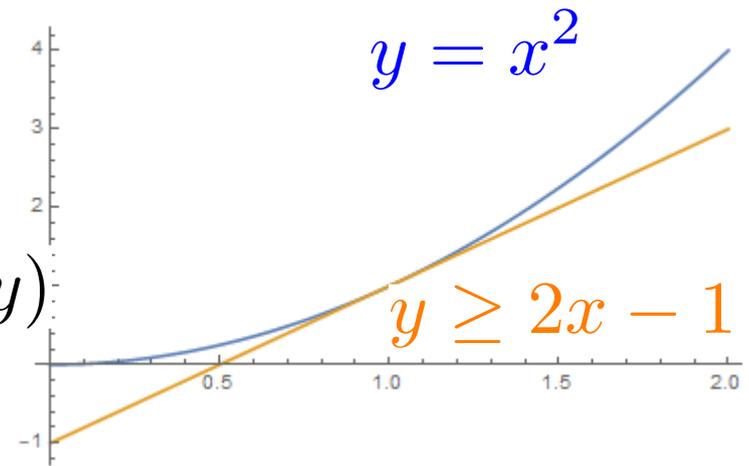
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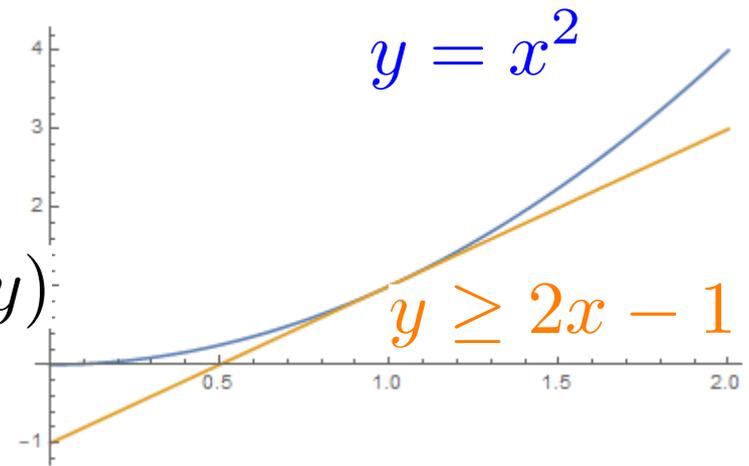
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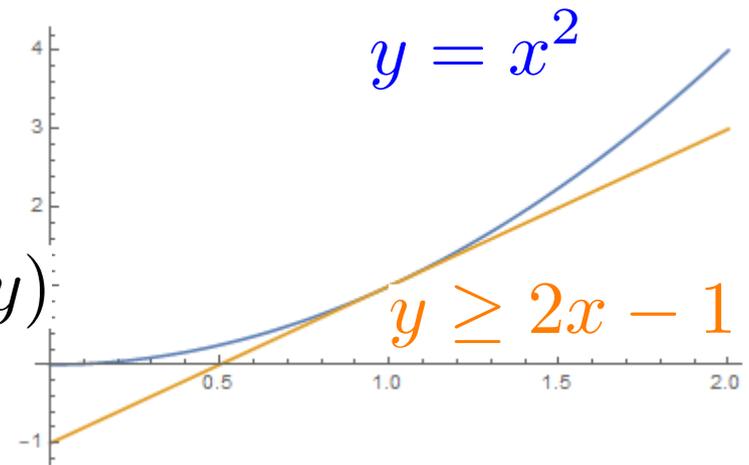
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Ideal generated
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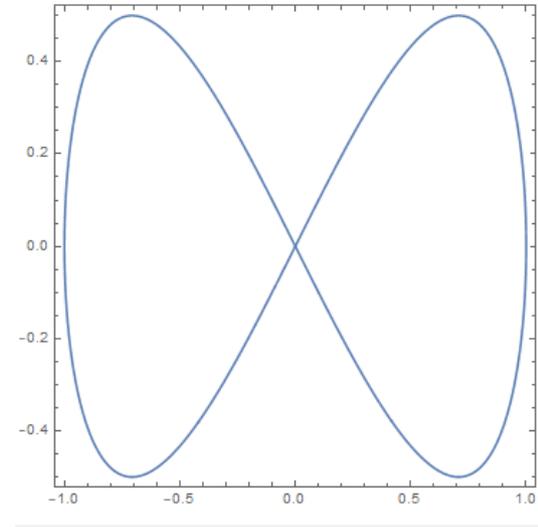
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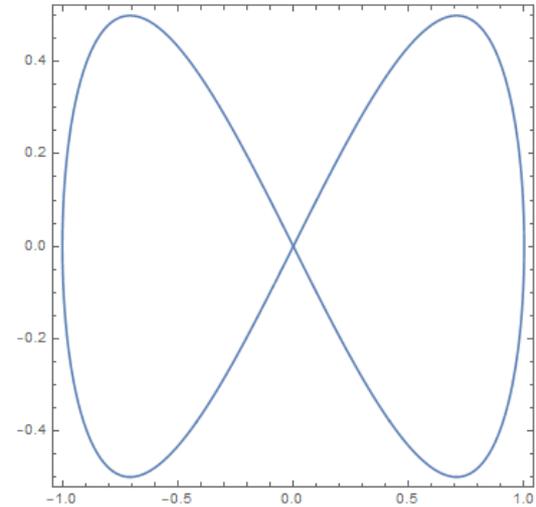
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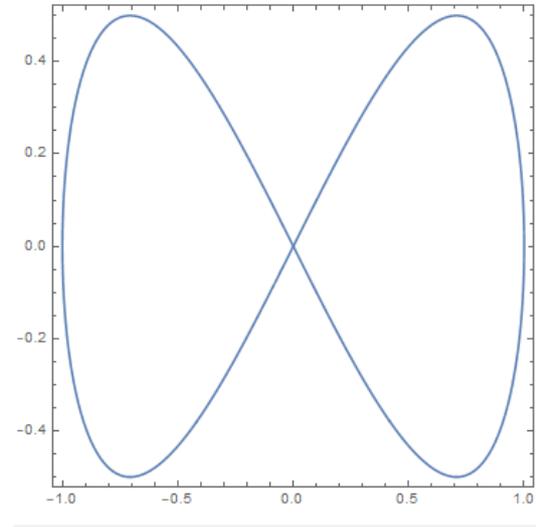
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Gröebner basis



Semidefinite Programming

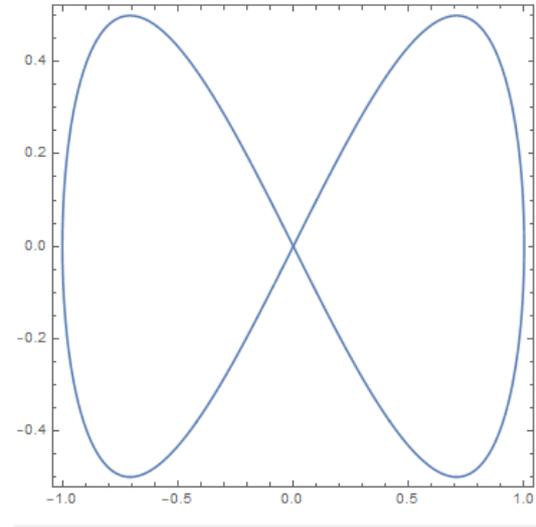
- But now the degree of the sos may be unbounded!!

$$f(x, y) = x^4 - x^2 + y^2 = 0$$

Simplification rule

$$x^4 \rightarrow x^2 - y^2$$

Gröebner basis



$$\begin{pmatrix} 1 & x & y & w_2^0 & w_1^1 & w_0^2 \\ x & w_2^0 & w_1^1 & w_3^0 & w_2^1 & w_1^2 \\ y & w_1^1 & w_0^2 & w_2^1 & w_1^2 & w_0^3 \\ w_2^0 & w_3^0 & w_2^1 & w_2^0 - w_0^2 & w_3^1 & w_2^2 \\ w_1^1 & w_2^1 & w_1^2 & w_3^1 & w_2^2 & w_1^3 \\ w_0^2 & w_1^2 & w_0^3 & w_2^2 & w_1^3 & w_0^4 \end{pmatrix}$$

Moment matrix, 2nd order

[Gouveia, Thomas, Convex Hulls of semialgebraic sets, 2012]

[Lasserre, 2001]



Semidefinite Programming

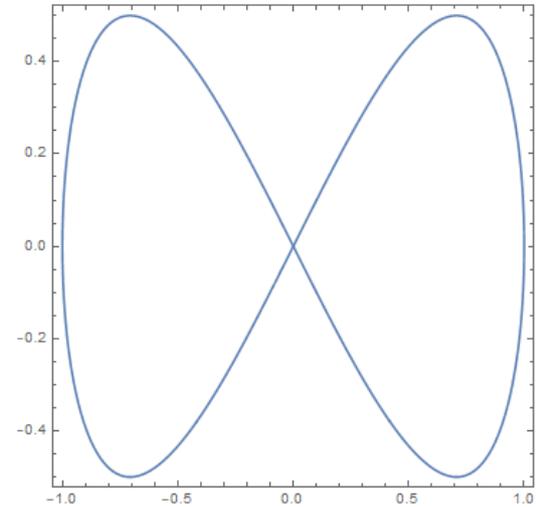
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w_i^j linearizes $x^i y^j$

Moment matrix, 2nd order

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Semidefinite Programming

- We can also add inequality constraints!*



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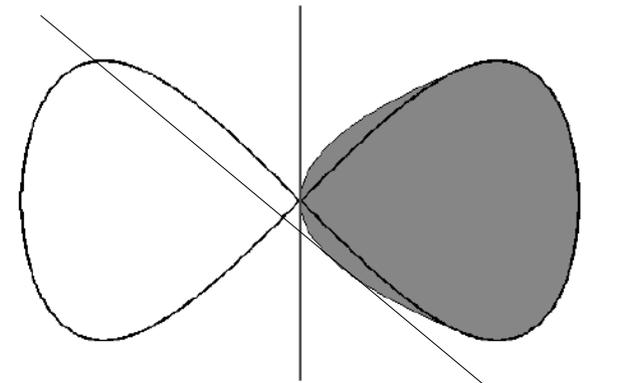


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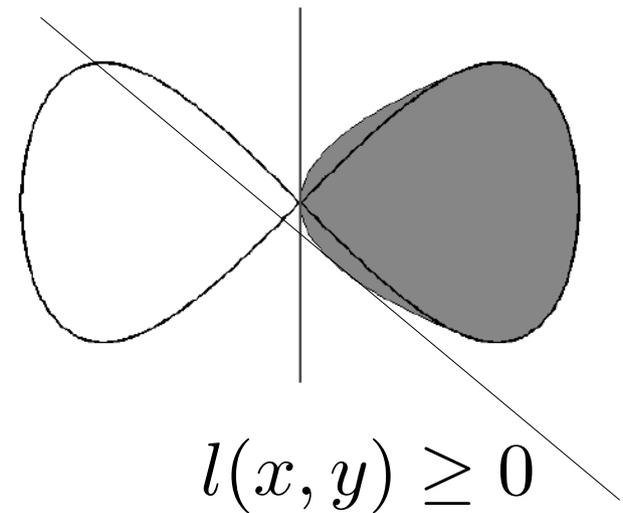
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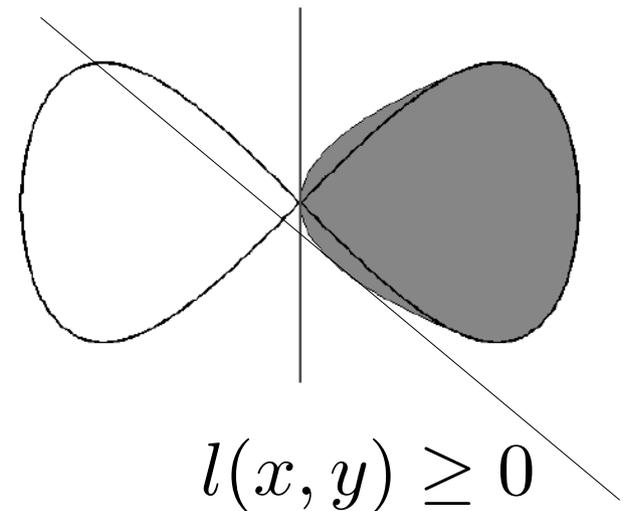
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Moment matrix, 2nd order

$$\begin{pmatrix} x & w_2^0 & w_1^1 \\ w_2^0 & w_3^0 & w_2^1 \\ w_1^1 & w_2^1 & w_1^2 \end{pmatrix}$$

Shifted moment matrix by x , 1st order



[Gouveia, Thomas, Convex Hulls of semialgebraic sets, 2012]

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Semidefinite Programming

- *In general, we can consider*



Semidefinite Programming

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Convex hulls of semialgebraic sets

$$S = \{\vec{x} : f_i(\vec{x}) = 0, g_j(\vec{x}) \geq 0\}$$



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sums-of-squares



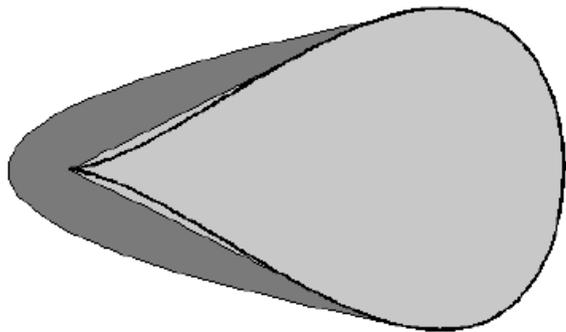
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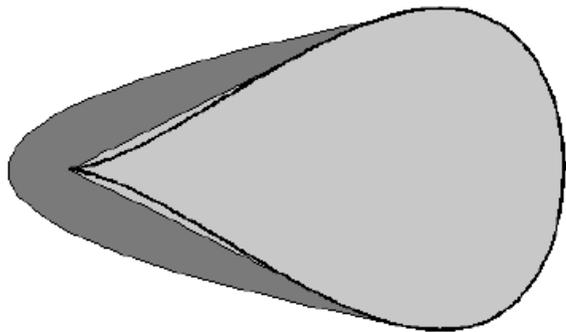
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sums-of-squares



Increasing the degree of
the sos σ_i gives more
representability power



Outline

- *Non-negative polynomials*
- *Sum-of-squares representations*
- *Characterizing many-body correlations*



A bit of Quantum Info

- *Bell-type experiment*



A bit of Quantum Info

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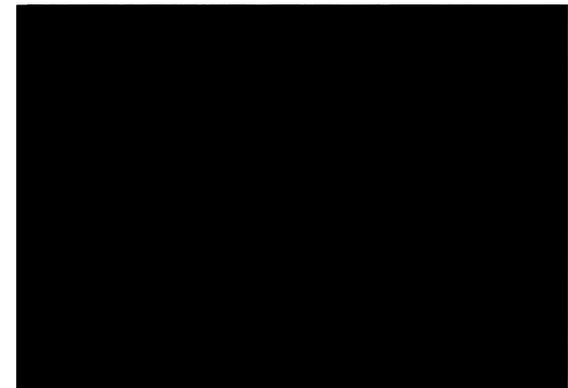
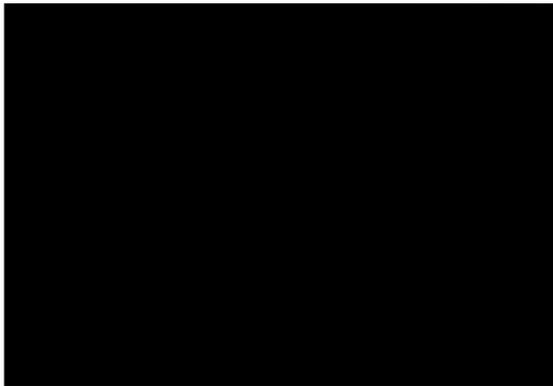
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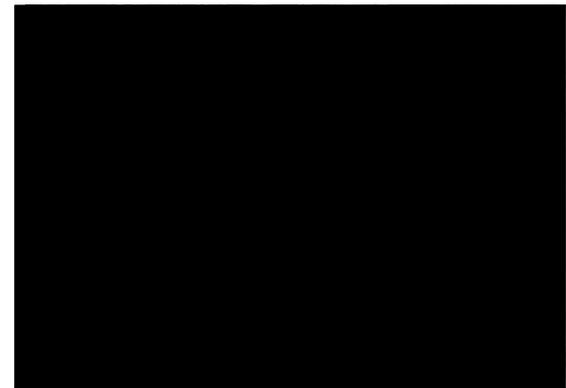
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x



a



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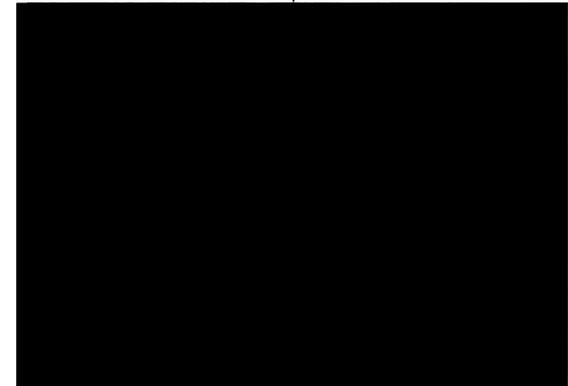
x



a



y



b



A bit of Quantum Info

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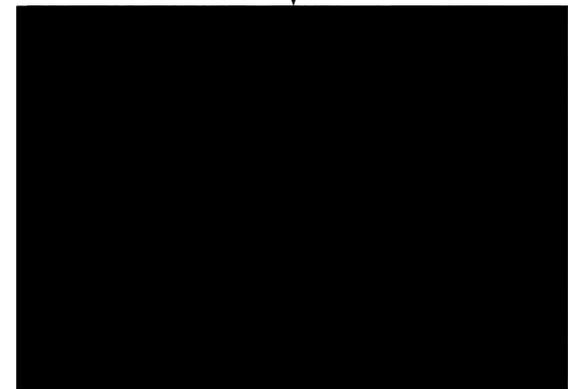
x



a



y



b

$p(ab|xy)$



Nonlocality in many-body quantum systems



Nonlocality in many-body quantum systems

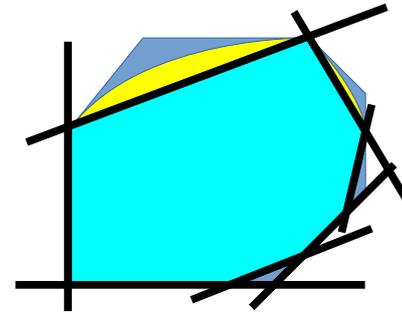
- *The complexity of the problem*



Nonlocality in many-body quantum systems

- The complexity of the problem

Finding all Bell inequalities \longleftrightarrow Convex Hull problem

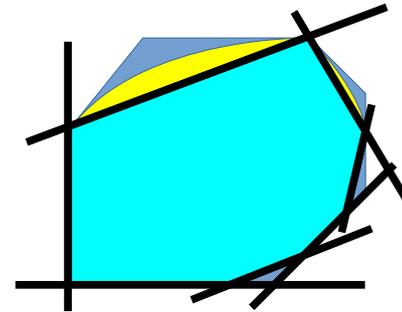


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Finding all Bell inequalities \longleftrightarrow Convex Hull problem

(n, m, d) scenario



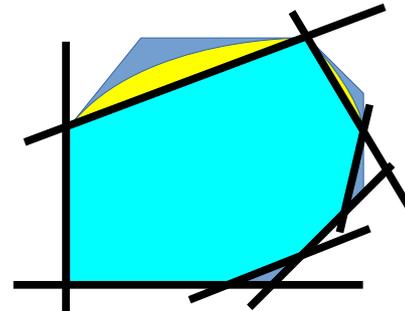
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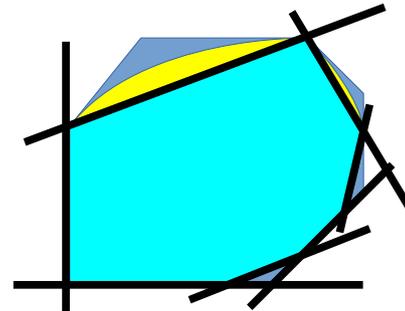
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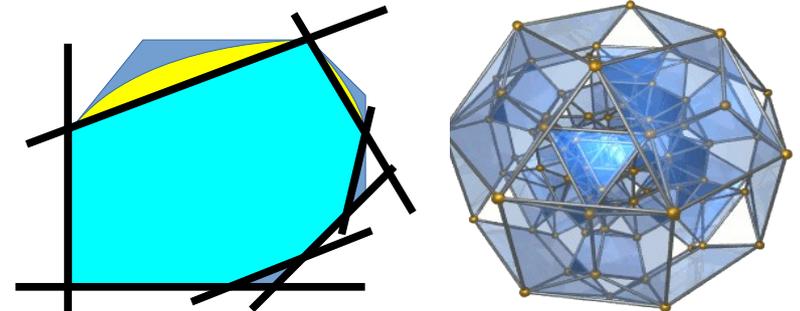
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[B. Chazelle, An optimal convex hull algorithm in any fixed dimension, *Discrete Comput. Geom.* **10** 377409 (1993)]



Nonlocality in many-body quantum systems

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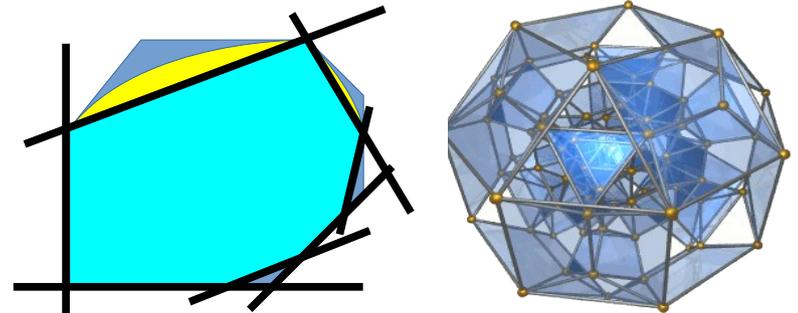
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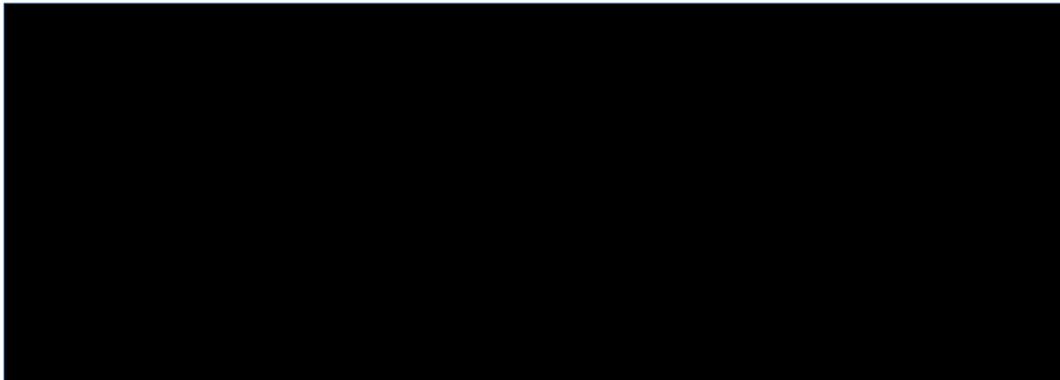
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Examples



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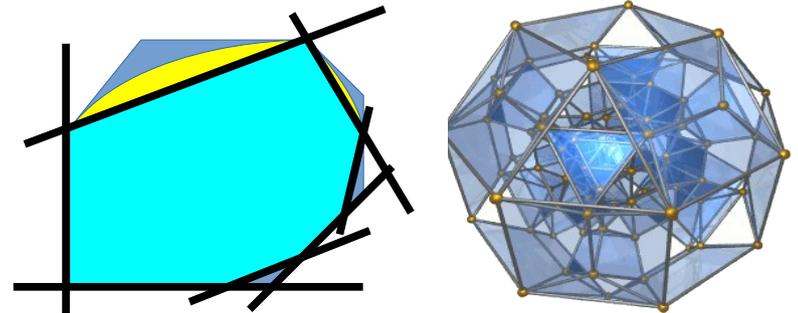
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$(2, 2, 2) \rightarrow O(\text{ms})$



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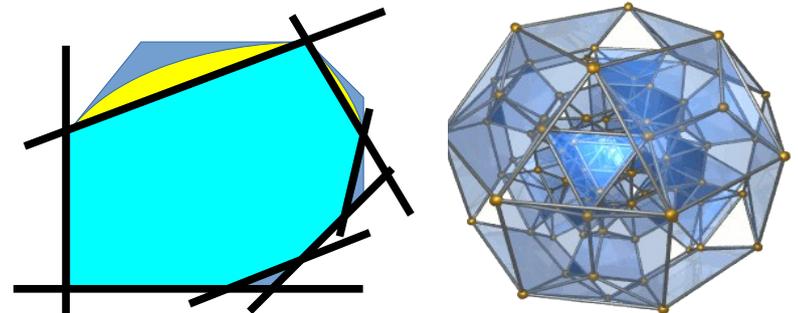
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Examples

$$(2, 2, 2) \longrightarrow O(ms)$$

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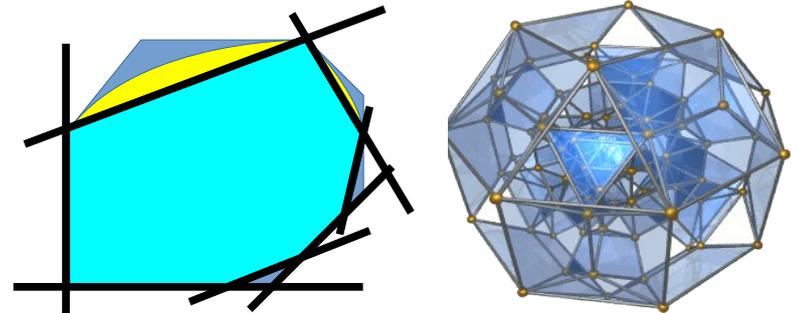
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$(2, 2, 2) \longrightarrow O(\text{ms})$
 $(3, 2, 2) \longrightarrow 5'$
 $(4, 2, 2) \longrightarrow 10^{67}$ years
 \vdots



[S. Dalí *The persistence of memory* (1931)]



Nonlocality in many-body quantum systems

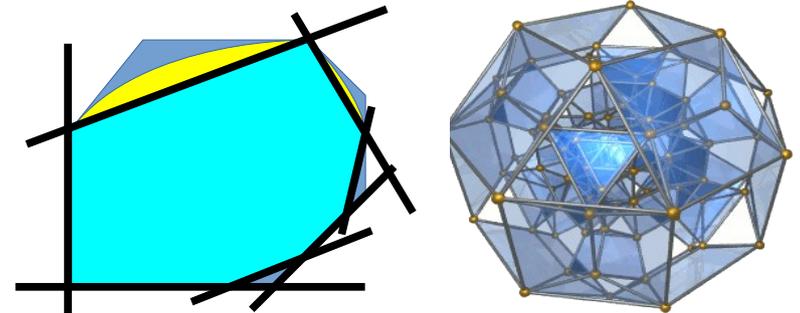
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Examples

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$(3, 2, 2) \longrightarrow 5'$

$(4, 2, 2) \longrightarrow 10^{67}$ years

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$(10^4, 2, 2) \longrightarrow 10^{10^{10^{4.67867\dots}}}$ basically any timescale you want



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Nonlocality in many-body quantum systems



Nonlocality in many-body quantum systems

- *Reducing the mathematical complexity*



Nonlocality in many-body quantum systems

- Reducing the mathematical complexity*

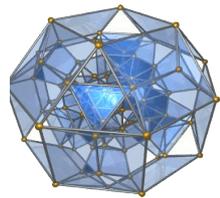
Polytope				
Dimension				
Vertices				



Nonlocality in many-body quantum systems

- Reducing the mathematical complexity

Polytope	\mathbb{P}_n			
Dimension	$3^n - 1$			
Vertices	2^{2n}			

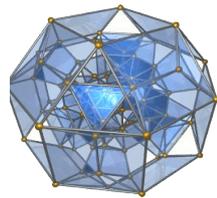


Nonlocality in many-body quantum systems

- Reducing the mathematical complexity

↖ 2-body

Polytope	\mathbb{P}_n	Lower order correlators		
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Vertices	2^{2n}			

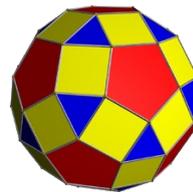
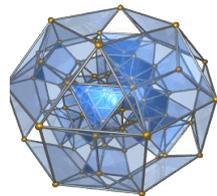


Nonlocality in many-body quantum systems

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↖ 2-body

Polytope	\mathbb{P}_n	Lower order correlators	\mathbb{P}_2		
Dimension	$3^n - 1$		$2n^2$		
Vertices	2^{2n}		2^{2n}		

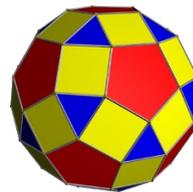
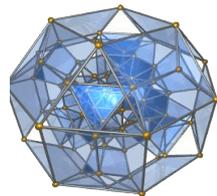


Nonlocality in many-body quantum systems

- Reducing the mathematical complexity

↖ 2-body

Polytope	\mathbb{P}_n Lower order correlators	\mathbb{P}_2 Action of a symmetry group	
Dimension	$3^n - 1$	$2n^2$	
Vertices	2^{2n}	2^{2n}	

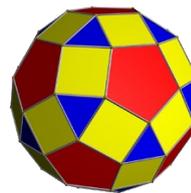
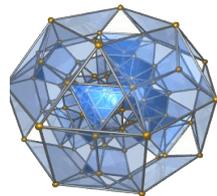


Nonlocality in many-body quantum systems

- Reducing the mathematical complexity

\swarrow 2-body \swarrow Cyclic group

Polytope	\mathbb{P}_n <i>Lower order correlators</i>	\mathbb{P}_2 <i>Action of a symmetry group</i>	
Dimension	$3^n - 1$	$2n^2$	
Vertices	2^{2n}	2^{2n}	



[JT, A. B. Sainz, T. Vértesi, M. Lewenstein, A. Acín, R. Augusiak
J. Phys. A: Math. Theor. **47** 424024 (2014)]

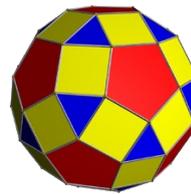
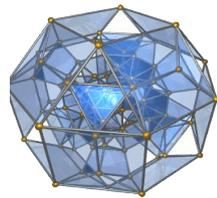


Nonlocality in many-body quantum systems

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Polytope	\mathbb{P}_n <i>Lower order correlators</i>	\mathbb{P}_2 <i>Action of a symmetry group</i>		
Dimension	$3^n - 1$	$2n^2$		
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↖ 2-body
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Science **344** 1256 (2014)]

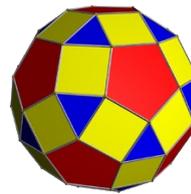
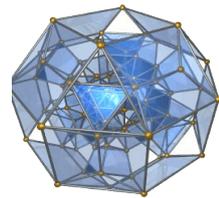


Nonlocality in many-body quantum systems

- Reducing the mathematical complexity

Polytope	\mathbb{P}_n <i>Lower order correlators</i>	\mathbb{P}_2 <i>Action of a symmetry group</i>	\mathbb{P}_2^S
Dimension	$3^n - 1$	$2n^2$	5
Vertices	2^{2n}	2^{2n}	$\binom{n+3}{3}$

\nearrow 2-body \nearrow Cyclic group
 \nearrow Symmetric group



$\text{CH}(\varphi(\mathbb{T}))$



$\nearrow \ell$

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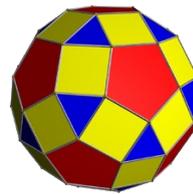
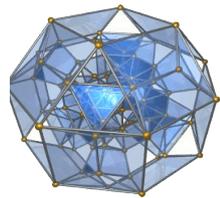
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Science **344** 1256 (2014)]



Nonlocality in many-body quantum systems

- Reducing the mathematical complexity

		2-body	Cyclic group	Deterministic Local Strategies Parametrization
Polytope	\mathbb{P}_n Lower order correlators	\mathbb{P}_2 Action of a symmetry group	\mathbb{P}_2^S Extremality analysis	
Dimension	$3^n - 1$	$2n^2$	5	
Vertices	2^{2n}	2^{2n}	$\binom{n+3}{3}$	



$\text{CH}(\varphi(\mathbb{T}))$



Parameter space \mathbb{T}



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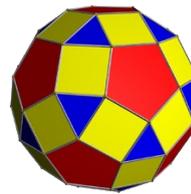
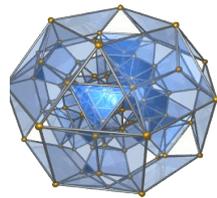
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Nonlocality in many-body quantum systems

- Reducing the mathematical complexity

		2-body	Cyclic group Symmetric group	Deterministic Local Strategies Parametrization
Polytope	\mathbb{P}_n Lower order correlators	\mathbb{P}_2 Action of a symmetry group	\mathbb{P}_2^S Extremality analysis	\mathbb{P}_2^S
Dimension	$3^n - 1$	$2n^2$	5	5
Vertices	2^{2n}	2^{2n}	$\binom{n+3}{3}$	$2(n^2 + 1)$



$CH(\varphi(\mathbb{T}))$



$CH(\varphi(\partial\mathbb{T}))$



Parameter space \mathbb{T}



[JT, A. B. Sainz, T. Vértesi, M. Lewenstein, A. Acín, R. Augusiak
J. Phys. A: Math. Theor. **47** 424024 (2014)]

[JT, R. Augusiak, A. B. Sainz, T. Vértesi, M. Lewenstein, A. Acín
Science **344** 1256 (2014)]



Some applications



Some applications

PRL **119**, 230402 (2017)

PHYSICAL REVIEW LETTERS

week ending
8 DECEMBER 2017

Bounding the Set of Classical Correlations of a Many-Body System

Matteo Fadel^{1,*} and Jordi Tura^{2,3,†}

We present a method to certify the presence of Bell correlations in experimentally observed statistics, and to obtain new Bell inequalities. Our approach is based on relaxing the conditions defining the set of correlations obeying a local hidden variable model, yielding a convergent hierarchy of semidefinite programs (SDP's). Because the size of these SDP's is independent of the number of parties involved, this technique allows us to characterize correlations in many-body systems. As an example, we illustrate our method with the experimental data presented in *Science* **352**, 441 (2016).

DOI: 10.1103/PhysRevLett.119.230402



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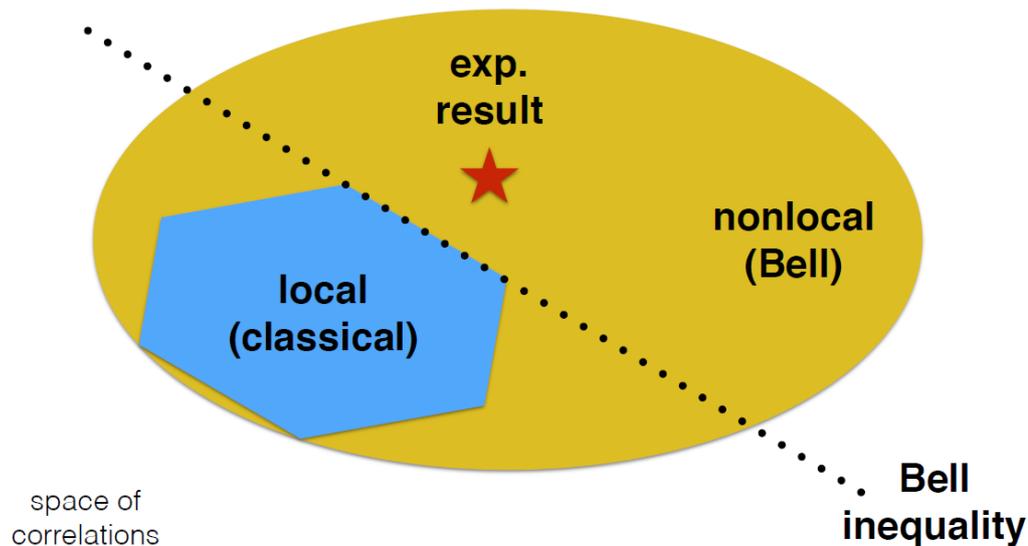


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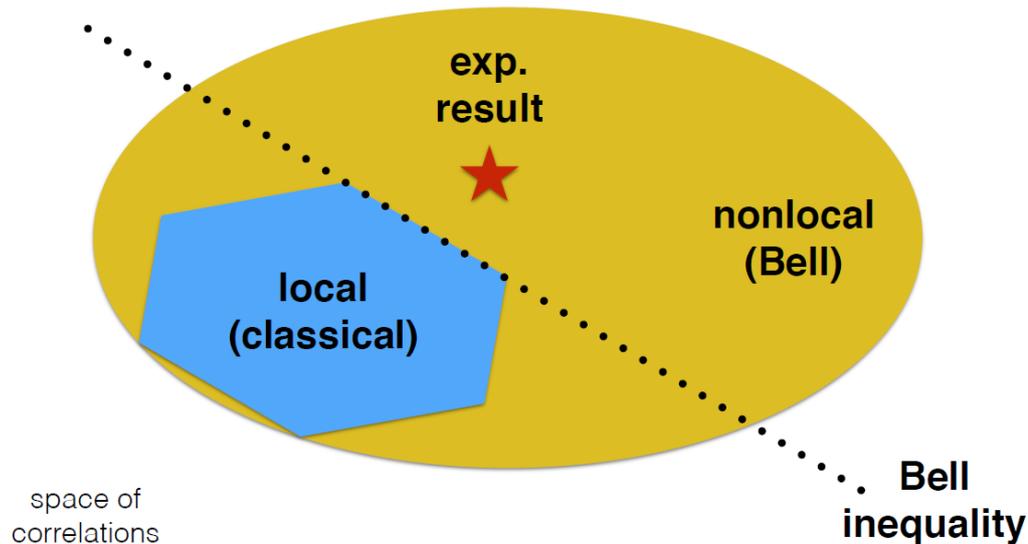
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- *Certify Bell correlations from experiments*



Some applications

PRL 119, 230402 (2017)

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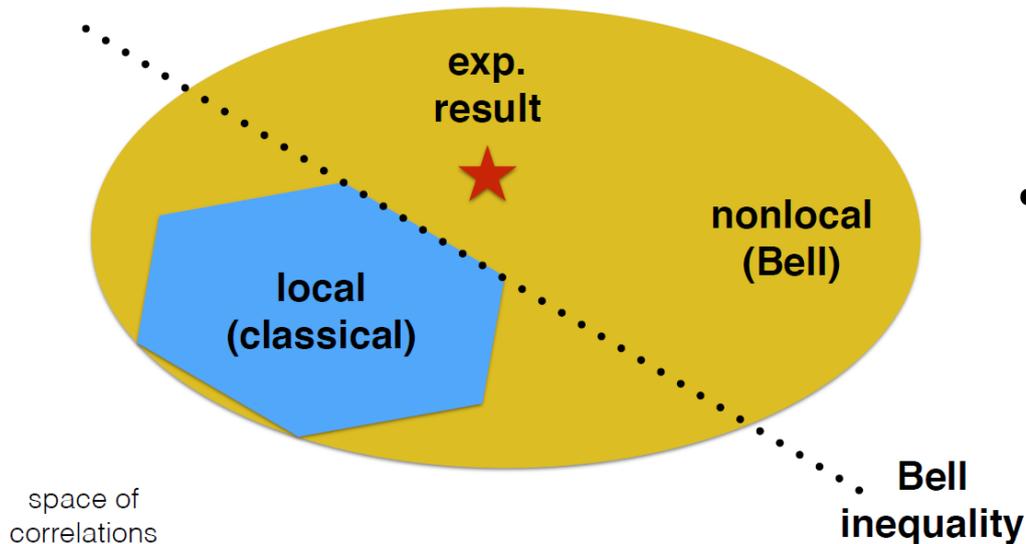
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- *Certify Bell correlations from experiments*
- *Find new Bell Inequalities*



Some applications



Some applications

We consider Bell inequalities of the form

$$\sum_k \sum_{j_1 \leq \dots \leq j_k} \alpha_{j_1 \dots j_k} S_{j_1 \dots j_k} + \beta C \geq 0$$



Some applications

We consider Bell inequalities of the form

$$\sum_k \sum_{j_1 \leq \dots \leq j_k} \alpha_{j_1 \dots j_k} \mathcal{S}_{j_1 \dots j_k} + \beta_C \geq 0$$

with $\mathcal{S}_{j_1 \dots j_k} = \sum_{\substack{i_1, \dots, i_k = 1 \\ \text{all } i\text{'s different}}}^N \langle \mathcal{M}_{j_1}^{(i_1)} \dots \mathcal{M}_{j_k}^{(i_k)} \rangle$



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Dimension depends on

- Order of the correlators
- #Measurements



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Dimension depends on

- Order of the correlators
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Does NOT depend on



Some applications

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Dimension depends on

- Order of the correlators
- #Measurements
- #outcomes

Does NOT depend on

- #Parties



Previous results



Previous results

Science 344, 1256 (2014)

Detecting nonlocality in many-body quantum states

J. Tura,¹ R. Augusiak,^{1*} A. B. Sainz,¹ T. Vértesi,² M. Lewenstein,^{1,3} A. Acín^{1,3}



Previous results

Science 344, 1256 (2014)

Detecting nonlocality in many-body quantum states

Example
$$-2S_0 + \frac{1}{2}S_{00} - S_{01} + \frac{1}{2}S_{11} + 2N \geq 0$$

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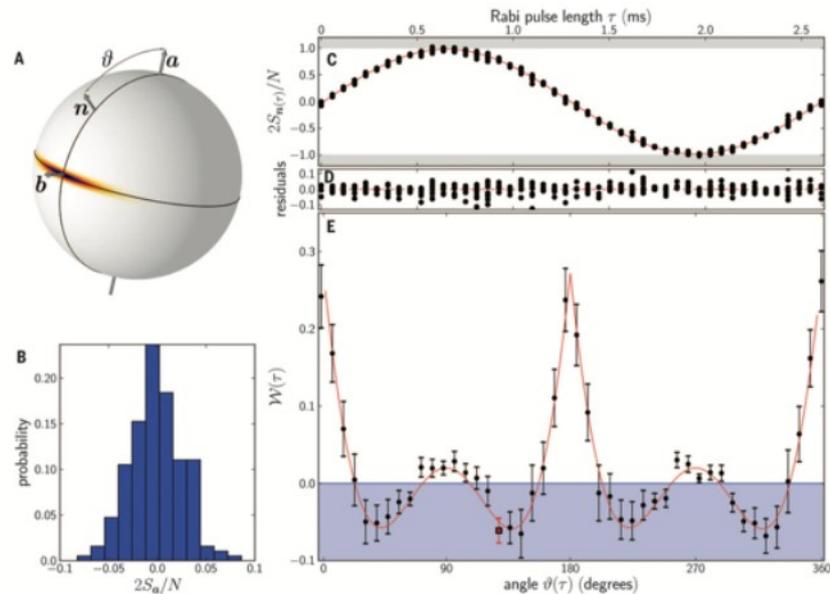
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Bell correlations in a Bose-Einstein condensate

Roman Schmied,^{1*} Jean-Daniel Bancal,^{2,4*} Baptiste Allard,^{1*} Matteo Fadel,¹ Valerio Scarani,^{2,3} Philipp Treutlein,^{1,†} Nicolas Sangouard^{4,†}

$$\hat{W} = - \left| \frac{\hat{S}_n}{N/2} \right| + (a \cdot n)^2 \frac{\hat{S}_a^2}{N/4} + 1 - (a \cdot n)^2$$



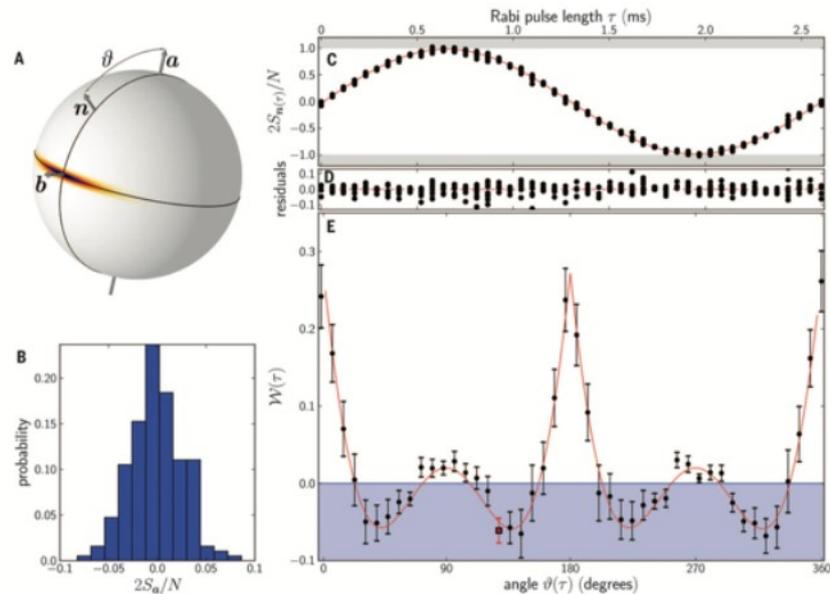
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PRL 118, 140401 (2017)

PHYSICAL REVIEW LETTERS

week ending
7 APRIL 2017



Bell Correlations in Spin-Squeezed States of 500 000 Atoms

Nils J. Engelsen, Rajiv Krishnakumar, Onur Hosten, and Mark A. Kasevich*



MAX PLANCK INSTITUTE
OF QUANTUM OPTICS

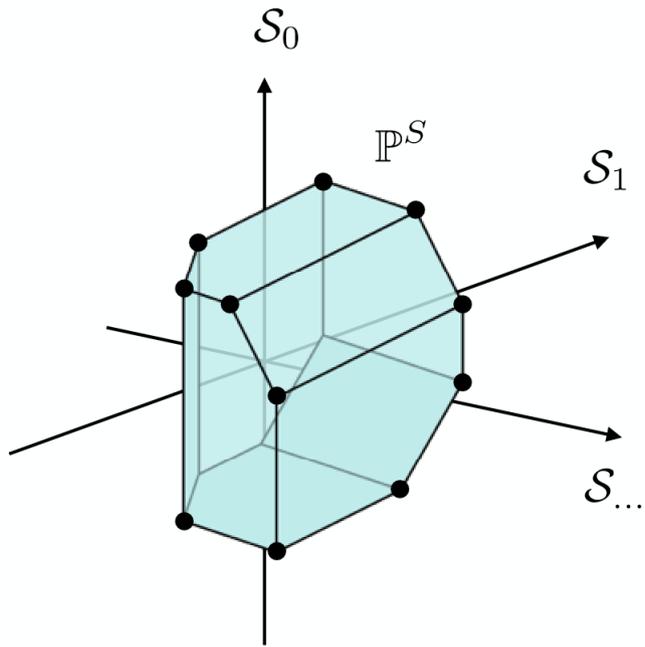
The Local polytope

- *Solving for a few values of N ...*



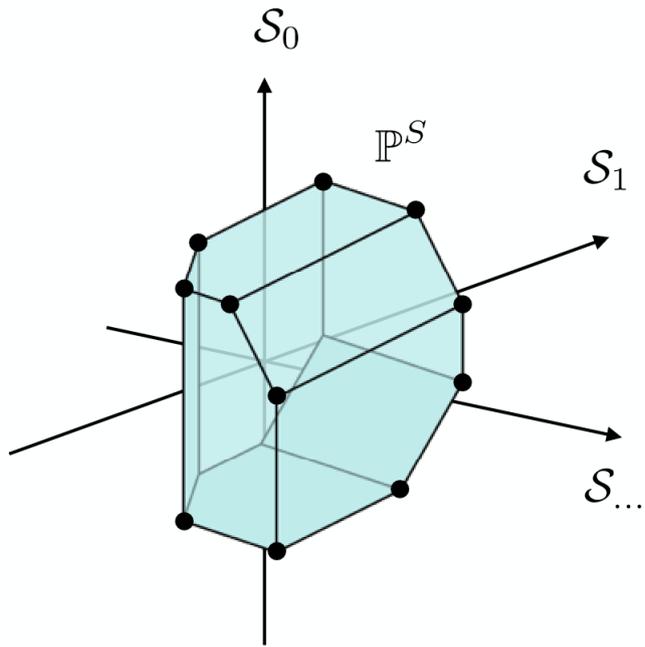
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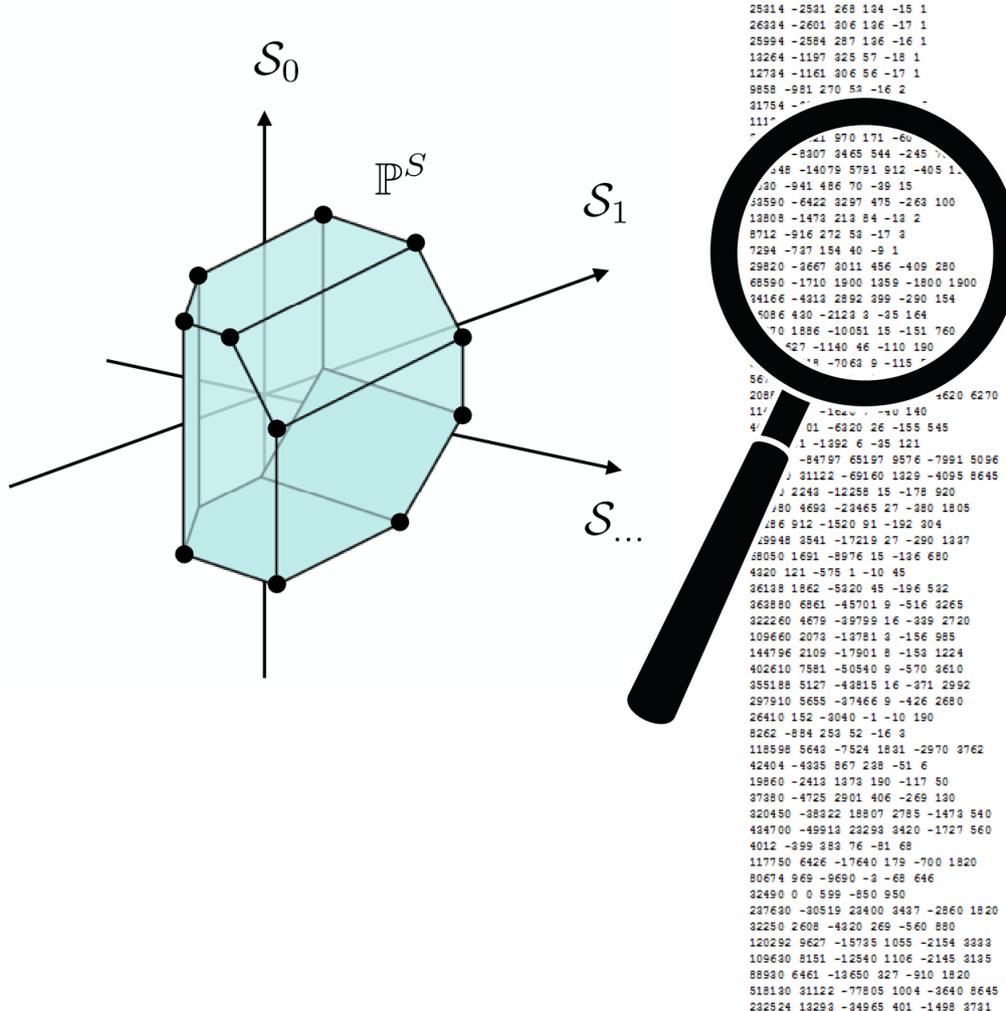


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25994 -2584 287 136 -16 1
13264 -1197 325 57 -18 1
12734 -1161 306 56 -17 1
9888 -981 270 53 -16 2
31754 -3128 867 167 -51 6
11138 -1083 308 57 -18 2
29470 -3021 970 171 -60 10
78560 -8307 3465 544 -245 70
128548 -14079 5791 912 -405 112
7830 -941 486 70 -39 15
53590 -6422 3297 475 -263 100
13808 -1473 213 84 -13 2
8712 -916 272 53 -17 3
7294 -737 154 40 -9 1
29820 -3667 3011 456 -409 280
68590 -1710 1900 1359 -1800 1900
34166 -4313 2892 399 -290 154
16086 430 -2123 3 -35 164
76470 1886 -10051 15 -151 760
7980 627 -1140 46 -110 190
53706 1418 -7063 9 -115 544
5674 138 -743 1 -11 56
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11490 437 -1620 7 -40 140
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661284 -84797 65197 9576 -7991 5096
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178980 4699 -23465 27 -380 1805
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129948 3541 -17219 27 -290 1337
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109660 2073 -13781 3 -156 985
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402610 7581 -50540 9 -570 3610
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297910 5655 -37466 9 -426 2680
26410 152 -3040 -1 -10 190
8262 -884 253 52 -16 3
118598 5643 -7524 1831 -2970 3762
42404 -4335 867 238 -51 6
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434700 -49913 23293 3420 -1727 560
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117750 6426 -17640 179 -700 1820
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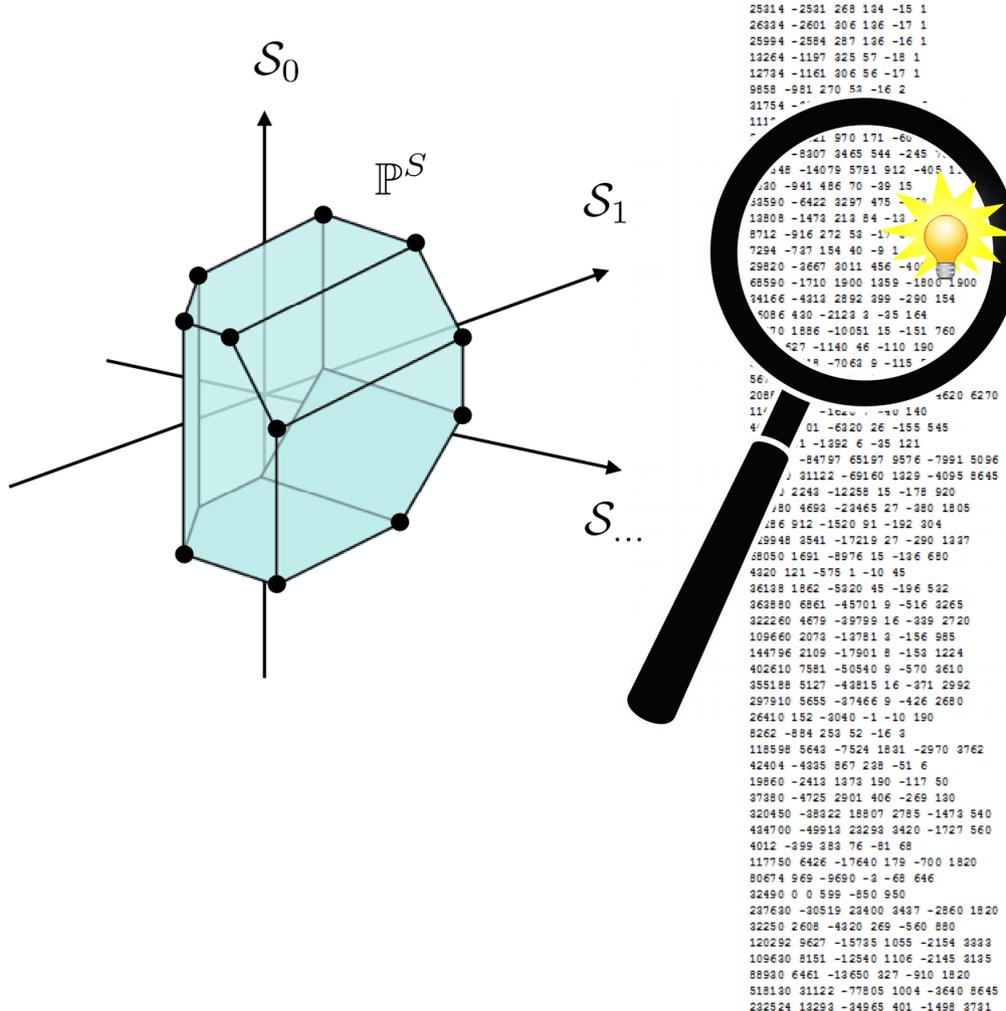
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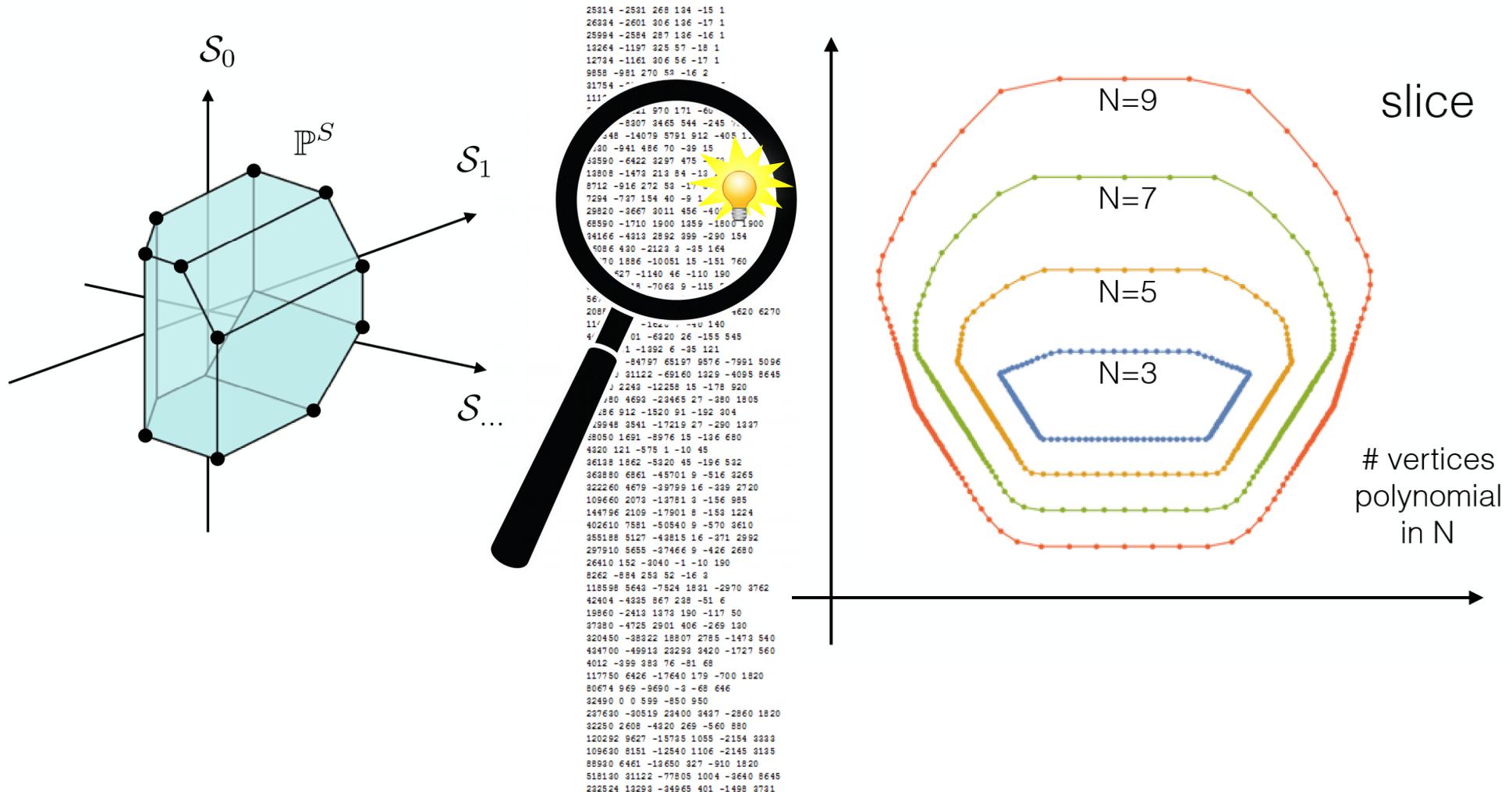
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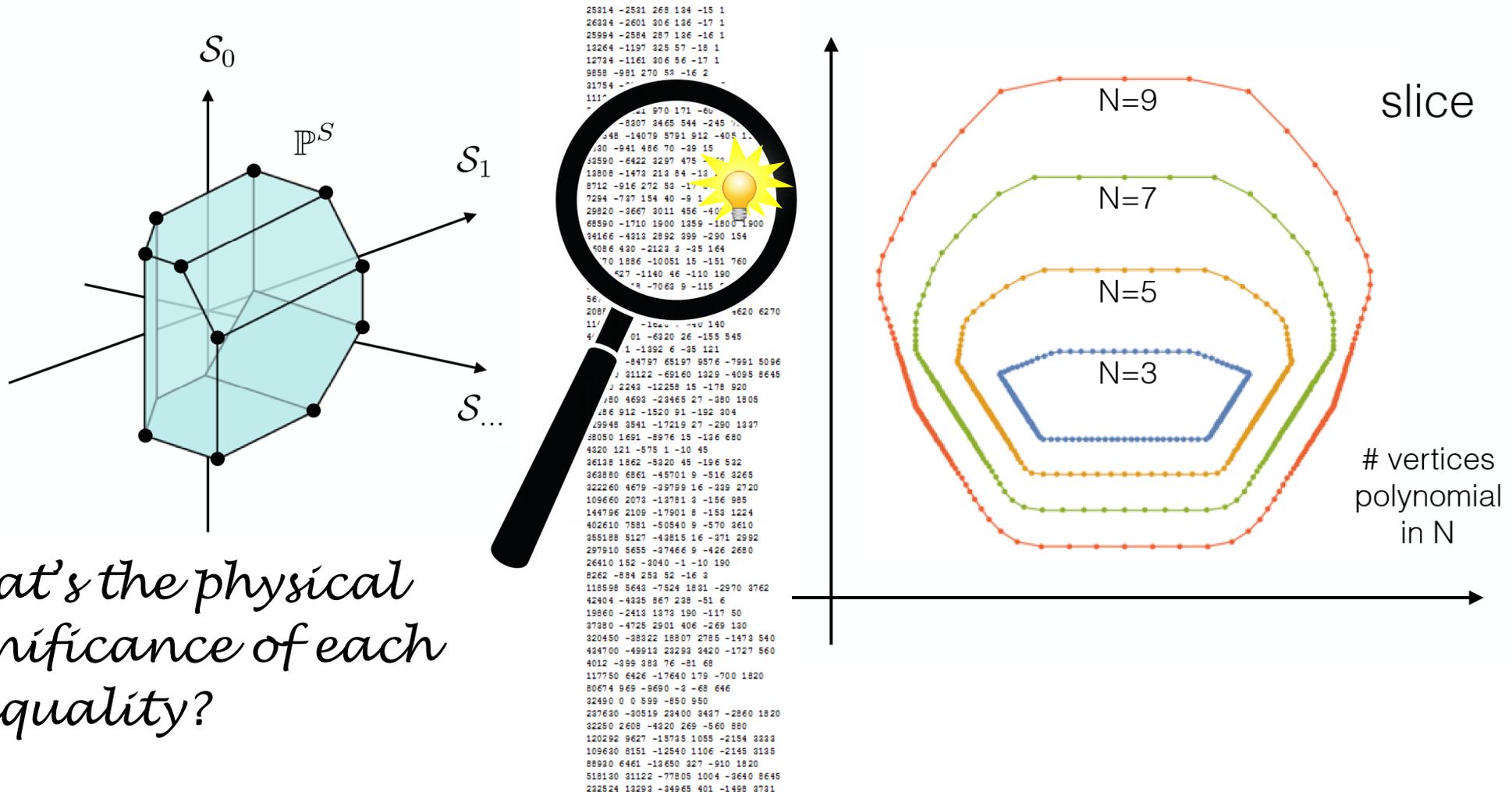
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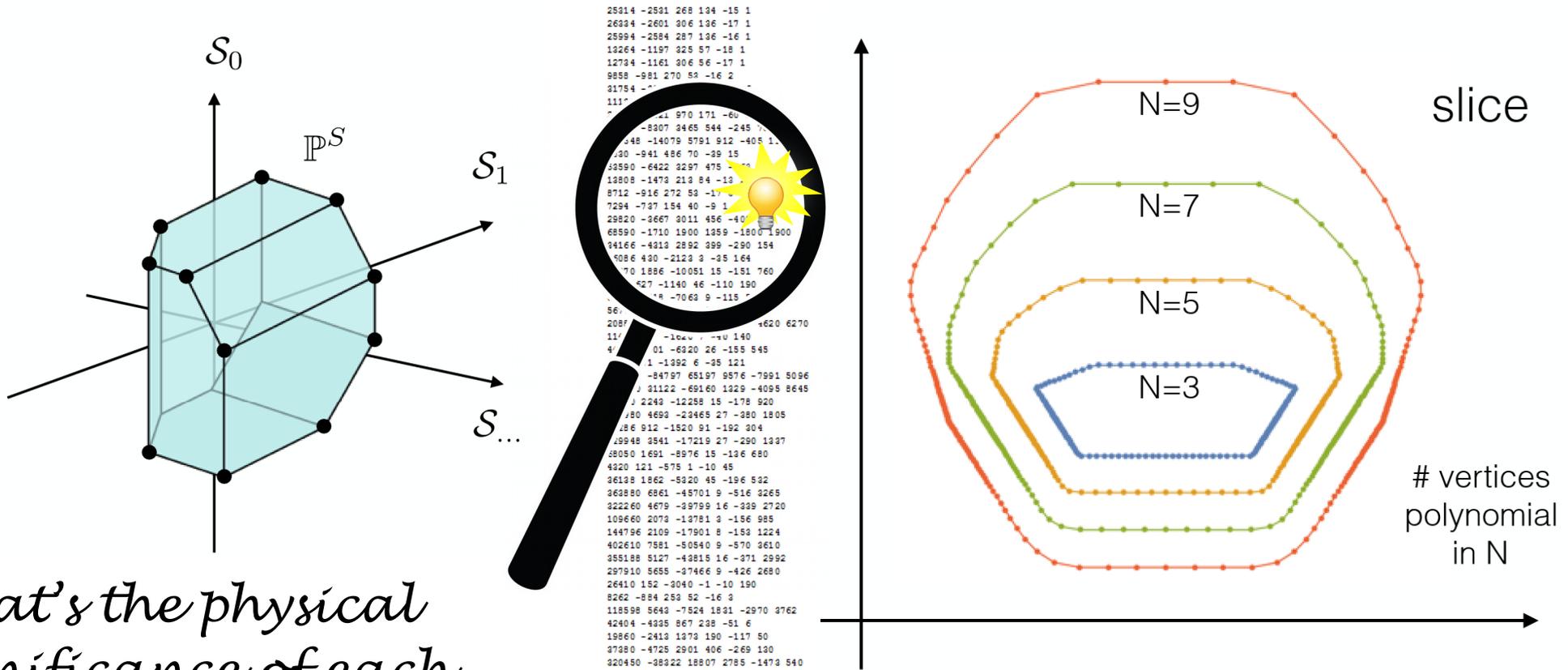
The Local polytope

- Solving for a few values of N ...



The Local polytope

- Solving for a few values of N ...



What's the physical significance of each inequality?
 Are they equally relevant?



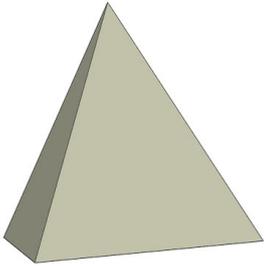
Bounding the LHV set

- *Main Observation*



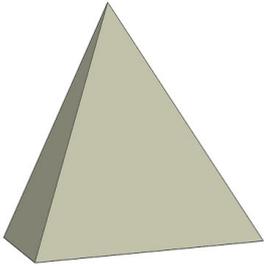
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Bounding the LHVVM set

- Main Observation

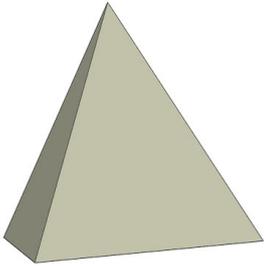


As the system becomes larger



Bounding the LHV set

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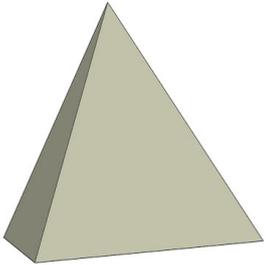


As the system becomes larger



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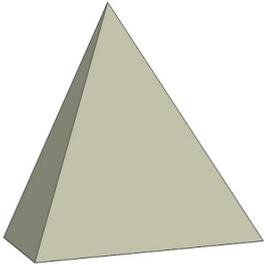
As the system becomes larger

Where does this extra structure come from?



Bounding the LHV set

- Main Observation



As the system becomes larger

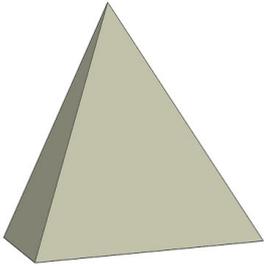
Where does this extra structure come from?

*Local
Deterministic
Strategy view:*



Bounding the LHV set

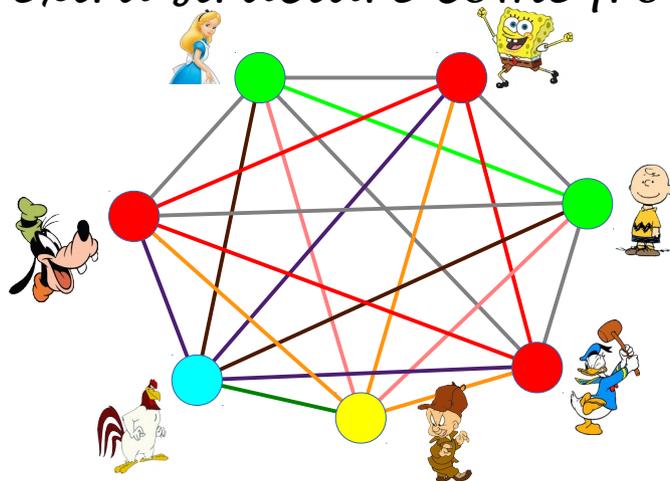
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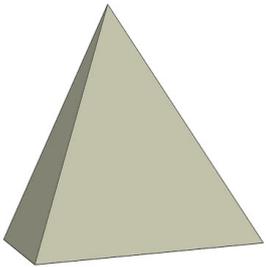
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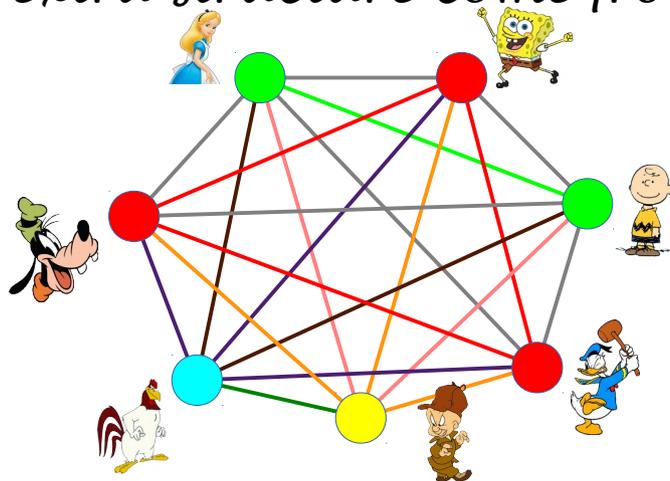
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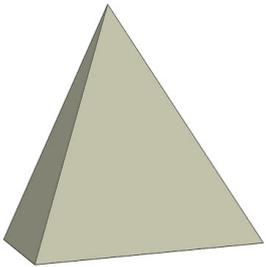


Permutational Invariance:



Bounding the LHV set

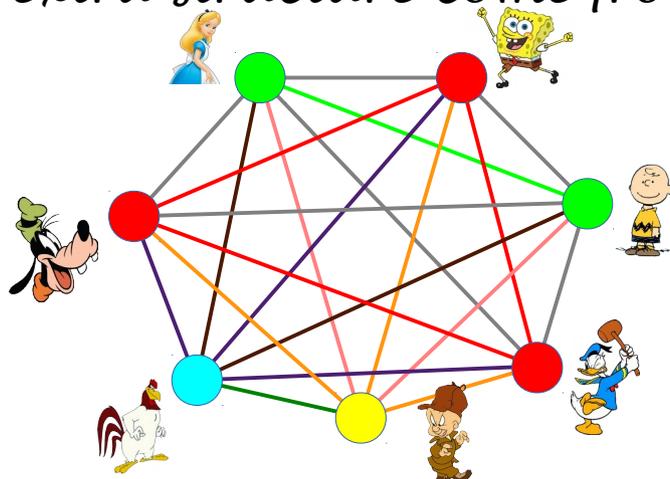
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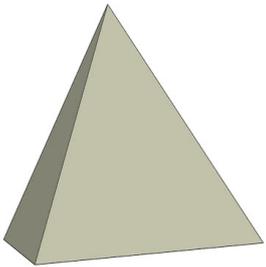
Permutational Invariance:

*Only amount of each color
becomes relevant*



Bounding the LHV set

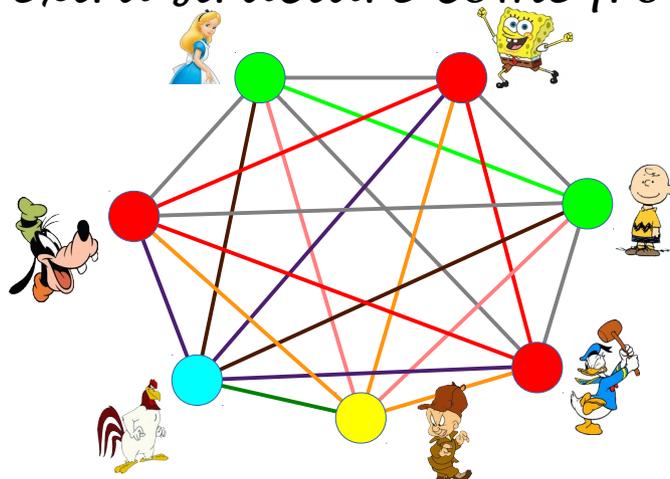
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Strategy view:



Permutational Invariance:

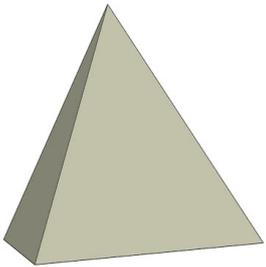
Only amount of each color
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$$x_i \geq 0$$



Bounding the LHV set

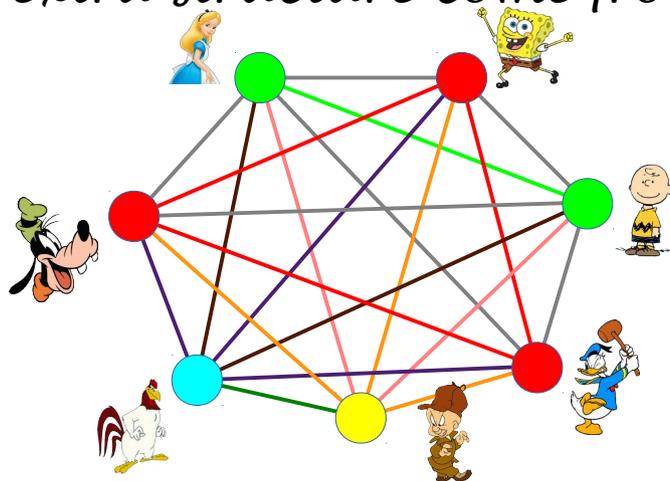
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As the system becomes larger

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Local
Deterministic
Strategy view:



Permutational Invariance:

Only amount of each color becomes relevant

$$x_i \geq 0 \quad \sum_i x_i = N$$



Bounding the LHV set

- *Algebraic structure at every LDS*



Bounding the LHV set

- Algebraic structure at every LDS

$$S_{kl} = S_k \cdot S_l - Z_{kl}$$



Bounding the LHV set

- Algebraic structure at every LDS

$$S_{kl} = S_k \cdot S_l - Z_{kl}$$

$$\begin{pmatrix} N \\ S_1 \\ S_0 \\ Z_{01} \end{pmatrix} = 2H^{\otimes 2} \begin{pmatrix} x_1 \\ x_2 \\ x_3 \\ x_4 \end{pmatrix}$$



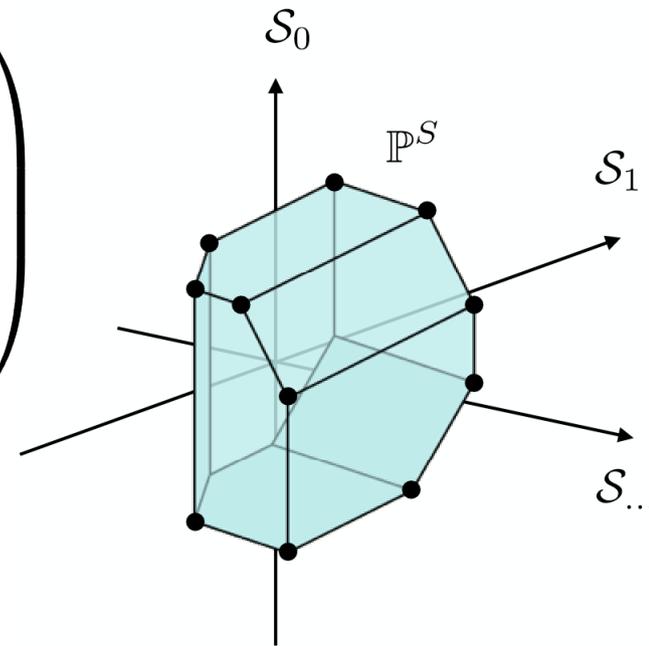
Bounding the LHV set

- Algebraic structure at every LDS

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$$\mathbb{P}^S = \text{CH} \left\{ \vec{S}(\vec{x}) \text{ s.t. } \sum_i x_i = N, x_i \in \mathbb{Z}_{\geq 0} \right\}$$

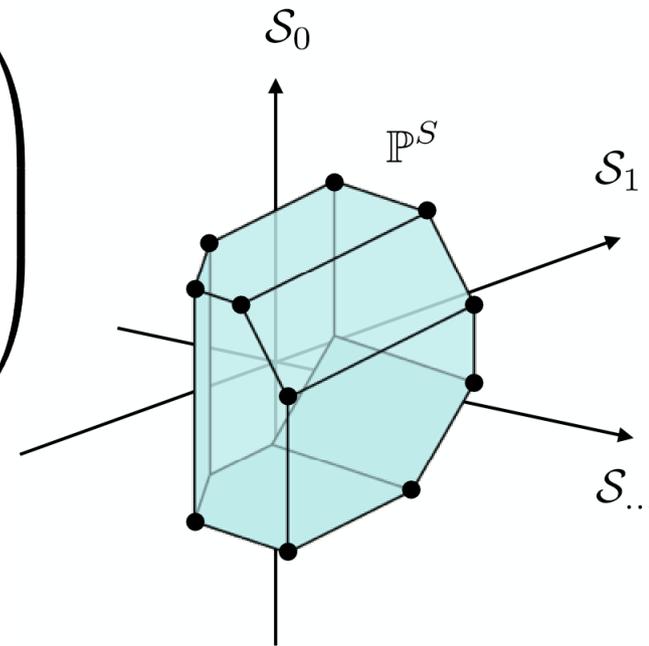


Bounding the LHV set

- Algebraic structure at every LDS

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$$\mathbb{P}^S = \text{CH} \left\{ \vec{S}(\vec{x}) \text{ s.t. } \sum_i x_i = N, x_i \in \mathbb{Z}_{\geq 0} \right\}$$

Goal: Define a manifold interpolating the vertices of the symmetric 2-body polytope, and compute its convex hull



First relaxation

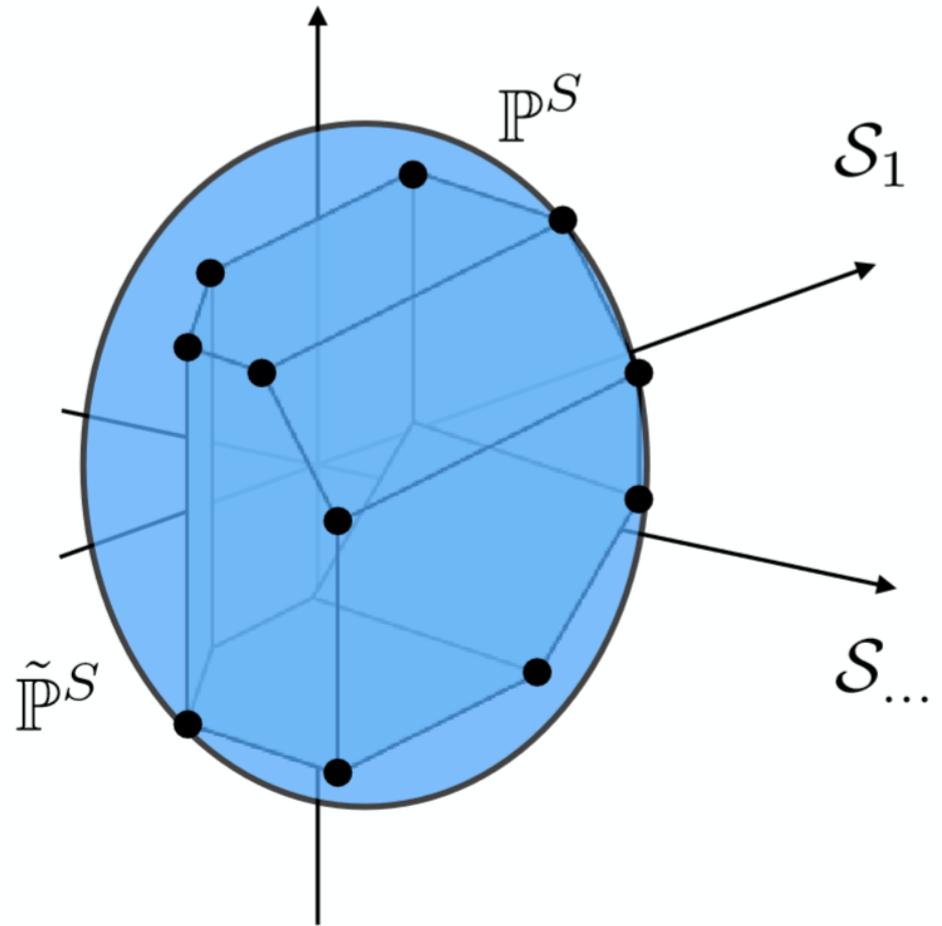


First relaxation

$$\mathbb{P}^S = \text{CH} \left\{ \vec{S}(\vec{x}) \text{ s.t. } \sum_i x_i = N, x_i \in \mathbb{Z}_{\geq 0} \right\}$$

$$\tilde{\mathbb{P}}^S = \text{CH} \left\{ \vec{S}_K(\vec{x}) \text{ s.t. } \sum_i x_i = N, x_i \in \mathbb{R}_{\geq 0} \right\}$$

$$\tilde{\mathbb{P}}^S \supseteq \mathbb{P}^S$$



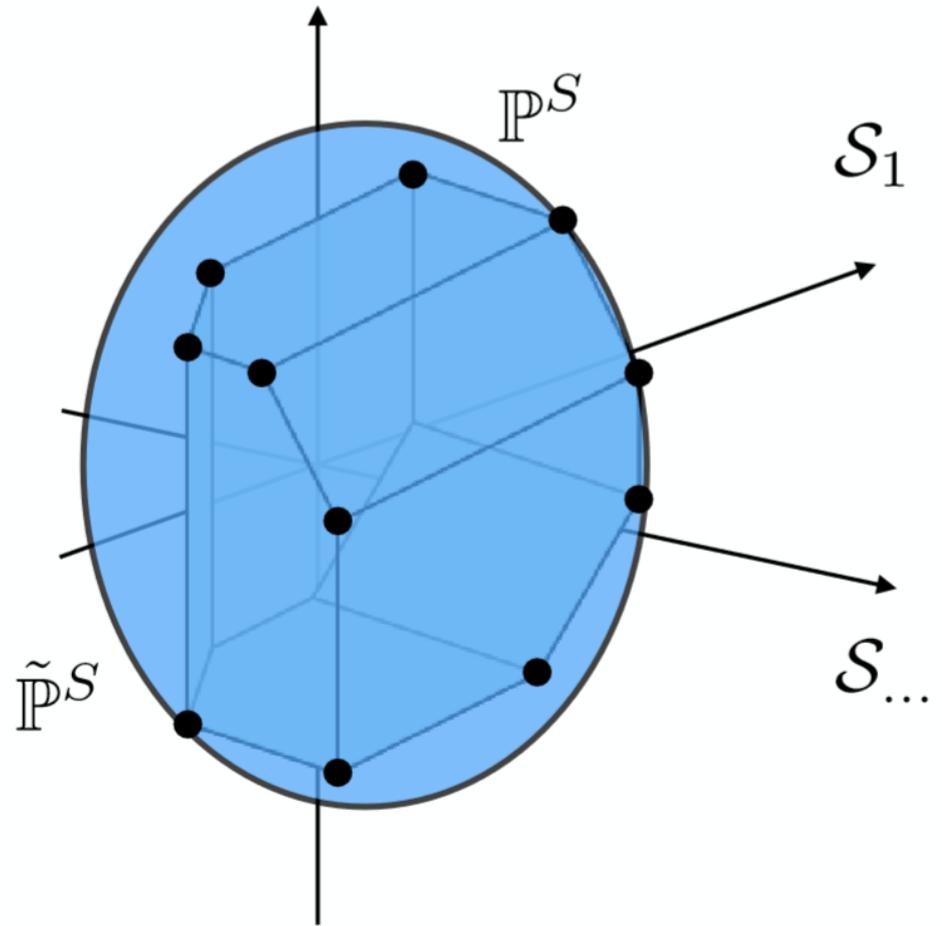
First relaxation

$$\mathbb{P}^S = \text{CH} \left\{ \vec{\mathcal{S}}(\vec{x}) \text{ s.t. } \sum_i x_i = N, x_i \in \mathbb{Z}_{\geq 0} \right\}$$

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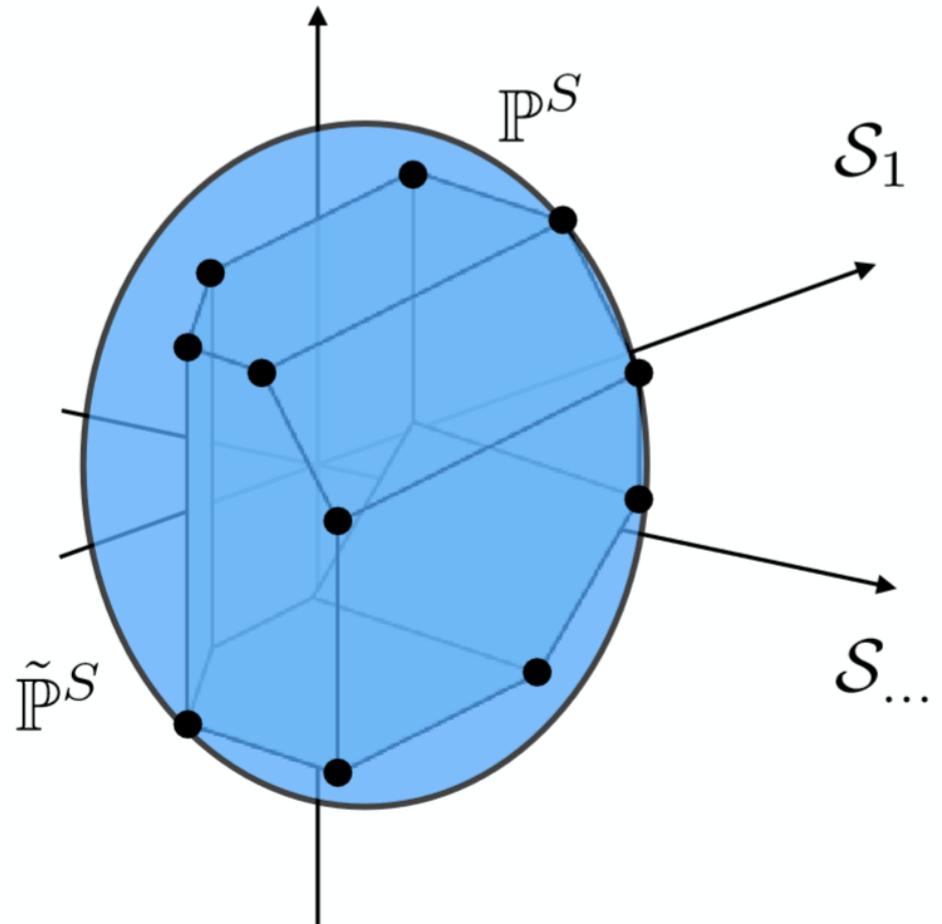
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Computing the convex hull of a semialgebraic set is NP-hard



Second relaxation



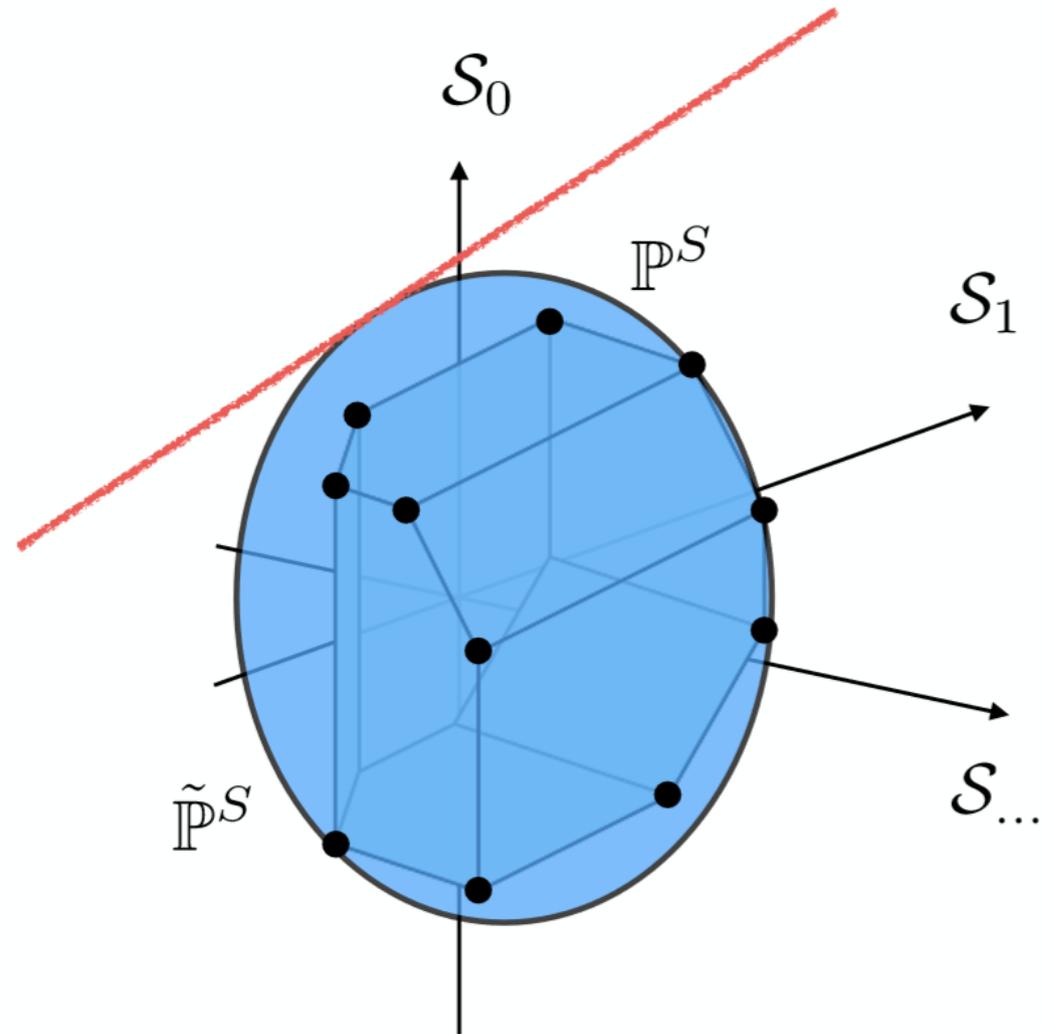
Second relaxation $l(\vec{\mathcal{S}})$

$\tilde{\mathbb{P}}^{\mathcal{S}}$ is defined by
$$\begin{cases} f_i(\vec{\mathcal{S}}_K) = 0 \\ g_j(\vec{\mathcal{S}}_K) \geq 0 \end{cases}$$

ansatz:
$$l(\vec{\mathcal{S}}) = \sum_{i=0}^m g_i(\vec{\mathcal{S}}) \sigma_i(\vec{\mathcal{S}})$$

$\sigma_i(\vec{\mathcal{S}})$ sos polynomials

NOTE: l is positive in $\tilde{\mathbb{P}}^{\mathcal{S}}$



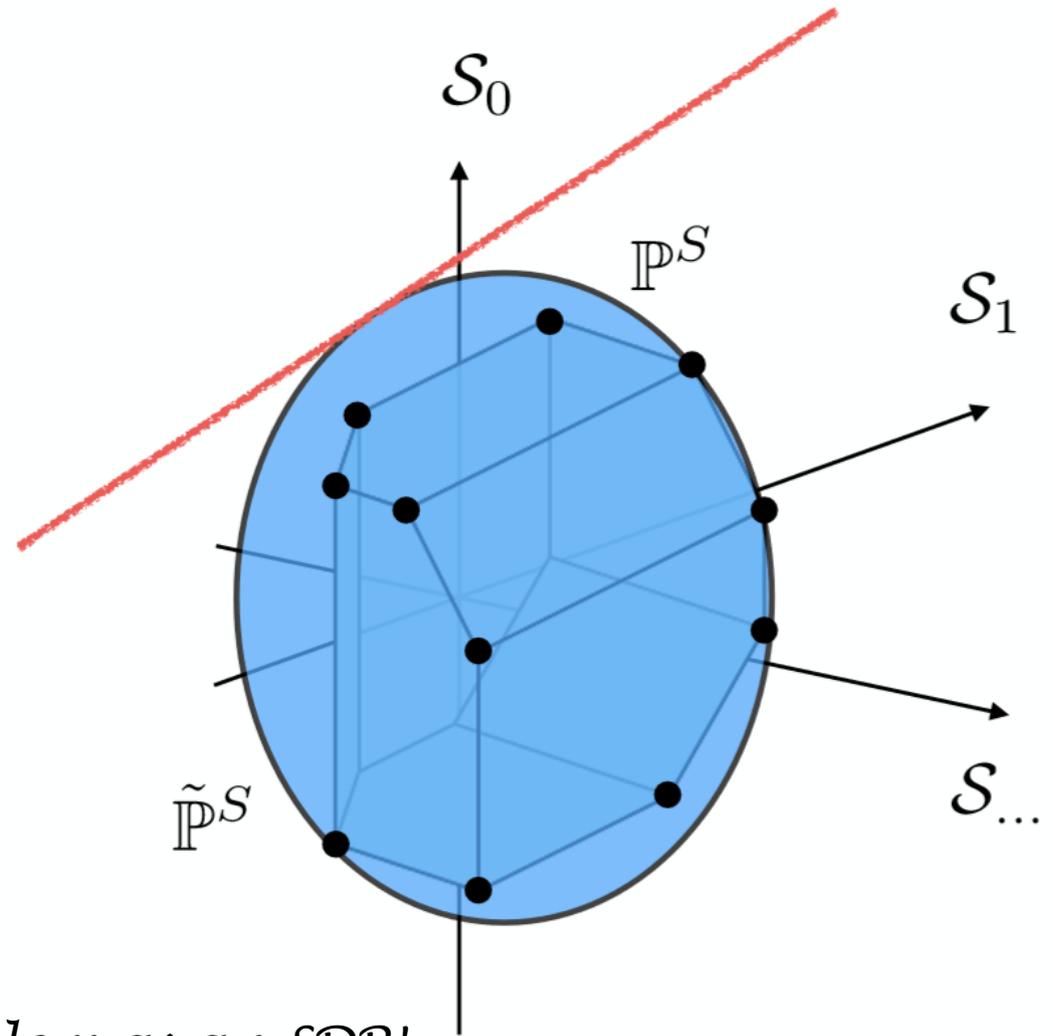
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Solve the sos representation problem as an SDP!



The SDP



The SDP

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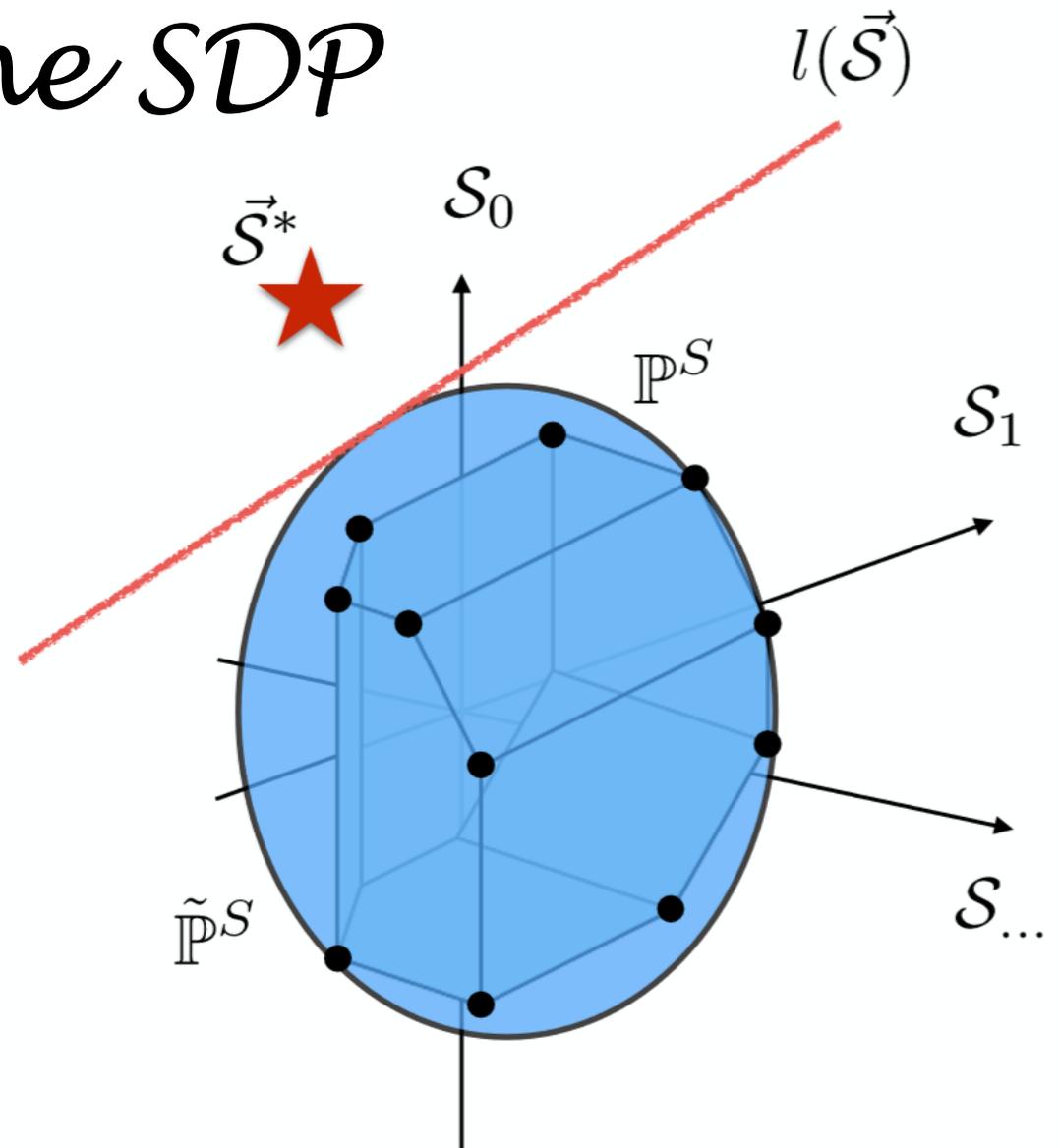
$$= \Gamma \cdot G$$

GOAL: prove $l(\vec{\mathcal{S}}^*) < 0$
for some $\sigma_i(\vec{\mathcal{S}})$

$$G \succeq 0 \quad \Gamma = \begin{pmatrix} 1 & \mathcal{S}_0 & \mathcal{S}_1 & \mathcal{S}_{00} & \dots \\ \mathcal{S}_0 & \mathcal{S}_0^2 & \mathcal{S}_0 \mathcal{S}_1 & \mathcal{S}_0 \mathcal{S}_{00} & \dots \\ \mathcal{S}_1 & \mathcal{S}_1 \mathcal{S}_0 & \mathcal{S}_1^2 & \mathcal{S}_1 \mathcal{S}_{00} & \dots \\ \dots & \dots & \dots & \dots & \dots \end{pmatrix}$$

SDP: $\Gamma \succeq 0$

subject to: $\mathcal{S}_0 = \mathcal{S}_0^*$, $\mathcal{S}_1 = \mathcal{S}_1^*$, ...



The SDP

$$l(\vec{\mathcal{S}}) = \sum_{i=0}^m g_i(\vec{\mathcal{S}}) \sigma_i(\vec{\mathcal{S}})$$

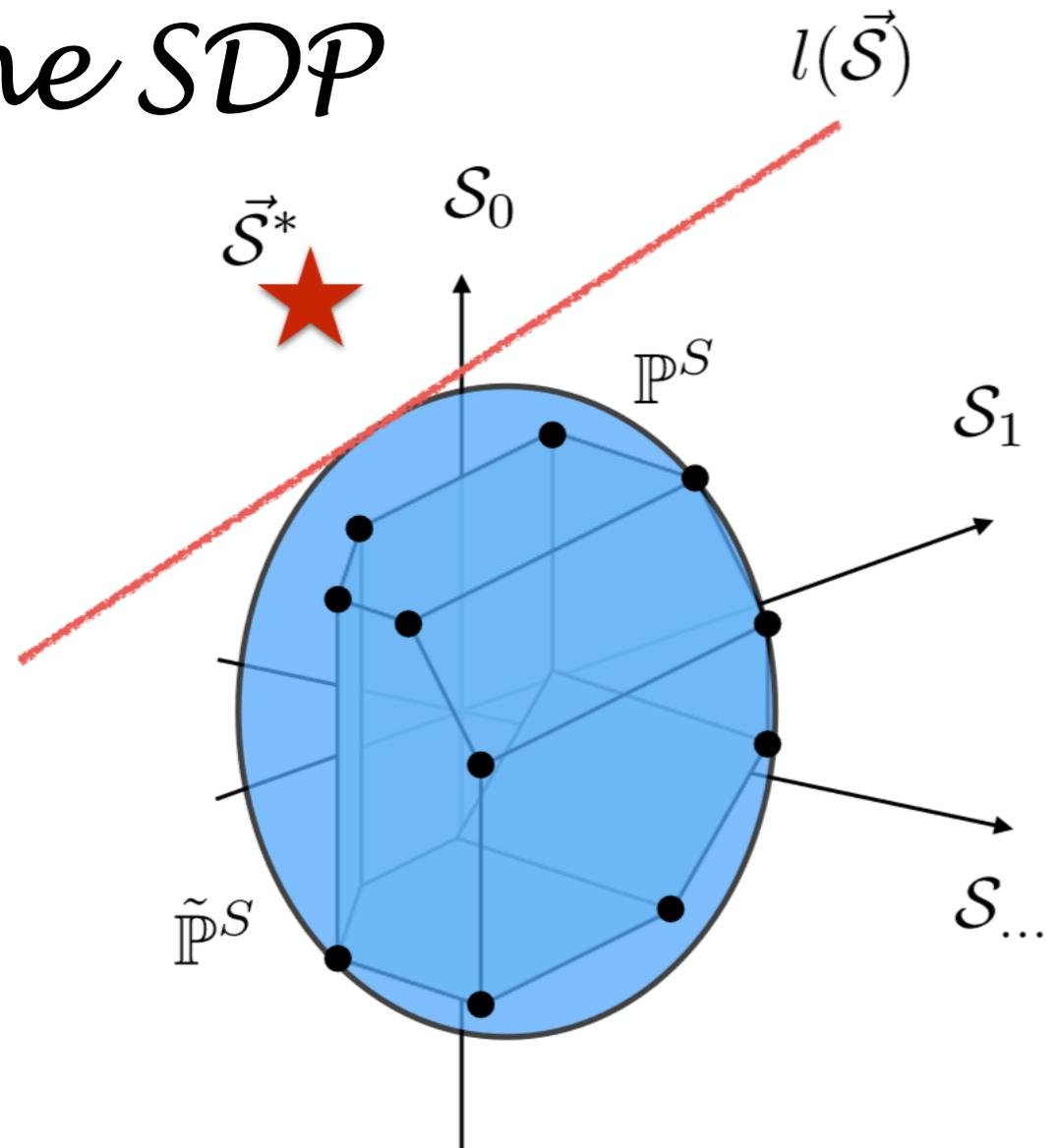
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(Experimental data point)



Results



Results

- *Testing all BI of a certain form with a single SDP*



Results

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- *If the experimental point is sufficiently nonlocal, the SDP outputs the Bell inequality that is violated, with a proof of its classical bound*



Results

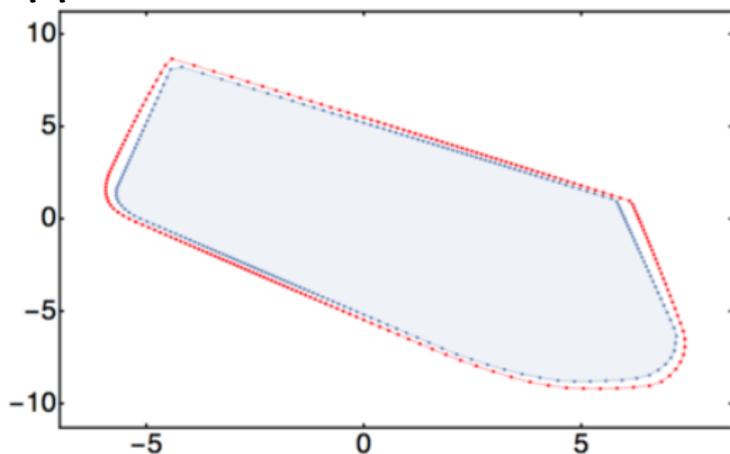
- *Testing all BI of a certain form with a single SDP*
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- *The complexity of the problem is independent of the system size*



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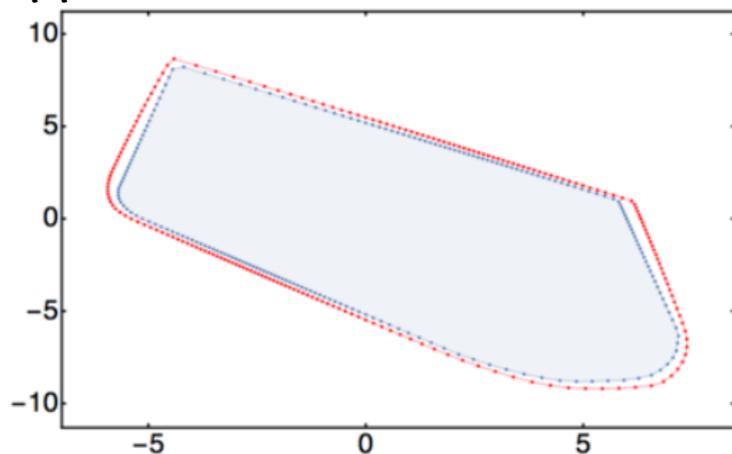
Approximation for $N=10$



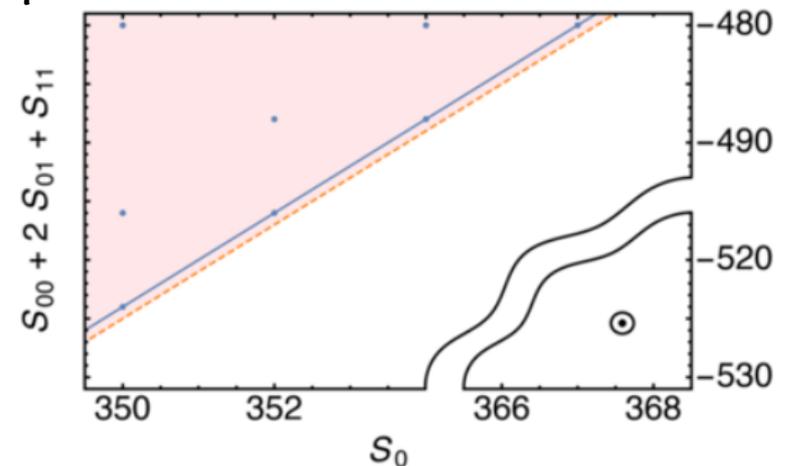
Results

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Approximation for $N=10$



Approximation with actual experimental data and $N=476$



Outlook



Outlook

- *Generalization to more outcomes, higher-order correlators...*



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Spin-nematic squeezing



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Polytope approach already impractical



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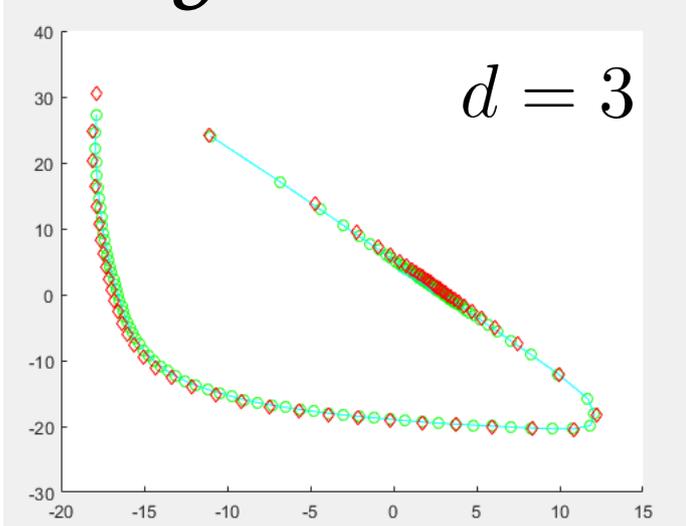
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[Ongoing work with A. Aloy and M. Fadel]



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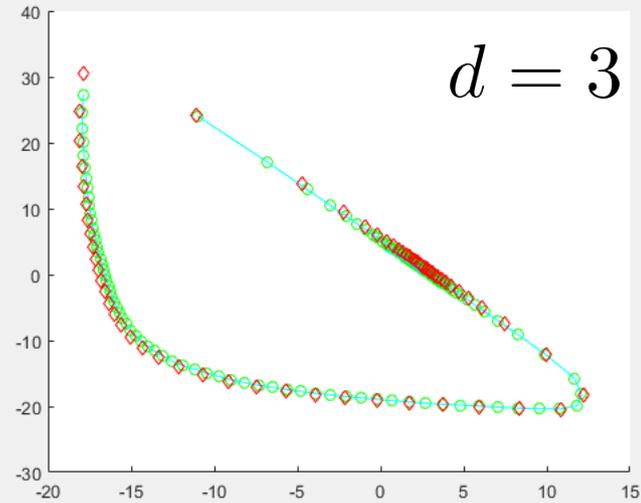
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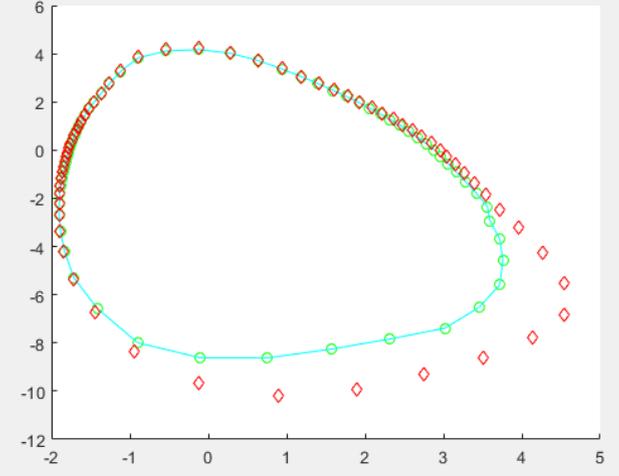
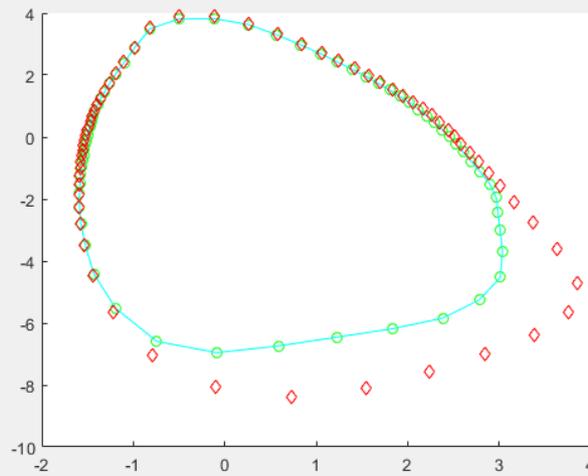
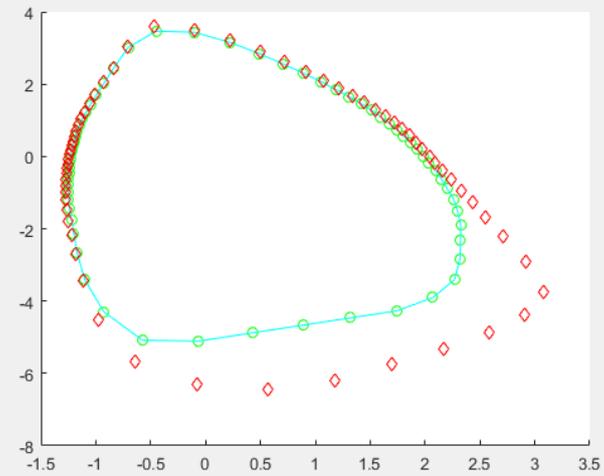
Outlook

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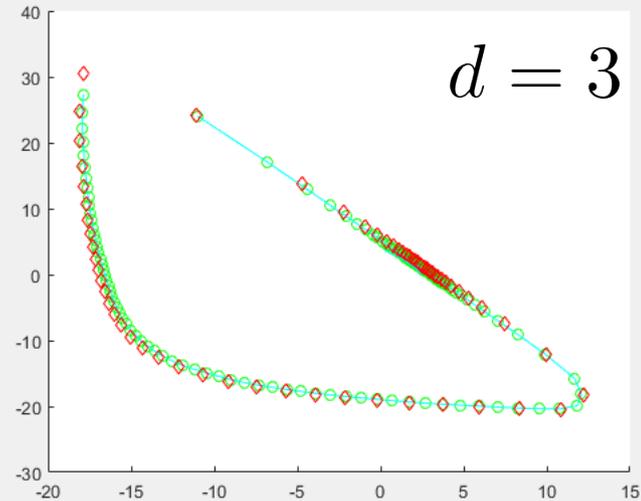
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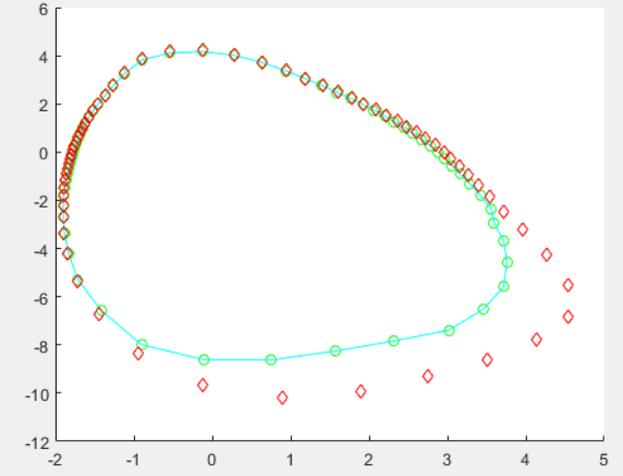
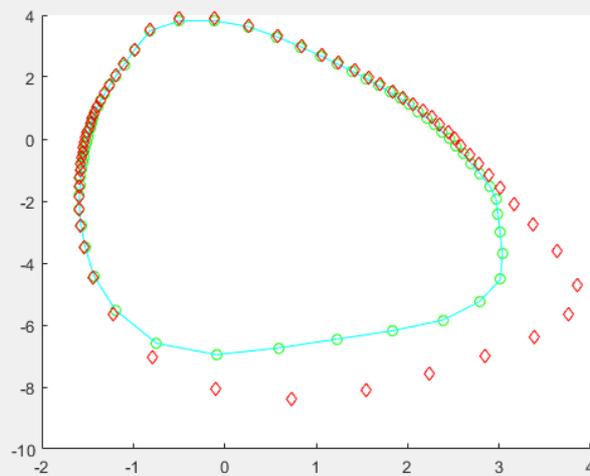
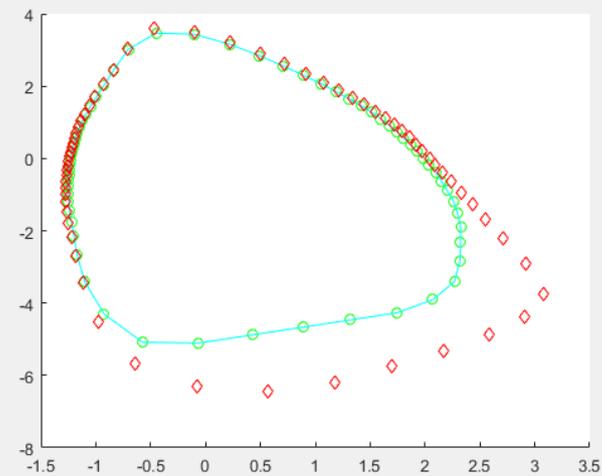
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Spin-nematic squeezing
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1	x_1	\square
2	$x_{01}x_{10} - x_{11}$	$\begin{array}{c} \square \\ \square \end{array} - \square$
3	$x_{001}x_{010}x_{100} - (x_{011}x_{100} + x_{010}x_{101} + x_{001}x_{110}) + 2x_{111}$	$\begin{array}{c} \square \\ \square \\ \square \end{array} - \begin{array}{c} \square \\ \square \end{array} + 2 \begin{array}{c} \square \\ \square \end{array}$
4	$x_{0001}x_{0010}x_{0100}x_{1000} + (x_{0110}x_{1001} + x_{0101}x_{1010} + x_{0011}x_{1100}) - (x_{0011}x_{0100}x_{1000} + x_{0010}x_{0101}x_{1000} + x_{0001}x_{0110}x_{1000} + x_{0010}x_{0100}x_{1001} + x_{0001}x_{0100}x_{1010} + x_{0001}x_{0010}x_{1100}) + 2(x_{0100}x_{1011} + x_{0010}x_{1101} + x_{0001}x_{1110} + x_{0111}x_{1000}) - 6x_{1111}$	$\begin{array}{c} \square \\ \square \\ \square \\ \square \end{array} + \begin{array}{c} \square \\ \square \end{array} - \begin{array}{c} \square \\ \square \\ \square \end{array} + 2 \begin{array}{c} \square \\ \square \\ \square \\ \square \end{array} - 6 \begin{array}{c} \square \\ \square \\ \square \\ \square \end{array}$

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Outlook



Outlook

- *Bell inequalities for Nonlocality depth*



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[arXiv.org](https://arxiv.org/abs/1802.09516) > [quant-ph](https://arxiv.org/abs/1802.09516) > [arXiv:1802.09516](https://arxiv.org/abs/1802.09516)

Quantum Physics

Bell correlations depth in many-body systems

[Flavio Baccari](#), [Jordi Tura](#), [Matteo Fadel](#), [Albert Aloy](#), [Jean-Daniel Bancal](#), [Nicolas Sangouard](#), [Maciej Lewenstein](#), [Antonio Acín](#), [Remigiusz Augusiak](#)

(Submitted on 26 Feb 2018)



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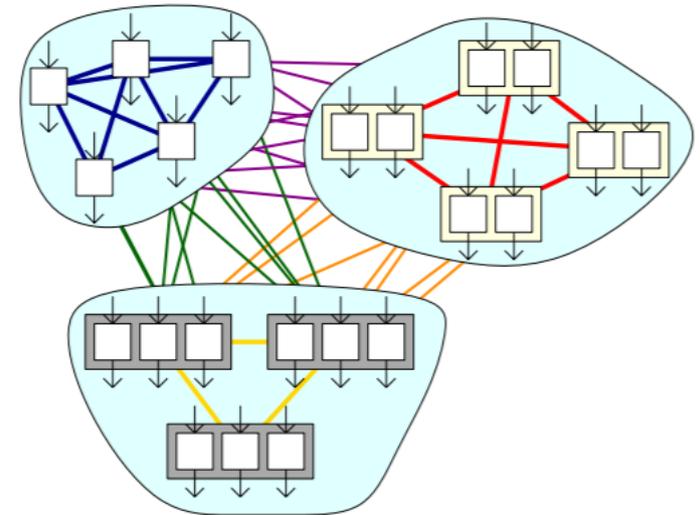
arXiv.org > quant-ph > arXiv:1802.09516

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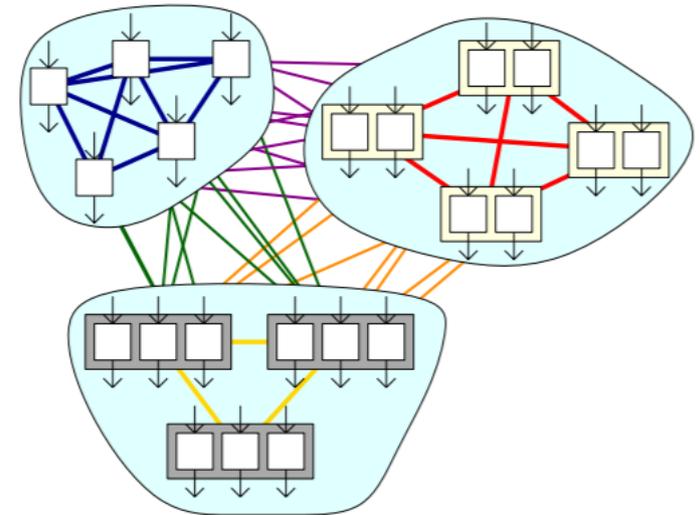
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- *Convergence analysis*



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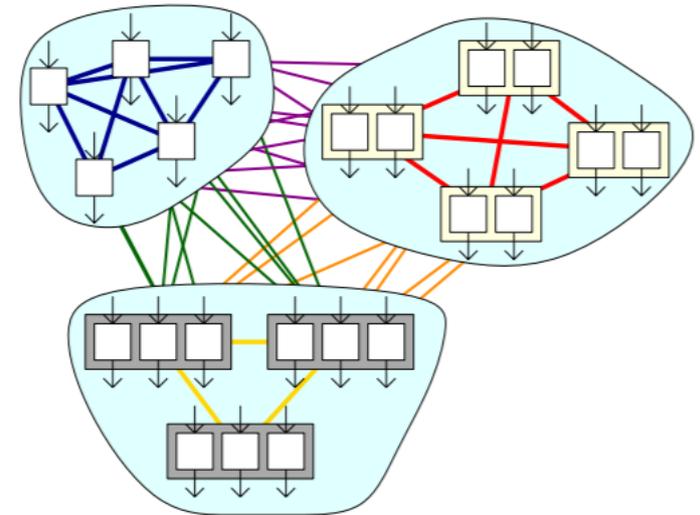
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(Submitted on 26 Feb 2018)

- *Convergence analysis*
- *Self-testing of spin-squeezed states*



Thanks for your attention!



Marie Skłodowska-Curie
Actions



European
Commission



MAX PLANCK INSTITUTE
OF QUANTUM OPTICS

Thanks for your attention!

