

# Constructing $k$ -uniform states of non-minimal support

Zahra Raissi, Christian Gogolin, Adam Teixido, Arnau Riera  
and Antonio Acín


ICFO - Institut de Ciències Fotoniques, Spain

Entanglement Days (Budapest), September 2018

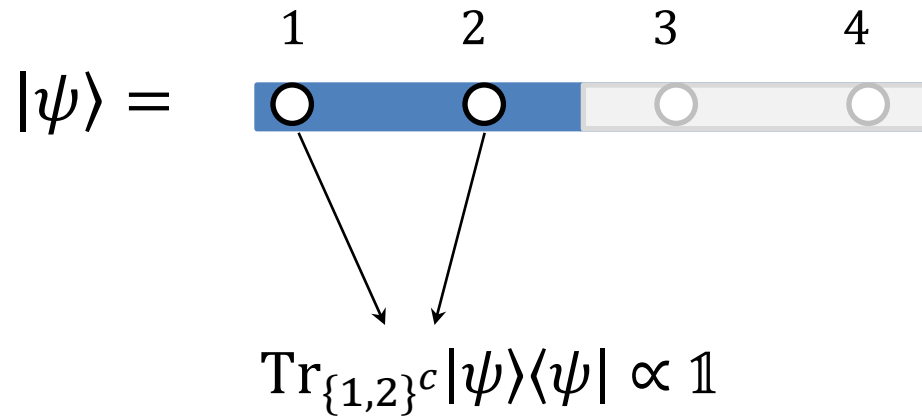
# $k$ -uniform states and Absolutely Maximally Entangled (AME) states

There is a fundamental question to ask, which states are useful for quantum information applications?

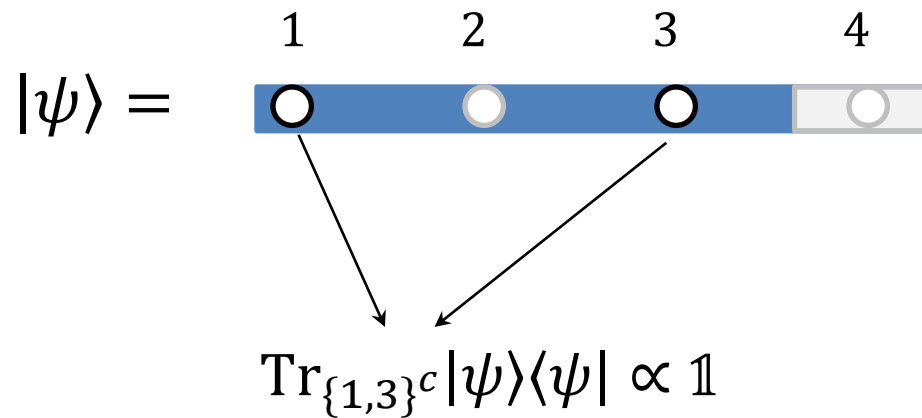
# What are AME states?

$$|\psi\rangle = \begin{array}{cccc} 1 & 2 & 3 & 4 \\ \bigcirc & \bigcirc & \bigcirc & \bigcirc \end{array}$$
A diagram representing a quantum state  $|\psi\rangle$  on four qubits. The qubits are labeled 1, 2, 3, and 4. Each qubit is represented by a white circle with a black outline. A solid blue horizontal bar highlights all four circles, indicating that the state has support on all four qubits.

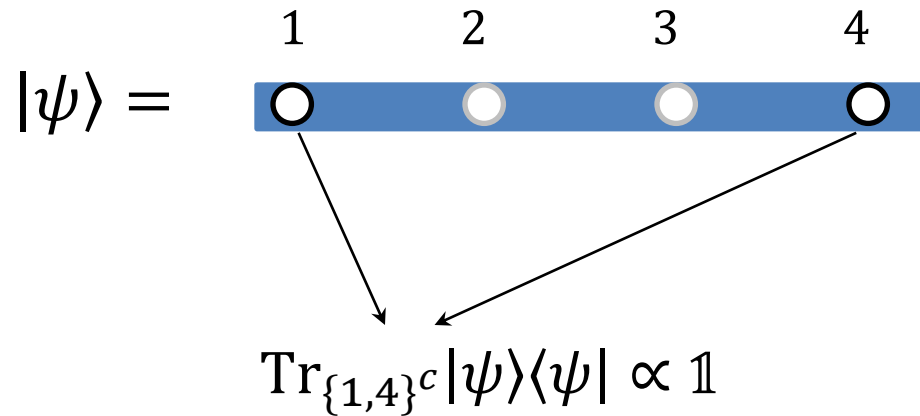
# What are AME states?




# What are AME states?



# What are AME states?



# What are AME states?

$|\psi\rangle =$    $\text{Tr}_{\{2,3\}^c} |\psi\rangle\langle\psi| \propto \mathbb{1}$

The diagram shows a horizontal bar representing a 4-qubit system. The bar is divided into four segments labeled 1, 2, 3, and 4. Each segment contains a small white circle representing a qubit. The segments for qubits 2 and 3 are highlighted in blue, indicating the support of the state. Two arrows point from the blue segments to the trace operation below.

# What are AME states?

$$|\psi\rangle = \begin{array}{cccc} 1 & 2 & 3 & 4 \\ \hline \text{○} & \text{○} & \text{○} & \text{○} \end{array}$$

$\text{Tr}_{\{3,4\}^c} |\psi\rangle\langle\psi| \propto \mathbb{1}$

*AME*( $n, q$ ) states:

A pure state of  $n$  parties with local dimension  $q$  is AME if reduced states on up to half of the systems are all maximally mixed, concretely

$$\text{AME}(n, q) :=$$

$$\{|\psi\rangle \in \mathcal{H}(n, q) : \forall S \subset \{1, \dots, n\} \mid |S| \leq \lfloor n/2 \rfloor \Rightarrow \text{Tr}_{S^c} |\psi\rangle\langle\psi| \propto \mathbb{1}\}$$



# Existence of AME states

- 
- [1] A. Higuchi, A. and Sudbery, Phys. Lett. A 273,213 (2000).
  - [2] A. J. Scott, Phys. Rev. A, 69, 052330 (2004).
  - [3] F. Huber, O. Gühne, and J. Siewert, Phys. Rev. Lett. 118, 200502 (2017).

# Existence of AME states

For qubits, ( $q = 2$ ):

---

[1] A. Higuchi, A. and Sudbery, Phys. Lett. A 273,213 (2000).

[2] A. J. Scott, Phys. Rev. A, 69, 052330 (2004).

[3] F. Huber, O. Gühne, and J. Siewert, Phys. Rev. Lett. 118, 200502 (2017).

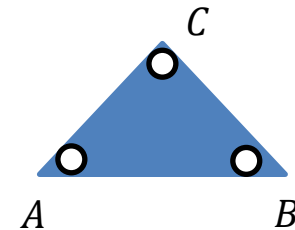
# Existence of AME states

For qubits, ( $q = 2$ ):  $n = 2, 3$

$$|\phi^+\rangle = |00\rangle + |11\rangle$$



$$|GHZ\rangle = |000\rangle + |111\rangle$$



---

[1] A. Higuchi, A. and Sudbery, Phys. Lett. A 273,213 (2000).

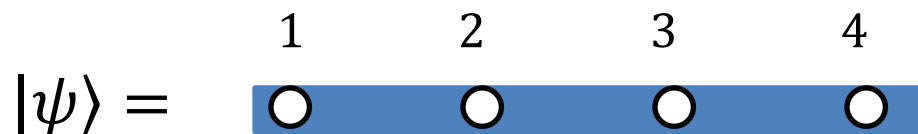
[2] A. J. Scott, Phys. Rev. A, 69, 052330 (2004).

[3] F. Huber, O. Gühne, and J. Siewert, Phys. Rev. Lett. 118, 200502 (2017).

# Existence of AME states

For qubits, ( $q = 2$ ):

$n = 2, 3, 4$



$$|\phi^\pm\rangle = |00\rangle \pm |11\rangle$$

$$|\psi^\pm\rangle = |01\rangle \pm |10\rangle$$

$$|\psi\rangle = |00\rangle_{12} |\phi^-\rangle_{34} + |11\rangle_{12} |\psi^-\rangle_{34} + |01\rangle_{12} |\psi^+\rangle_{34} + |10\rangle_{12} |\phi^+\rangle_{34} \quad [1]$$

[1] A. Higuchi, A. and Sudbery, Phys. Lett. A 273,213 (2000).

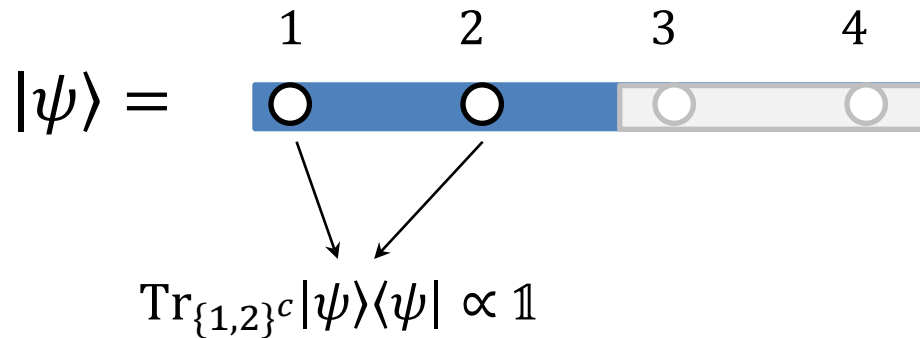
[2] A. J. Scott, Phys. Rev. A, 69, 052330 (2004).

[3] F. Huber, O. Gühne, and J. Siewert, Phys. Rev. Lett. 118, 200502 (2017).

# Existence of AME states

For qubits, ( $q = 2$ ):

$n = 2, 3, 4$



$$|\phi^\pm\rangle = |00\rangle \pm |11\rangle$$

$$|\psi^\pm\rangle = |01\rangle \pm |10\rangle$$

$$|\psi\rangle = |00\rangle_{12} |\phi^-\rangle_{34} + |11\rangle_{12} |\psi^-\rangle_{34} + |01\rangle_{12} |\psi^+\rangle_{34} + |10\rangle_{12} |\phi^+\rangle_{34} \quad [1]$$

[1] A. Higuchi, A. and Sudbery, Phys. Lett. A 273,213 (2000).

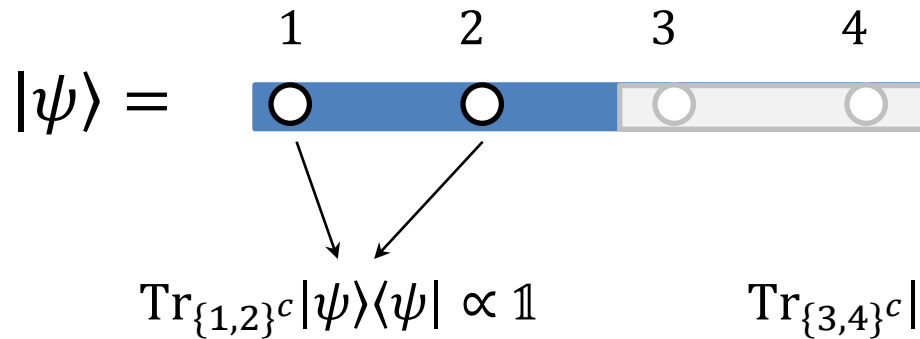
[2] A. J. Scott, Phys. Rev. A, 69, 052330 (2004).

[3] F. Huber, O. Gühne, and J. Siewert, Phys. Rev. Lett. 118, 200502 (2017).

# Existence of AME states

For qubits, ( $q = 2$ ):

$n = 2, 3, 4$



$$|\phi^\pm\rangle = |00\rangle \pm |11\rangle$$

$$|\psi^\pm\rangle = |01\rangle \pm |10\rangle$$

$$|\psi\rangle = |00\rangle_{12} |\phi^-\rangle_{34} + |11\rangle_{12} |\psi^-\rangle_{34} + |01\rangle_{12} |\psi^+\rangle_{34} + |10\rangle_{12} |\phi^+\rangle_{34} \quad [1]$$

[1] A. Higuchi, A. and Sudbery, Phys. Lett. A 273,213 (2000).

[2] A. J. Scott, Phys. Rev. A, 69, 052330 (2004).

[3] F. Huber, O. Gühne, and J. Siewert, Phys. Rev. Lett. 118, 200502 (2017).

# Existence of AME states

For qubits, ( $q = 2$ ):  $n = 2, 3, 4$

$$|\psi\rangle = \begin{array}{cccc} 1 & 2 & 3 & 4 \\ \bigcirc & \bigcirc & \bigcirc & \bigcirc \end{array} \quad \begin{array}{l} |\phi^\pm\rangle = |00\rangle \pm |11\rangle \\ |\psi^\pm\rangle = |01\rangle \pm |10\rangle \end{array}$$

$$\text{Tr}_{\{2,3\}^c} |\psi\rangle\langle\psi| \propto \mathbb{1} \quad \text{Tr}_{\{1,4\}^c} |\psi\rangle\langle\psi| \propto \mathbb{1}$$

$$|\psi\rangle = |00\rangle_{12} |\phi^-\rangle_{34} + |11\rangle_{12} |\psi^-\rangle_{34} + |01\rangle_{12} |\psi^+\rangle_{34} + |10\rangle_{12} |\phi^+\rangle_{34} \quad [1]$$

$$|\psi\rangle = |00\rangle_{23} |\phi^+\rangle_{14} - |01\rangle_{23} |\psi^-\rangle_{14} + |10\rangle_{23} |\psi^+\rangle_{14} - |11\rangle_{23} |\phi^-\rangle_{14}$$

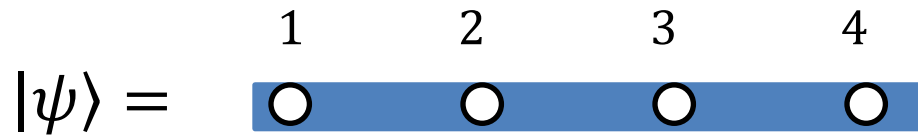
[1] A. Higuchi, A. and Sudbery, Phys. Lett. A 273,213 (2000).

[2] A. J. Scott, Phys. Rev. A, 69, 052330 (2004).

[3] F. Huber, O. Gühne, and J. Siewert, Phys. Rev. Lett. 118, 200502 (2017).

# Existence of AME states

For qubits, ( $q = 2$ ):  $n = 2, 3, 4$



$$|\phi^\pm\rangle = |00\rangle \pm |11\rangle$$

$$|\psi^\pm\rangle = |01\rangle \pm |10\rangle$$

$$\text{Tr}_{\{1,3\}^c} |\psi\rangle\langle\psi| \not\propto \mathbb{1}$$

$$\text{Tr}_{\{2,4\}^c} |\psi\rangle\langle\psi| \not\propto \mathbb{1}$$

$$|\psi\rangle = |00\rangle_{12} |\phi^-\rangle_{34} + |11\rangle_{12} |\psi^-\rangle_{34} + |01\rangle_{12} |\psi^+\rangle_{34} + |10\rangle_{12} |\phi^+\rangle_{34} \quad [1]$$

$$|\psi\rangle = |00\rangle_{23} |\phi^+\rangle_{14} - |01\rangle_{23} |\psi^-\rangle_{14} + |10\rangle_{23} |\psi^+\rangle_{14} - |11\rangle_{23} |\phi^-\rangle_{14}$$

$$|\psi\rangle = |\phi^+\rangle_{13} |\phi^+\rangle_{24} - |\psi^-\rangle_{13} |\psi^+\rangle_{24}$$

[1] A. Higuchi, A. and Sudbery, Phys. Lett. A 273,213 (2000).

[2] A. J. Scott, Phys. Rev. A, 69, 052330 (2004).

[3] F. Huber, O. Gühne, and J. Siewert, Phys. Rev. Lett. 118, 200502 (2017).



# Existence of AME states

For qubits, ( $q = 2$ ):  $n = 2, 3, 4$  [2]

$$|\psi\rangle = \begin{array}{cccc} 1 & 2 & 3 & 4 \\ \bigcirc & \bigcirc & \bigcirc & \bigcirc \end{array} \quad \begin{array}{l} |\phi^\pm\rangle = |00\rangle \pm |11\rangle \\ |\psi^\pm\rangle = |01\rangle \pm |10\rangle \end{array}$$

$$\text{Tr}_{\{1,3\}^c} |\psi\rangle\langle\psi| \not\propto \mathbb{1} \quad \text{Tr}_{\{2,4\}^c} |\psi\rangle\langle\psi| \not\propto \mathbb{1}$$

$$|\psi\rangle = |00\rangle_{12} |\phi^-\rangle_{34} + |11\rangle_{12} |\psi^-\rangle_{34} + |01\rangle_{12} |\psi^+\rangle_{34} + |10\rangle_{12} |\phi^+\rangle_{34} \quad [1]$$

$$|\psi\rangle = |00\rangle_{23} |\phi^+\rangle_{14} - |01\rangle_{23} |\psi^-\rangle_{14} + |10\rangle_{23} |\psi^+\rangle_{14} - |11\rangle_{23} |\phi^-\rangle_{14}$$

$$|\psi\rangle = |\phi^+\rangle_{13} |\phi^+\rangle_{24} - |\psi^-\rangle_{13} |\psi^+\rangle_{24}$$

[1] A. Higuchi, A. and Sudbery, Phys. Lett. A 273,213 (2000).

[2] A. J. Scott, Phys. Rev. A, 69, 052330 (2004).

[3] F. Huber, O. Gühne, and J. Siewert, Phys. Rev. Lett. 118, 200502 (2017).

# Existence of AME states

For qubits, ( $q = 2$ ):  $n = 2, 3, 4, 5, 6, 7, 8, 9, \dots$  [2]

$$|\psi\rangle = \begin{array}{cccc} 1 & 2 & 3 & 4 \\ \bigcirc & \bigcirc & \bigcirc & \bigcirc \end{array} \quad \begin{array}{l} |\phi^\pm\rangle = |00\rangle \pm |11\rangle \\ |\psi^\pm\rangle = |01\rangle \pm |10\rangle \end{array}$$

$$\text{Tr}_{\{1,3\}^c} |\psi\rangle\langle\psi| \not\propto \mathbb{1} \quad \text{Tr}_{\{2,4\}^c} |\psi\rangle\langle\psi| \not\propto \mathbb{1}$$

$$|\psi\rangle = |00\rangle_{12} |\phi^-\rangle_{34} + |11\rangle_{12} |\psi^-\rangle_{34} + |01\rangle_{12} |\psi^+\rangle_{34} + |10\rangle_{12} |\phi^+\rangle_{34} \quad [1]$$

$$|\psi\rangle = |00\rangle_{23} |\phi^+\rangle_{14} - |01\rangle_{23} |\psi^-\rangle_{14} + |10\rangle_{23} |\psi^+\rangle_{14} - |11\rangle_{23} |\phi^-\rangle_{14}$$

$$|\psi\rangle = |\phi^+\rangle_{13} |\phi^+\rangle_{24} - |\psi^-\rangle_{13} |\psi^+\rangle_{24}$$

[1] A. Higuchi, A. and Sudbery, Phys. Lett. A 273,213 (2000).

[2] A. J. Scott, Phys. Rev. A, 69, 052330 (2004).

[3] F. Huber, O. Gühne, and J. Siewert, Phys. Rev. Lett. 118, 200502 (2017).

# Existence of AME states

For qubits, ( $q = 2$ ):  $n = 2, 3, 4, 5, 6, 7, 8, 9, \dots$  [2, 3]

$$|\psi\rangle = \begin{array}{cccc} 1 & 2 & 3 & 4 \\ \bigcirc & \bigcirc & \bigcirc & \bigcirc \end{array} \quad \begin{array}{l} |\phi^\pm\rangle = |00\rangle \pm |11\rangle \\ |\psi^\pm\rangle = |01\rangle \pm |10\rangle \end{array}$$

$$\text{Tr}_{\{1,3\}^c} |\psi\rangle\langle\psi| \not\propto \mathbb{1} \quad \text{Tr}_{\{2,4\}^c} |\psi\rangle\langle\psi| \not\propto \mathbb{1}$$

$$|\psi\rangle = |00\rangle_{12} |\phi^-\rangle_{34} + |11\rangle_{12} |\psi^-\rangle_{34} + |01\rangle_{12} |\psi^+\rangle_{34} + |10\rangle_{12} |\phi^+\rangle_{34} \quad [1]$$

$$|\psi\rangle = |00\rangle_{23} |\phi^+\rangle_{14} - |01\rangle_{23} |\psi^-\rangle_{14} + |10\rangle_{23} |\psi^+\rangle_{14} - |11\rangle_{23} |\phi^-\rangle_{14}$$

$$|\psi\rangle = |\phi^+\rangle_{13} |\phi^+\rangle_{24} - |\psi^-\rangle_{13} |\psi^+\rangle_{24}$$

[1] A. Higuchi, A. and Sudbery, Phys. Lett. A 273,213 (2000).

[2] A. J. Scott, Phys. Rev. A, 69, 052330 (2004).

[3] F. Huber, O. Gühne, and J. Siewert, Phys. Rev. Lett. 118, 200502 (2017).

# $k$ -uniform states

- Since AME states may not always exist, one can loosen the criteria for maximal mixedness,

$$|\psi\rangle = \begin{array}{cccccc} & 1 & 2 & 3 & 4 & \dots & n \\ & \bigcirc & \bigcirc & \bigcirc & \bigcirc & \dots & \bigcirc \end{array}$$

$AME(n, q)$  states:

A pure state  $|\psi\rangle$  of  $n$  parties with local dimension  $q$  is AME if for all  $S \subset \{1, 2, \dots, n\}$ ,

$$|S| \leq \lfloor n/2 \rfloor \implies \text{Tr}_{S^c} |\psi\rangle\langle\psi| \propto \mathbb{1}$$

$k$ -UNI( $n, q$ ) states:

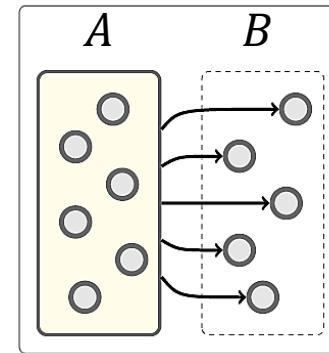
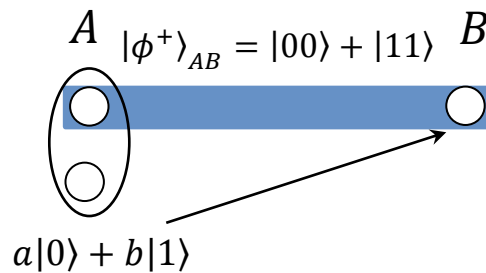
A pure state  $|\psi\rangle$  of  $n$  parties with local dimension  $q$  is  $k$ -uniform if for all  $S \subset \{1, 2, \dots, n\}$ ,

$$|S| \leq k \implies \text{Tr}_{S^c} |\psi\rangle\langle\psi| \propto \mathbb{1}$$

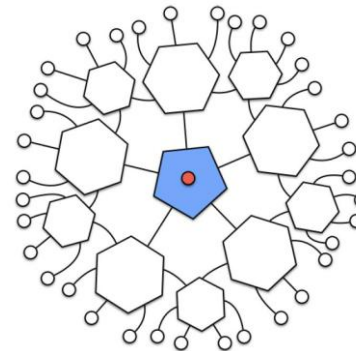
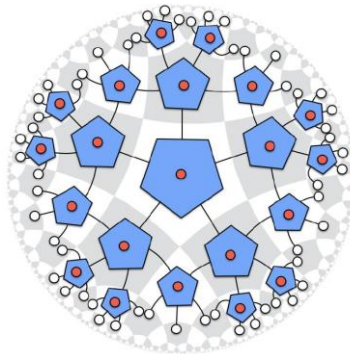
- Obviously, an AME state is a  $k = \lfloor n/2 \rfloor$ -uniform state.

# Why are $k$ -uniform states interesting?

- Natural generalization of EPR and GHZ states
- Resource for multipartite parallel teleportation [1]



- Holographic models implementing the AdS/CFT correspondence [2]



[1] W. Helwig, W. Cui, J. I. Latorre, A. Riera, and H.K. Lo, Phys. Rev. A, 86, 052335 (2012).

[2] F. Patawski, B. Yoshida, D. Harlow, and J. Preskill, Journal of High Energy Physics 06, 149 (2015).

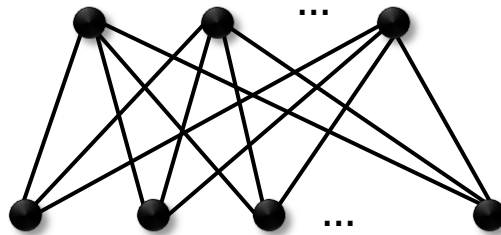
# Content of this talk

Classical error  
correcting codes



$k$ -uniform states of  
minimal support

Graph states:



orthonormal basis ( $k$ -uniform basis):

$$|\psi_{\vec{a}}\rangle := M(\vec{v}) |\psi\rangle$$

<i>Classical Part</i>	<i>Quantum Part</i>
All terms	orthonormal basis
$\ell - \text{UNI}_{\min}(n_{\text{cl}}, q)$	$\ell' - \text{UNI}_{\min}(n_{\text{q}}, q)$

$k$ -uniform state of non-minimal support

## $k$ -uniform states of minimal support

There is a connection between  $k$ -uniform states of minimal support and classical maximum distance separable (MDS) error correcting codes.

Using that we can provide explicit closed form expressions for this set of  $k$ -uniform states.

Classical error  
correcting codes   $k$ -uniform states

# $k$ -uniform states of minimal support from MDS codes

- We classify the  $k$ -uniform states according to the **number of their terms**, when they are expanded in product basis.

$$|\psi\rangle = \sum_{j_1, \dots, j_n=0}^{q-1} c_{j_1, \dots, j_n} |j_1, \dots, j_n\rangle \longrightarrow \# \text{ terms} = q^k$$



# $k$ -uniform states of minimal support from MDS codes

- We classify the  $k$ -uniform states according to the **number of their terms**, when they are expanded in product basis.

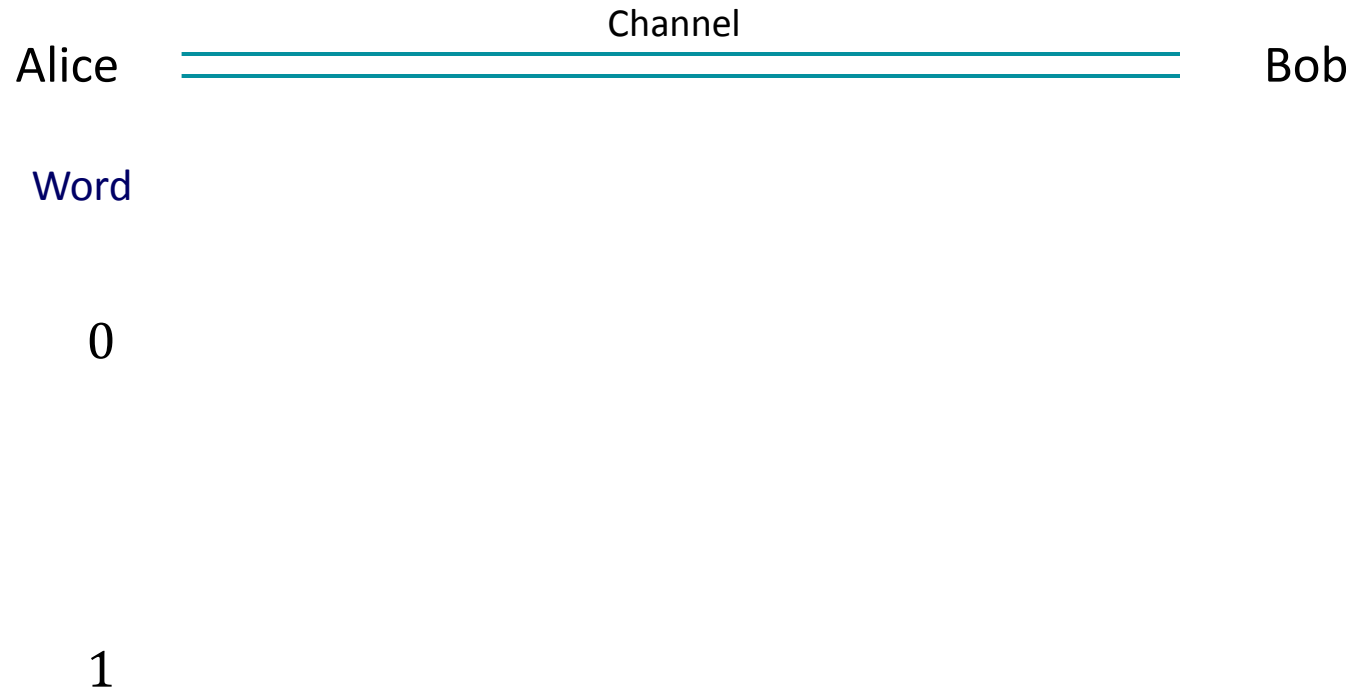
$$|\psi\rangle = \sum_{j_1, \dots, j_n=0}^{q-1} c_{j_1, \dots, j_n} |j_1, \dots, j_n\rangle \quad \longrightarrow \quad \# \text{ terms} = q^k$$

- Construct  $k$ -uniform states with **minimal support** from classical **MDS codes**

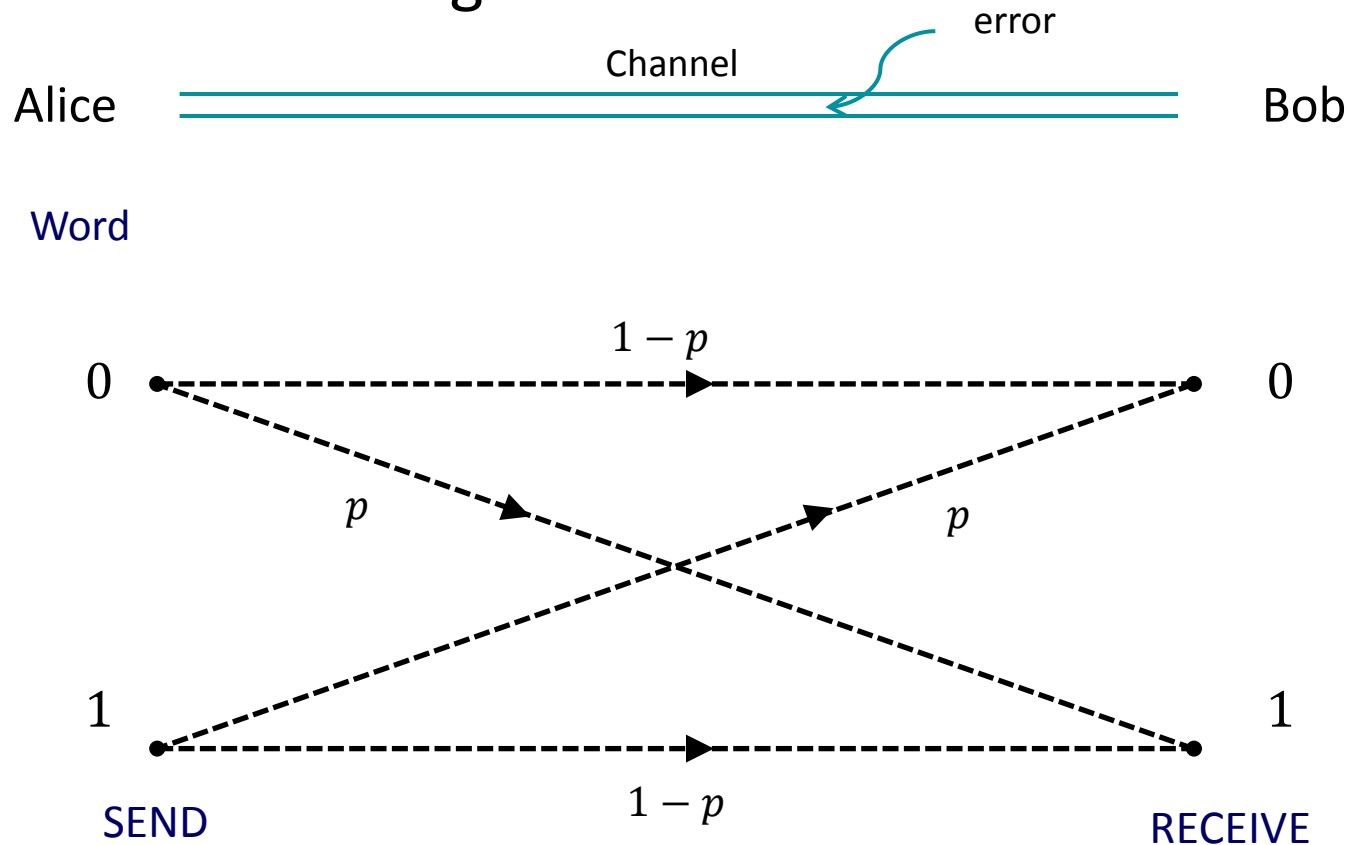
$$\begin{cases} n \leq q + 2 & \text{when } q \text{ is even and } k = 1 \text{ or } k = q - 1 \\ n \leq q + 1 & \text{in all other cases} \end{cases}$$

- Construct **basis**, develop **stabilizer** formalism and **AME graph states**

# Classical error correcting codes



# Classical error correcting codes



# Classical error correcting codes



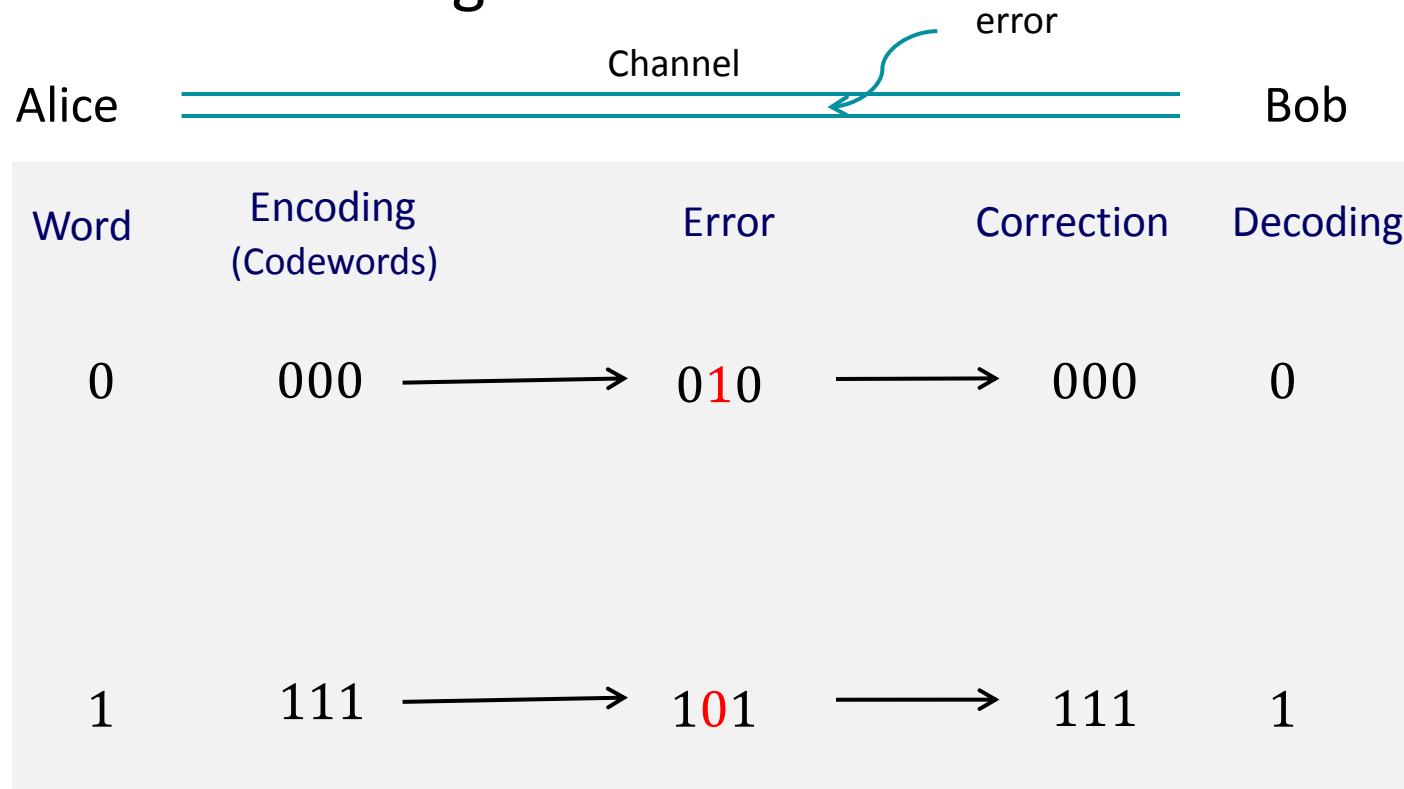
Word	Encoding (Codewords)
0	000
1	111

# Classical error correcting codes

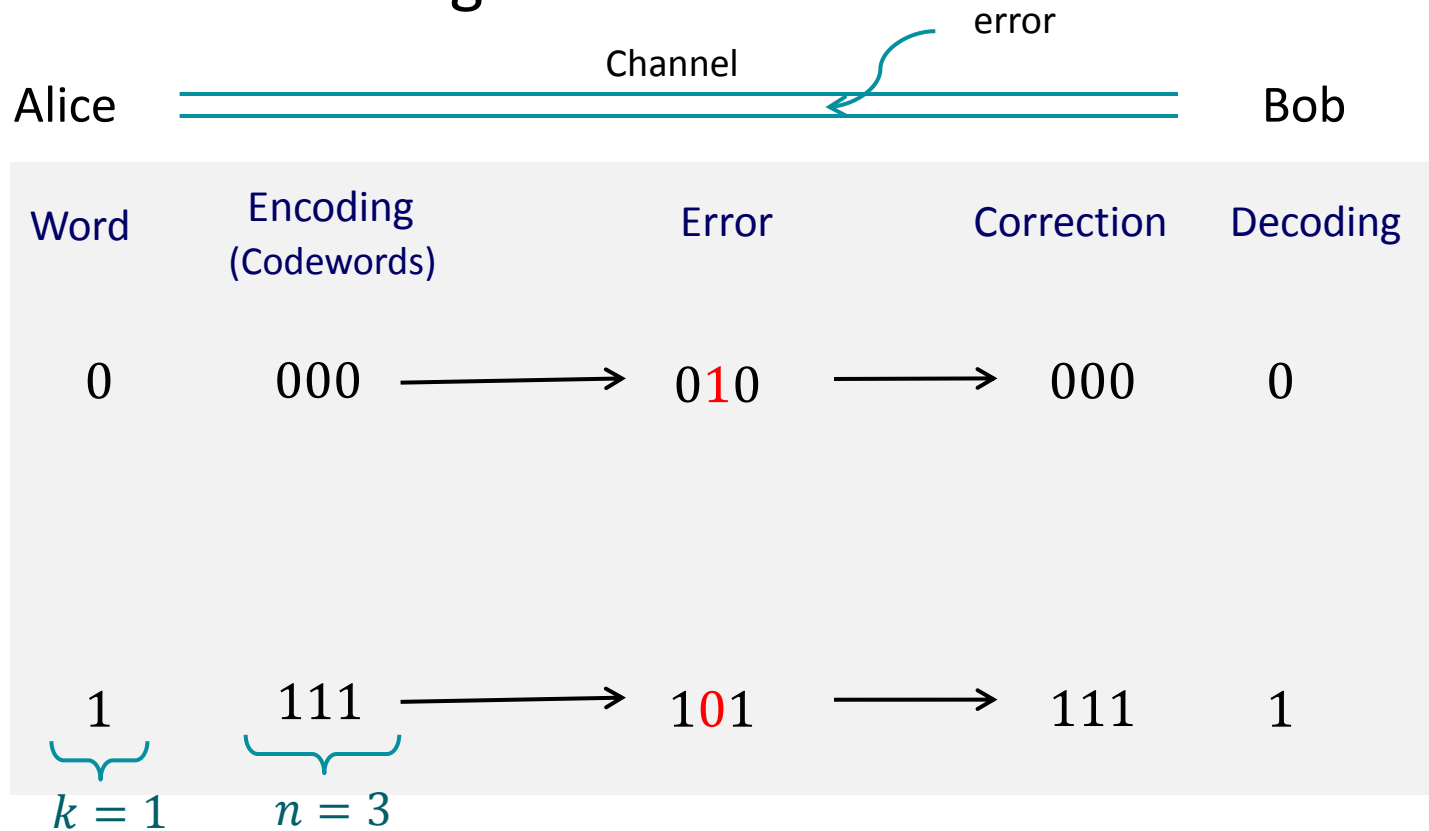


Word	Encoding (Codewords)	Error
0	000	→ 010
1	111	→ 101

# Classical error correcting codes

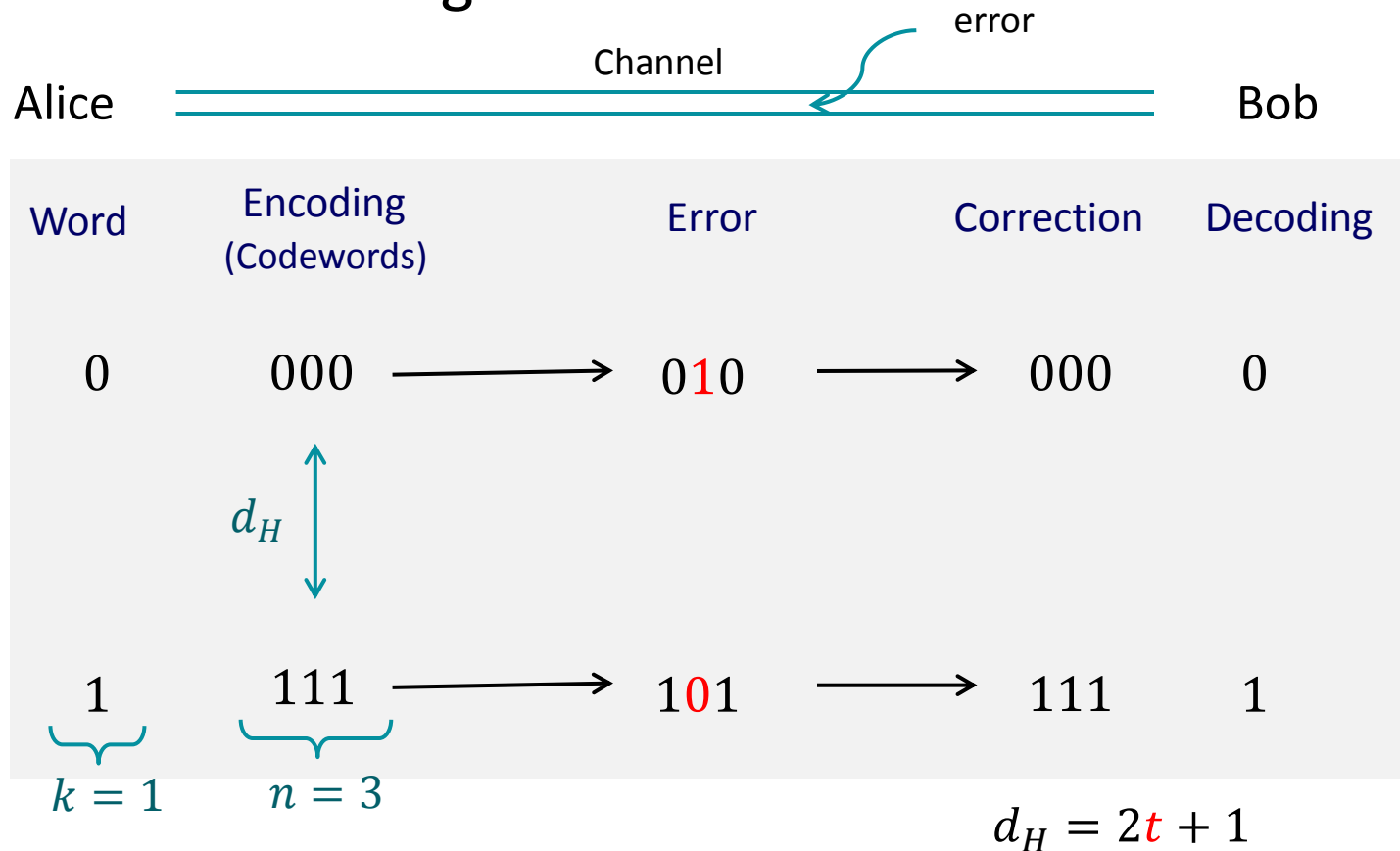


# Classical error correcting codes



$$\vec{m} = (x_1, \dots, x_k) \longrightarrow \vec{c} = (\underbrace{x_1, \dots, x_k}_{\text{Message symbols}}, \underbrace{x_{k+1}, \dots, x_n}_{\text{Check symbols}})$$

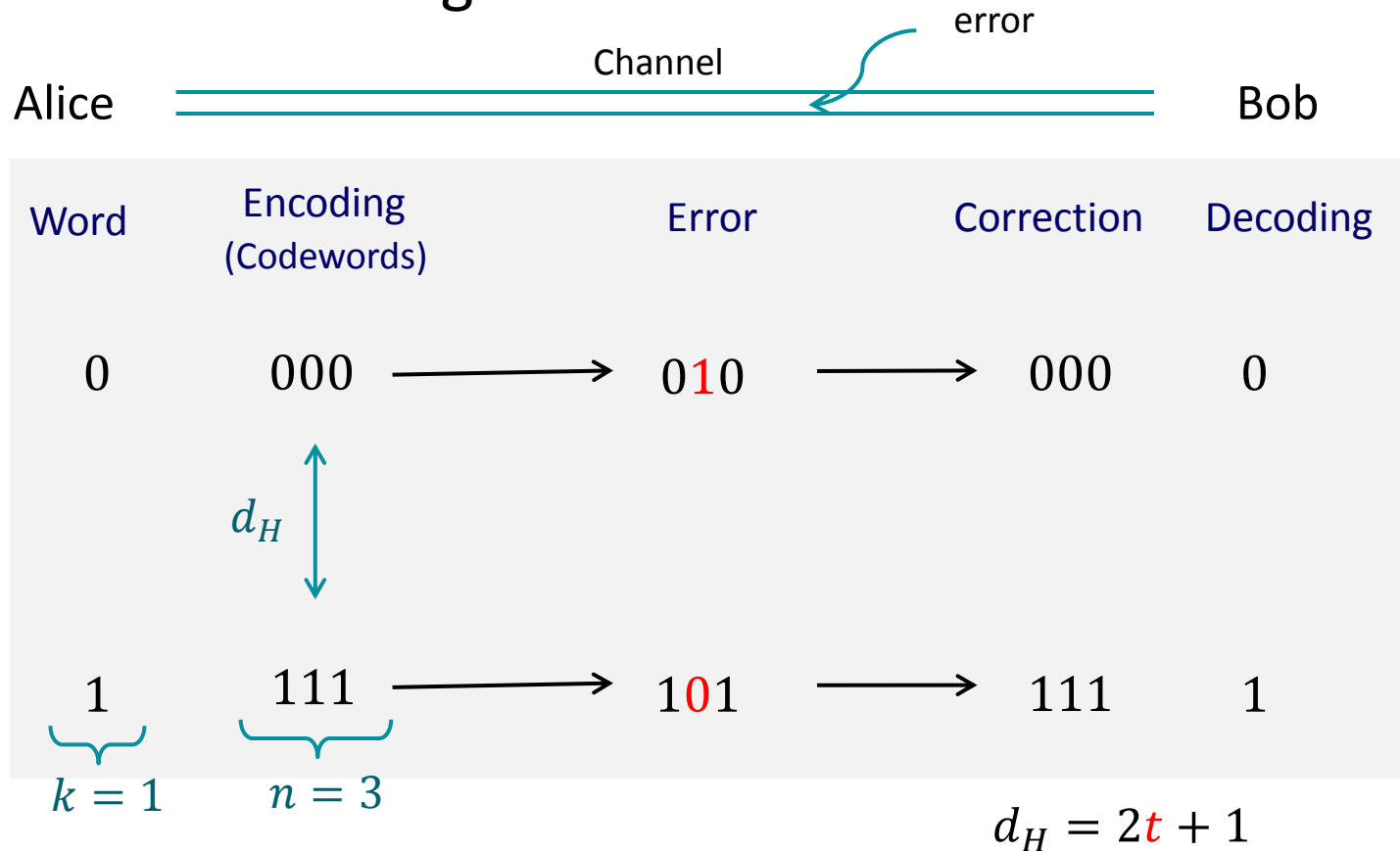
# Classical error correcting codes



$$\vec{m} = (x_1, \dots, x_k) \longrightarrow \vec{c} = (\underbrace{x_1, \dots, x_k}_{\text{Message symbols}}, \underbrace{x_{k+1}, \dots, x_n}_{\text{Check symbols}})$$



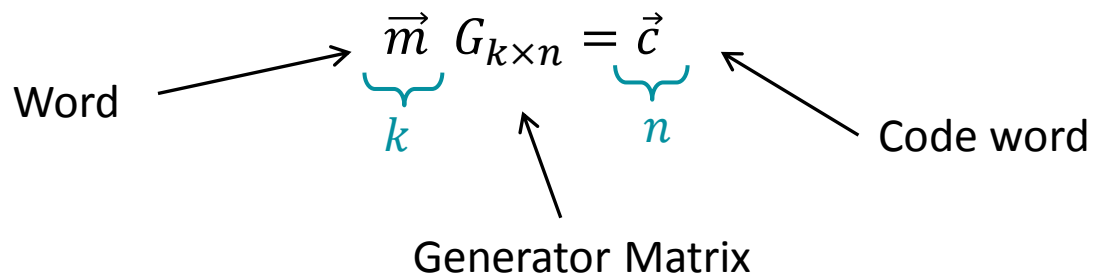
# Classical error correcting codes



$$[n = 3, k = 1, d_H = 3]_{q=2}$$

# MDS codes

- Constructing classical error correcting codes  $[n, k, d_H]_q$  [1]

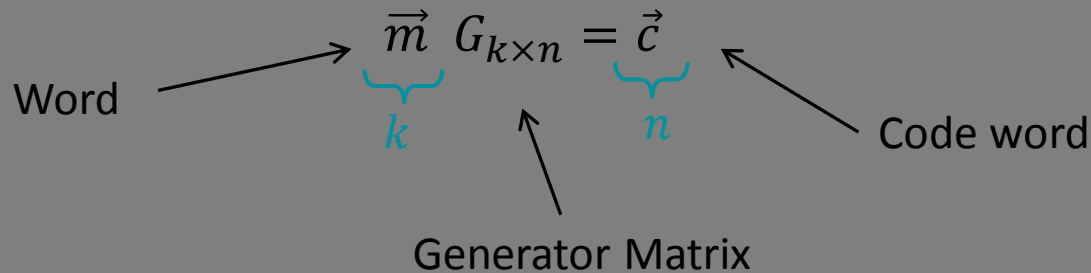


[1] F.J. MacWilliams, N.J.A. Sloane, The theory of error-correction codes (1977).

[2] R. Singleton, IEEE Trans. Inf. Theor., 10, 116 (2006).

# MDS codes

- Constructing classical error correcting codes  $[n, k, d_H]_q$



**This only makes sense if you can take linear combinations of the code words.**

What do we need? **Finite fields**

**Prime numbers**,  $q = p$ : Integers modulo  $q$  form a field

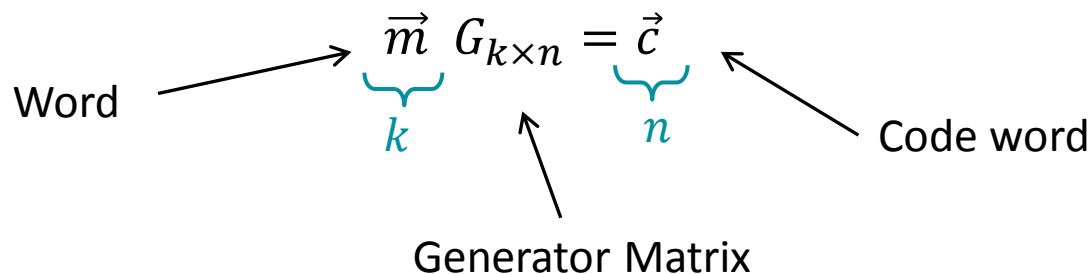
$$\text{e.g. } \text{GF}(5) = \{0, 1, 2, 3, 4\} \pmod{5}$$

**Prime powers**,  $q = p^m$ : It works with the representation in terms of polynomials based on the **irreducible polynomial**

$$\begin{aligned} \text{e.g. } \text{GF}(2^2) &= \{0, 1, x, x + 1\} \quad \text{generated by} \quad x^2 = x + 1 \\ &\equiv \{0, 1, 2, 3\} \end{aligned}$$

# MDS codes

- Constructing classical error correcting codes  $[n, k, d_H]_q$  [1]



- $G_{k \times n}$  has standard form (by taking linear combination of the codewords)

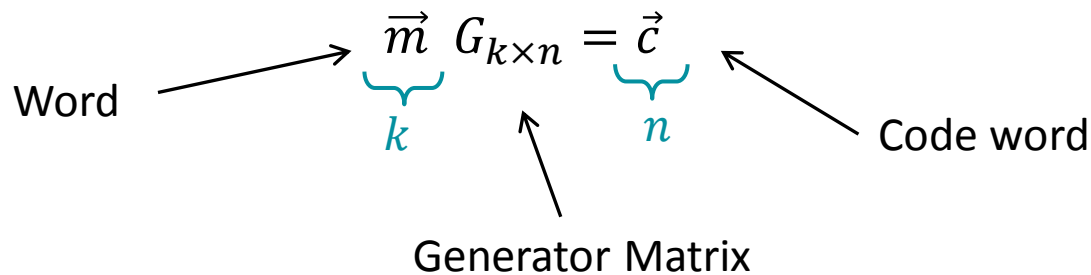
$$G_{k \times n} = [\mathbb{1}_k | A]$$

[1] F.J. MacWilliams, N.J.A. Sloane, The theory of error-correction codes (1977).

[2] R. Singleton, IEEE Trans. Inf. Theor., 10, 116 (2006).

# MDS codes

- Constructing classical error correcting codes  $[n, k, d_H]_q$  [1]



- $G_{k \times n}$  has standard form (by taking linear combination of the codewords)

$$G_{k \times n} = [\mathbb{1}_k | A]$$

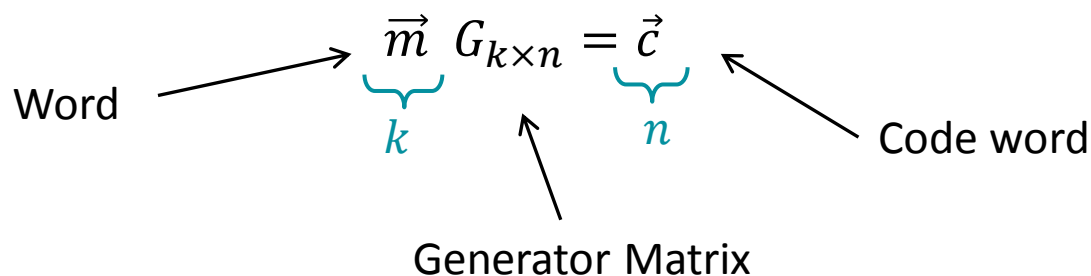
- Singleton bound for any linear code:  $d_H \leq n - k + 1$  [2]

[1] F.J. MacWilliams, N.J.A. Sloane, The theory of error-correction codes (1977).

[2] R. Singleton, IEEE Trans. Inf. Theor., 10, 116 (2006).

# MDS codes

- Constructing classical error correcting codes  $[n, k, d_H]_q$  [1]



- $G_{k \times n}$  has standard form (by taking linear combination of the codewords)

$$G_{k \times n} = [\mathbb{1}_k | A]$$

- Singleton bound for any linear code:  $d_H \leq n - k + 1$  [2]
- Code is Maximum Distance Separable (MDS) if  $d_H = n - k + 1 \rightarrow$  Code is MDS iff any subset of  $k$  columns of  $G_{k \times n}$  is linearly independent.

[1] F.J. MacWilliams, N.J.A. Sloane, The theory of error-correction codes (1977).

[2] R. Singleton, IEEE Trans. Inf. Theor., 10, 116 (2006).

# An example of $k$ -uniform state

- From any MDS code a  $k$ -uniform state can be constructed by taking the **equally weighted superposition** of all the codewords

Classical MDS codes   $k$ -uniform states

# An example of $k$ -uniform state

- From any MDS code a  $k$ -uniform state can be constructed by taking the **equally weighted superposition** of all the codewords

Classical MDS codes  $\longrightarrow$   $k$ -uniform states

- Generator matrix of an MDS code  $[6,2,5]_5$

$$G_{2 \times 6} = \left[ \begin{array}{cc|cccc} 1 & 0 & 1 & 1 & 1 & 1 \\ 0 & 1 & 1 & 2 & 3 & 4 \end{array} \right]$$

- Yields minimal support 2-uniform state for  $n = 6$ ,  $q = 5$ :

$$|\psi\rangle = \sum_{\vec{m} \in GF(5)^2} |\vec{m} G_{2 \times 6}\rangle = \sum_{i,j=0}^4 |i, j, i+j, i+2j, i+3j, i+4j\rangle$$

(All additions and multiplications modulo  $q$ .)



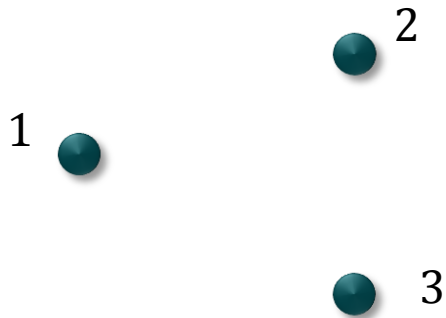
## Graph state and basis

Description of the  $k$ -uniform states of minimal support within the [graph state](#) formalism.

Starting from a single  $k$ -uniform state of minimal support  $|\psi\rangle \in \mathcal{H}(n, q)$ , a [complete orthonormal basis](#) for  $\mathcal{H}(n, q)$  can be constructed.

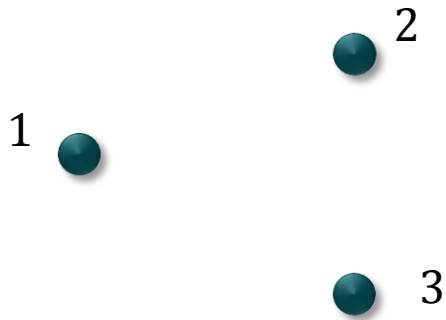
# Graph states: Introduction

- initialize each qubit as the state,  $|+\rangle = \frac{|0\rangle + \dots + |q-1\rangle}{\sqrt{q}} \equiv \bullet$  Qudits



# Graph states: Introduction

- initialize each qubit as the state,  $|+\rangle = \frac{|0\rangle + \dots + |q-1\rangle}{\sqrt{q}} \equiv \bullet$  Qudits



- Perform  $CZ$  operations between any two qubits that are connected by an edge.

$$CZ_{ij} = \sum_l |l\rangle\langle l|_i \otimes Z^l_j = \sum_{l,m} \omega^{lm} |l\rangle\langle l|_i \otimes |m\rangle\langle m|_j$$



Adjacency matrix  $\Gamma$

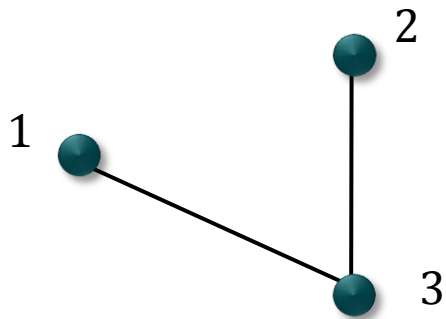
$\Gamma_{ij} > 0$   $i, j$  connected

$\Gamma_{ij} = 0$  not connected

- Multiple edges are possible because we are considering qudits

# Graph states: Introduction

- initialize each qubit as the state,  $|+\rangle = \frac{|0\rangle + \dots + |q-1\rangle}{\sqrt{q}} \equiv \bullet$  Qudits



- Perform  $CZ$  operations between any two qubits that are connected by an edge.

$$CZ_{ij} = \sum_l |l\rangle\langle l|_i \otimes Z^l_j = \sum_{l,m} \omega^{lm} |l\rangle\langle l|_i \otimes |m\rangle\langle m|_j$$



Adjacency matrix  $\Gamma$

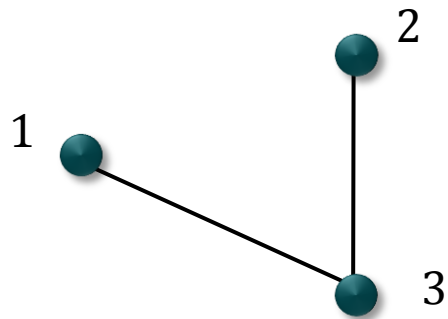
$\Gamma_{ij} > 0$   $i, j$  connected

$\Gamma_{ij} = 0$  not connected

- Multiple edges are possible because we are considering qudits

# Graph states: Introduction

- initialize each qubit as the state,  $|+\rangle = \frac{|0\rangle + \dots + |q-1\rangle}{\sqrt{q}} \equiv \bullet$  Qudits



$$|\Gamma\rangle = \prod_{i>j} CZ_{ij}^{\Gamma_{ij}} |+\rangle^{\otimes n}$$

Stabilizers:

$$g_l^\Gamma = X_l \prod_m (Z_m)^{\Gamma_{lm}} \quad 1 \leq l \leq n$$

- Perform  $CZ$  operations between any two qubits that are connected by an edge.

$$CZ_{ij} = \sum_l |l\rangle\langle l|_i \otimes Z^l_j = \sum_{l,m} \omega^{lm} |l\rangle\langle l|_i \otimes |m\rangle\langle m|_j$$



Adjacency matrix  $\Gamma$

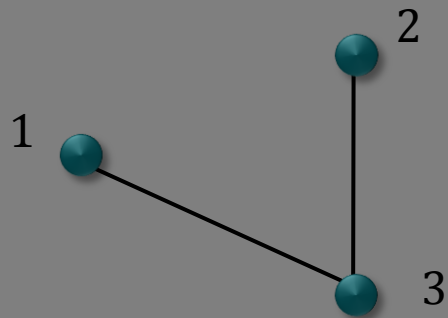
$\Gamma_{ij} > 0$   $i, j$  connected

$\Gamma_{ij} = 0$  not connected

- Multiple edges are possible because we are considering qudits

# Graph states: Introduction

- initialize each qubit as the state,  $|+\rangle = \frac{|0\rangle + \dots + |q-1\rangle}{\sqrt{q}} \equiv \bullet$  Qudits



$$|\Gamma\rangle = \prod_{i>j} CZ_{ij}^{\Gamma_{ij}} |+\rangle^{\otimes n}$$

Stabilizers:

$$g_l^\Gamma = X_l \prod_m (Z_m)^{\Gamma_{lm}} \quad 1 \leq l \leq n$$

$X$  and  $Z$  that generalize the Pauli operators  $\sigma_X$  and  $\sigma_Z$  to Hilbert spaces of dimension  $q \leq 2$

$$\begin{cases} X|j\rangle = |j+1 \bmod q\rangle \\ Z|j\rangle = \omega^j |j\rangle \end{cases} \quad \omega := e^{\frac{2\pi i}{q}}$$

Stabilizers of the example:

$$\begin{cases} g_1 = X \otimes \mathbb{1} \otimes Z \\ g_2 = \mathbb{1} \otimes X \otimes Z \\ g_3 = Z \otimes Z \otimes X \end{cases}$$

# Stabilizers

- The minimal support  $k$ -uniform state are linear combinations of the rows of  $G_{k \times n}$

$$k - \text{UNI}_{\min}(n, q) \ni |\psi\rangle = \sum_{\vec{v}} |\vec{v} G_{k \times n}\rangle \quad \longrightarrow \quad G_{k \times n} = [\mathbb{1}_k | A_{k \times n-k}]$$

$$\left\{ \begin{array}{l} s_1 = X \quad \mathbb{1} \quad \dots \quad \mathbb{1} \\ s_2 = \mathbb{1} \quad X \quad \dots \quad \mathbb{1} \\ \vdots \\ s_k = \mathbb{1} \quad \mathbb{1} \quad \dots \quad X \end{array} \right. \begin{array}{l} \xleftarrow{k} \quad \xrightarrow{n-k} \\ X^{a_{11}} \quad X^{a_{12}} \quad \dots \quad X^{a_{1(n-k)}} \\ X^{a_{21}} \quad X^{a_{22}} \quad \dots \quad X^{a_{2(n-k)}} \\ \dots \\ X^{a_{k1}} \quad X^{a_{k2}} \quad \dots \quad X^{a_{k(n-k)}} \end{array}$$

# Stabilizers

- The minimal support  $k$ -uniform state are linear combinations of the rows of  $G_{k \times n}$

$$k - \text{UNI}_{\min}(n, q) \ni |\psi\rangle = \sum_{\vec{v}} |\vec{v} G_{k \times n}\rangle \longrightarrow G_{k \times n} = [\mathbb{1}_k | A_{k \times n-k}]$$

$$\left\{ \begin{array}{l} s_1 = X \quad \mathbb{1} \quad \dots \quad \mathbb{1} \quad X^{a_{11}} \quad X^{a_{12}} \quad \dots \quad X^{a_{1(n-k)}} \\ s_2 = \mathbb{1} \quad X \quad \dots \quad \mathbb{1} \quad X^{a_{21}} \quad X^{a_{22}} \quad \dots \quad X^{a_{2(n-k)}} \\ \vdots \\ s_k = \mathbb{1} \quad \mathbb{1} \quad \dots \quad X \quad X^{a_{k1}} \quad X^{a_{k2}} \quad \dots \quad X^{a_{k(n-k)}} \end{array} \right. \quad H_{n-k \times n} = [-A^T | \mathbb{1}_{n-k}]$$

$$G_{k \times n} (H_{n-k \times n})^T = 0 \longrightarrow s_l \psi |\psi\rangle = \sum_{\vec{v}} \omega^{H_{n-k \times n} (G_{k \times n})^T \vec{v}} |\vec{v} G_{k \times n}\rangle = |\psi\rangle$$

$$\left\{ \begin{array}{l} s_{k+1} = Z^{-a_{11}} \quad Z^{-a_{21}} \quad \dots \quad Z^{-a_{k1}} \quad Z \quad \mathbb{1} \quad \dots \quad \mathbb{1} \\ s_{k+2} = Z^{-a_{12}} \quad Z^{-a_{22}} \quad \dots \quad Z^{-a_{k2}} \quad \mathbb{1} \quad Z \quad \dots \quad \mathbb{1} \\ \vdots \\ s_n = Z^{-a_{1(n-k)}} \quad Z^{-a_{2(n-k)}} \quad \dots \quad Z^{-a_{k(n-k)}} \quad \mathbb{1} \quad \mathbb{1} \quad \dots \quad Z \end{array} \right.$$



# Stabilizers

- The minimal support  $k$ -uniform state are linear combinations of the rows of  $G_{k \times n}$

$$k - \text{UNI}_{\min}(n, q) \ni |\psi\rangle = \sum_{\vec{v}} |\vec{v} G_{k \times n}\rangle \longrightarrow G_{k \times n} = [\mathbb{1}_k | A_{k \times n-k}]$$

$$\left\{ \begin{array}{l} s_1 = X \quad \mathbb{1} \quad \dots \quad \mathbb{1} \quad X^{a_{11}} \quad X^{a_{12}} \quad \dots \quad X^{a_{1(n-k)}} \\ s_2 = \mathbb{1} \quad X \quad \dots \quad \mathbb{1} \quad X^{a_{21}} \quad X^{a_{22}} \quad \dots \quad X^{a_{2(n-k)}} \\ \vdots \\ s_k = \mathbb{1} \quad \mathbb{1} \quad \dots \quad X \quad X^{a_{k1}} \quad X^{a_{k2}} \quad \dots \quad X^{a_{k(n-k)}} \end{array} \right. \quad H_{n-k \times n} = [-A^T | \mathbb{1}_{n-k}]$$

$$G_{k \times n} (H_{n-k \times n})^T = 0 \longrightarrow s_l \psi |\psi\rangle = \sum_{\vec{v}} \omega^{H_{n-k \times n} (G_{k \times n})^T \vec{v}} |\vec{v} G_{k \times n}\rangle = |\psi\rangle$$

$$\left\{ \begin{array}{l} s_{k+1} = Z^{-a_{11}} \quad Z^{-a_{21}} \quad \dots \quad Z^{-a_{k1}} \quad Z \quad \mathbb{1} \quad \dots \quad \mathbb{1} \\ s_{k+2} = Z^{-a_{12}} \quad Z^{-a_{22}} \quad \dots \quad Z^{-a_{k2}} \quad \mathbb{1} \quad Z \quad \dots \quad \mathbb{1} \\ \vdots \\ s_n = Z^{-a_{1(n-k)}} \quad Z^{-a_{2(n-k)}} \quad \dots \quad Z^{-a_{k(n-k)}} \quad \mathbb{1} \quad \mathbb{1} \quad \dots \quad Z \end{array} \right. \quad \begin{array}{l} \text{on the last } n - k \text{ qudits} \\ F^{-1} Z F = X \\ F^{-1} X F = Z^{-1} \end{array}$$

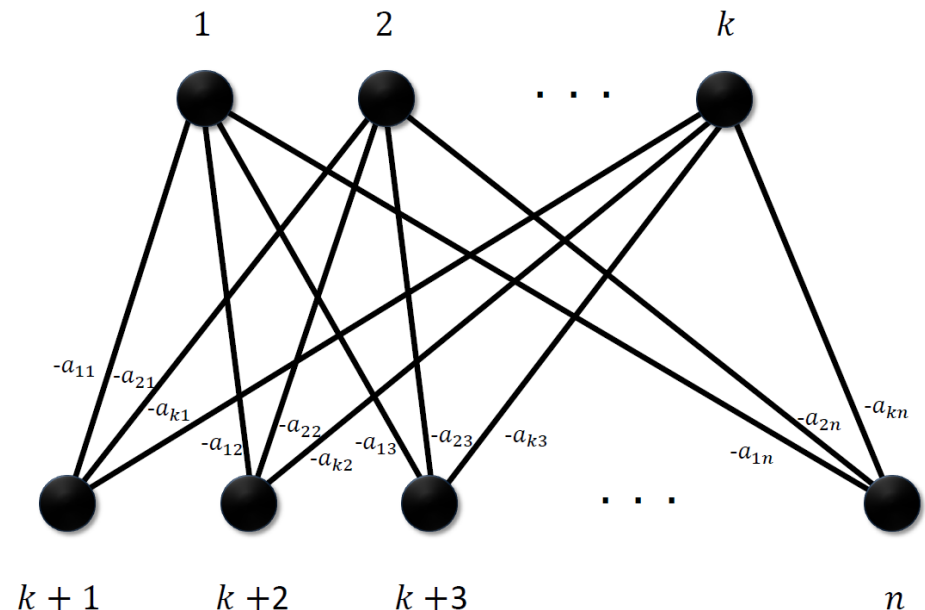
# Stabilizers formalism within the graph states

- The minimal support  $k$ -uniform state constructed from the MDS codes

$$k - \text{UNI}_{\min}(n, q) \ni |\psi\rangle = \sum_{\vec{v}} |\vec{v}G_{k \times n}\rangle \quad \longrightarrow \quad G_{k \times n} = [\mathbb{1}_k | A_{k \times n-k}]$$

$$g_l^\Gamma = X_l \prod_{m=0}^{n-1} (Z_m)^{\Gamma_{lm}} \psi \quad 1 \leq l \leq n$$

$$\Gamma\psi = - \begin{bmatrix} 0 & A^T \\ A^T & 0 \end{bmatrix}_{n \times n}$$



- A **complete bipartite graph** shows the structure of the graph represent the  $k$ -uniform state of minimal support

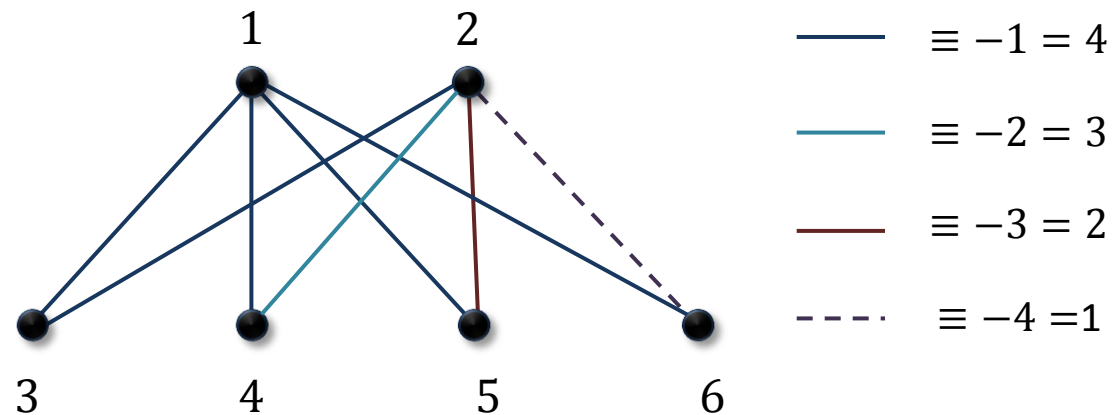
# An example of $k$ -uniform state

- Generator matrix of an MDS code  $[6,2,5]_5$

$$2 - \text{UNI}_{\min}(6,5) \ni |\psi\rangle = \sum_{i,j=0}^4 |i, j, i+j, i+2j, i+3j, i+4j\rangle$$

$$G_{2 \times 6} = [\mathbb{1}_2 | A] = \begin{bmatrix} 1 & 0 & 1 & 1 & 1 & 1 \\ 0 & 1 & 1 & 2 & 3 & 4 \end{bmatrix}$$

$s_1 =$	$X \quad \mathbb{1}$	$X \quad X \quad X \quad X$
$s_2 =$	$\mathbb{1} \quad X$	$X \quad X^2 \quad X^3 \quad X^4$
$s_3 =$	$Z^{-1} \quad Z^{-1}$	$Z \quad \mathbb{1} \quad \mathbb{1} \quad \mathbb{1}$
$s_4 =$	$Z^{-1} \quad Z^{-2}$	$\mathbb{1} \quad Z \quad \mathbb{1} \quad \mathbb{1}$
$s_5 =$	$Z^{-1} \quad Z^{-3}$	$\mathbb{1} \quad \mathbb{1} \quad Z \quad \mathbb{1}$
$s_6 =$	$Z^{-1} \quad Z^{-4}$	$\mathbb{1} \quad \mathbb{1} \quad \mathbb{1} \quad Z$



$$F^{-1}ZF = X$$

$$F^{-1}XF = Z^{-1}$$

# Basis

- Given a  $k$ -uniform state of minimal support  $|\psi\rangle \in \mathcal{H}(n, q)$  the  $q^n$  states

$$|\psi_{\vec{a}}\rangle := M(\vec{v}) |\psi\rangle$$

form a **complete orthonormal basis** of minimal support  $k$ -uniform state.

$$\begin{aligned} M(\vec{v}) &:= M(\vec{v}_Z) \otimes M(\vec{v}_X) \\ &= \underbrace{Z^{v_1} \otimes Z^{v_{k+1}} \otimes \dots \otimes Z^{v_k}}_k \otimes \underbrace{X^{a_{v_{k+1}}} \otimes X^{v_{k+2}} \otimes \dots \otimes X^{v_n}}_{n-k} \end{aligned}$$

$$\longrightarrow \langle \psi | M(\vec{v})^\dagger M(\vec{w}) | \psi \rangle = \prod_i \delta_{a_i, b_i}$$

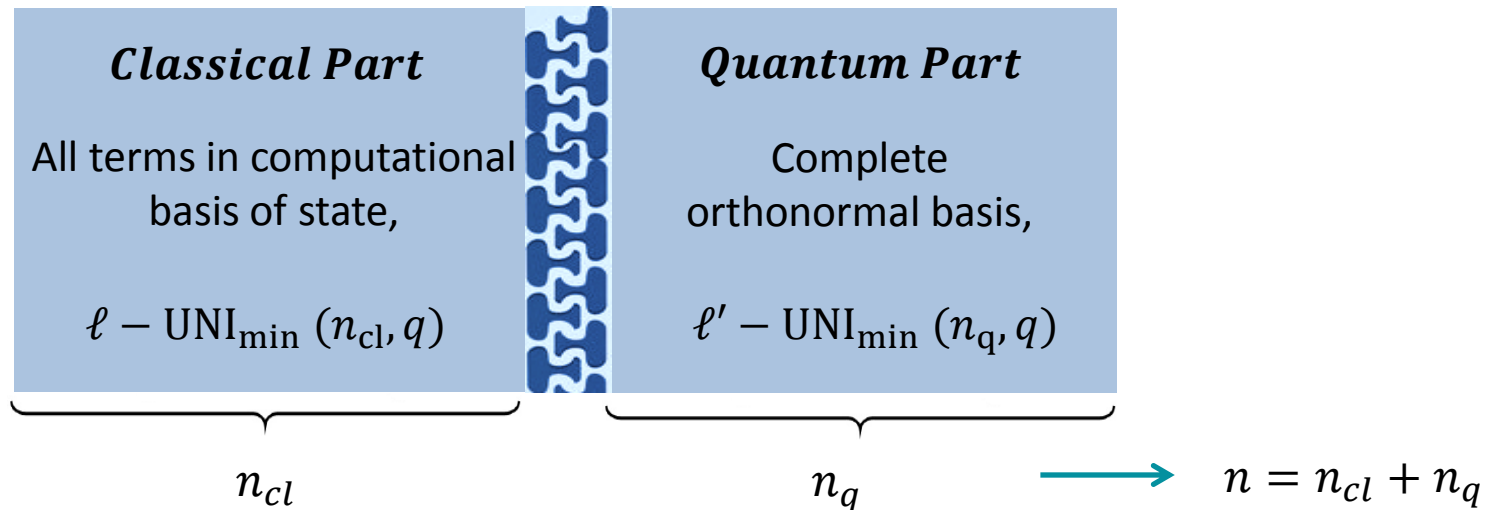
# Q-Clutch and constructing non-minimal support $k$ -uniform states

A systematic method to construct a set of non-minimal support  $k$  –  $\text{UNI}(n, q)$  states. We call this method **Q-Clutch**.



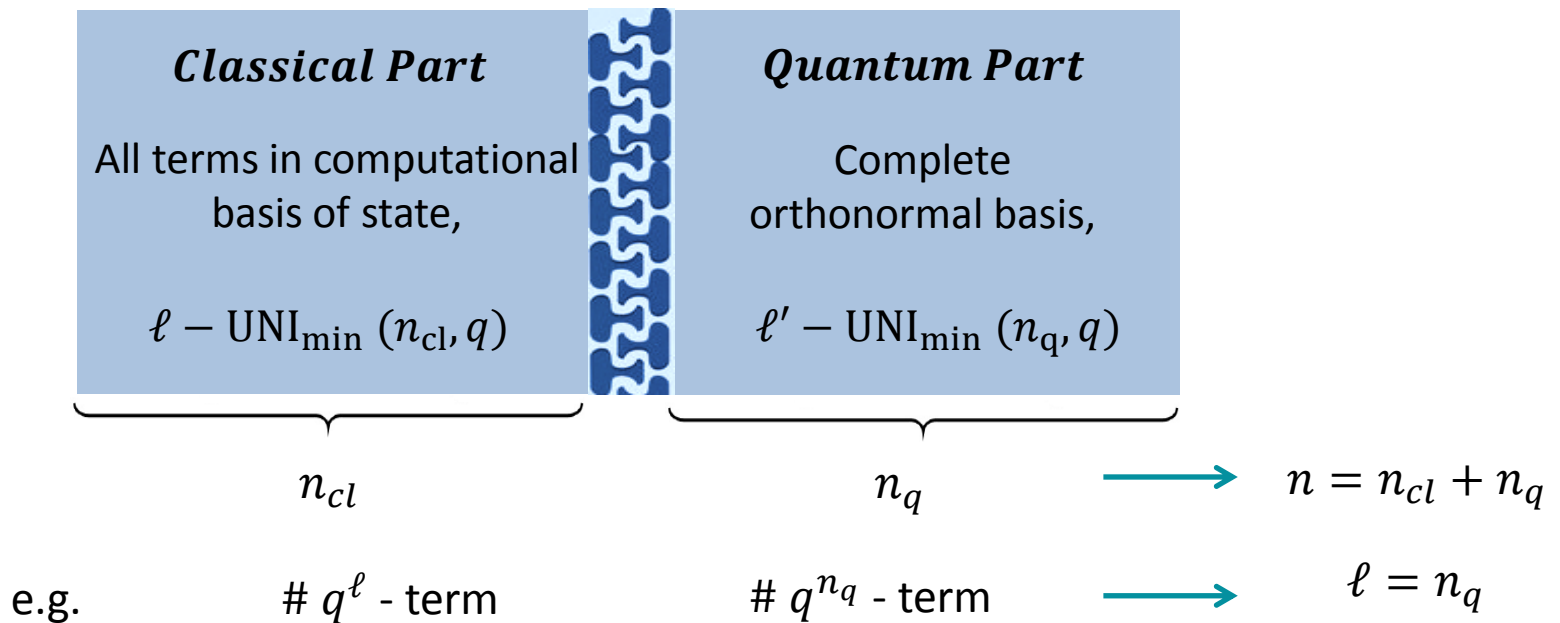
# $k$ -uniform states of non-minimal support

- The Q-Clutch connects **all the terms** of a  $\ell$ -uniform state to the quantum state of  $\ell'$ -uniform **basis** and lead to construction of **non-minimal support  $k$ -uniform state**



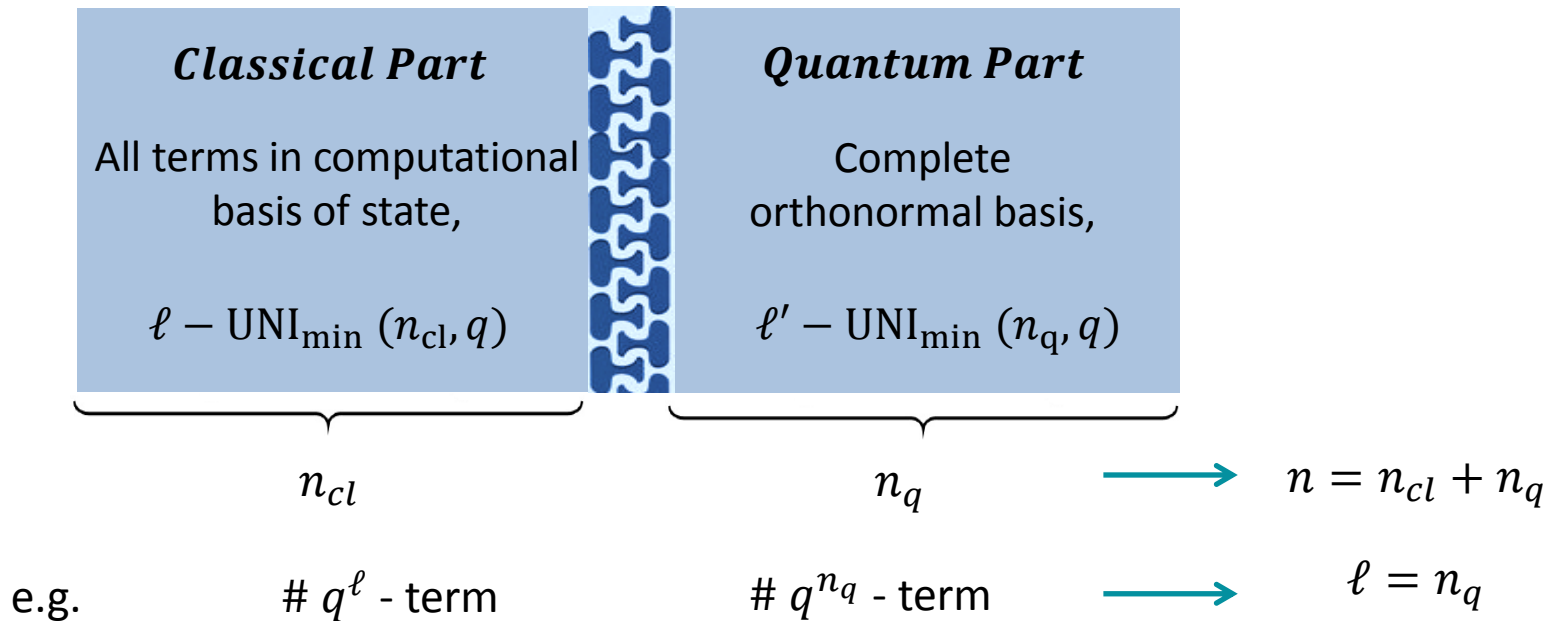
# $k$ -uniform states of non-minimal support

- The Q-Clutch connects **all the terms** of a  $\ell$ -uniform state to the quantum state of  $\ell'$ -uniform **basis** and lead to construction of **non-minimal support  $k$ -uniform state**



# $k$ -uniform states of non-minimal support

- The Q-Clutch connects **all the terms** of a  $\ell$ -uniform state to the quantum state of  $\ell'$ -uniform **basis** and lead to construction of **non-minimal support  $k$ -uniform state**



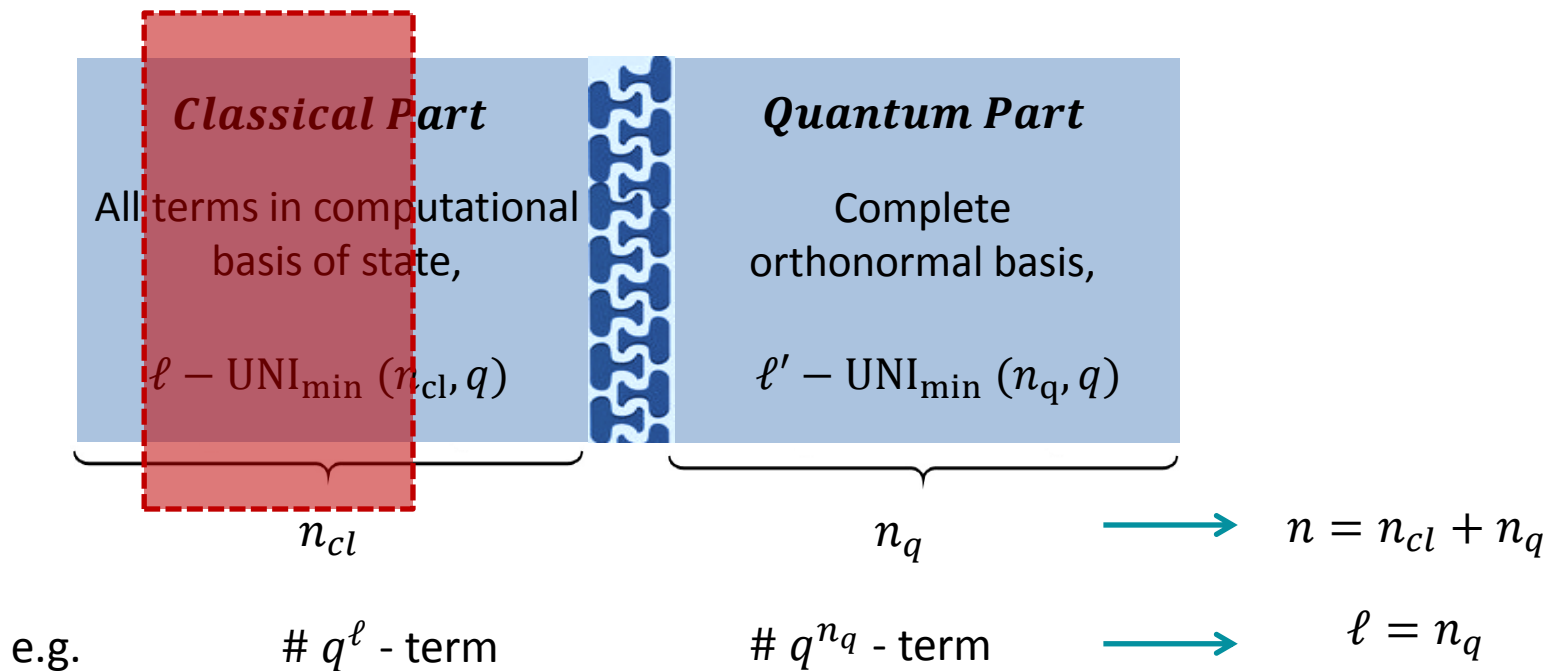
Hence we have non-minimal support  $k$ -uniform state:

$$k = \min\{\ell + 1, \ell' + 1\}$$



# $k$ -uniform states of non-minimal support

- The Q-Clutch connects **all the terms** of a  $\ell$ -uniform state to the quantum state of  $\ell'$ -uniform **basis** and lead to construction of **non-minimal support  $k$ -uniform state**

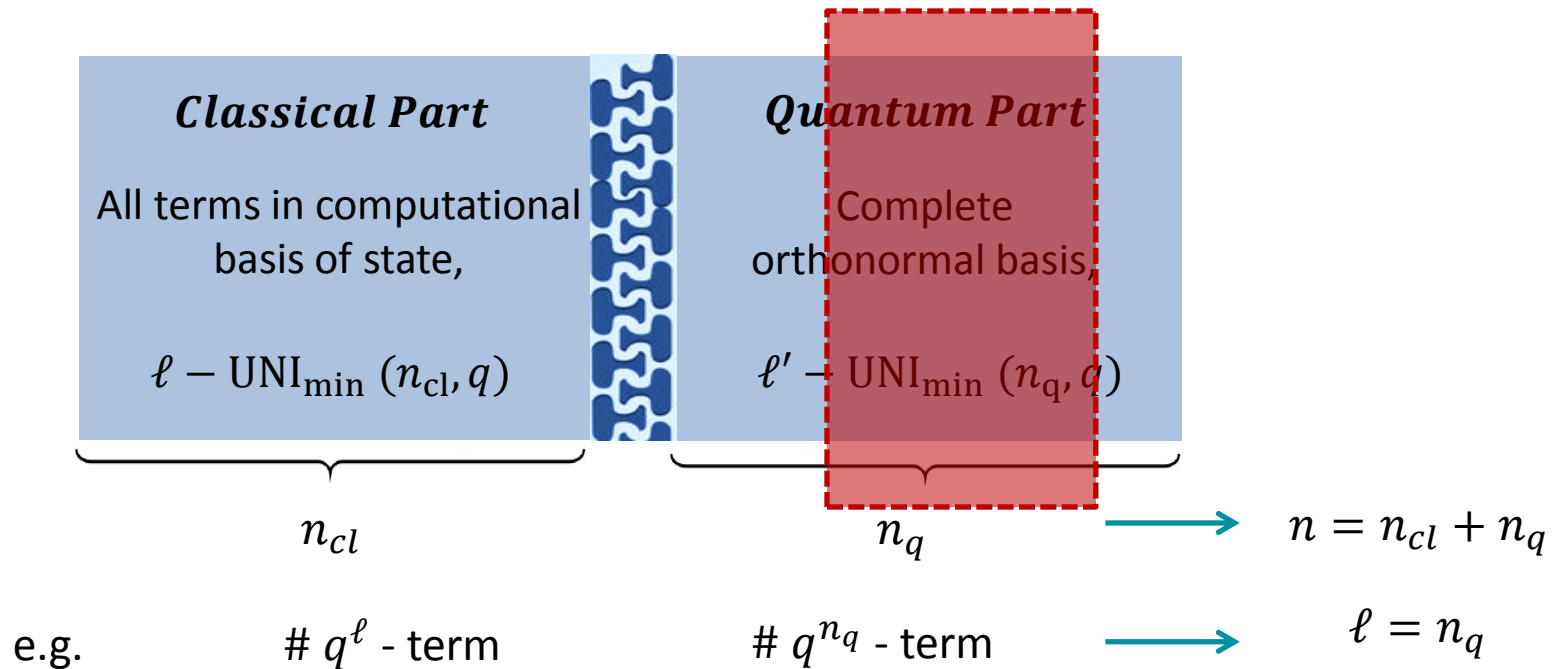


Hence we have non-minimal support  $k$ -uniform state:

$$k = \min\{\ell + 1, \ell' + 1\}$$

# $k$ -uniform states of non-minimal support

- The Q-Clutch connects **all the terms** of a  $\ell$ -uniform state to the quantum state of  $\ell'$ -uniform **basis** and lead to construction of **non-minimal support  $k$ -uniform state**

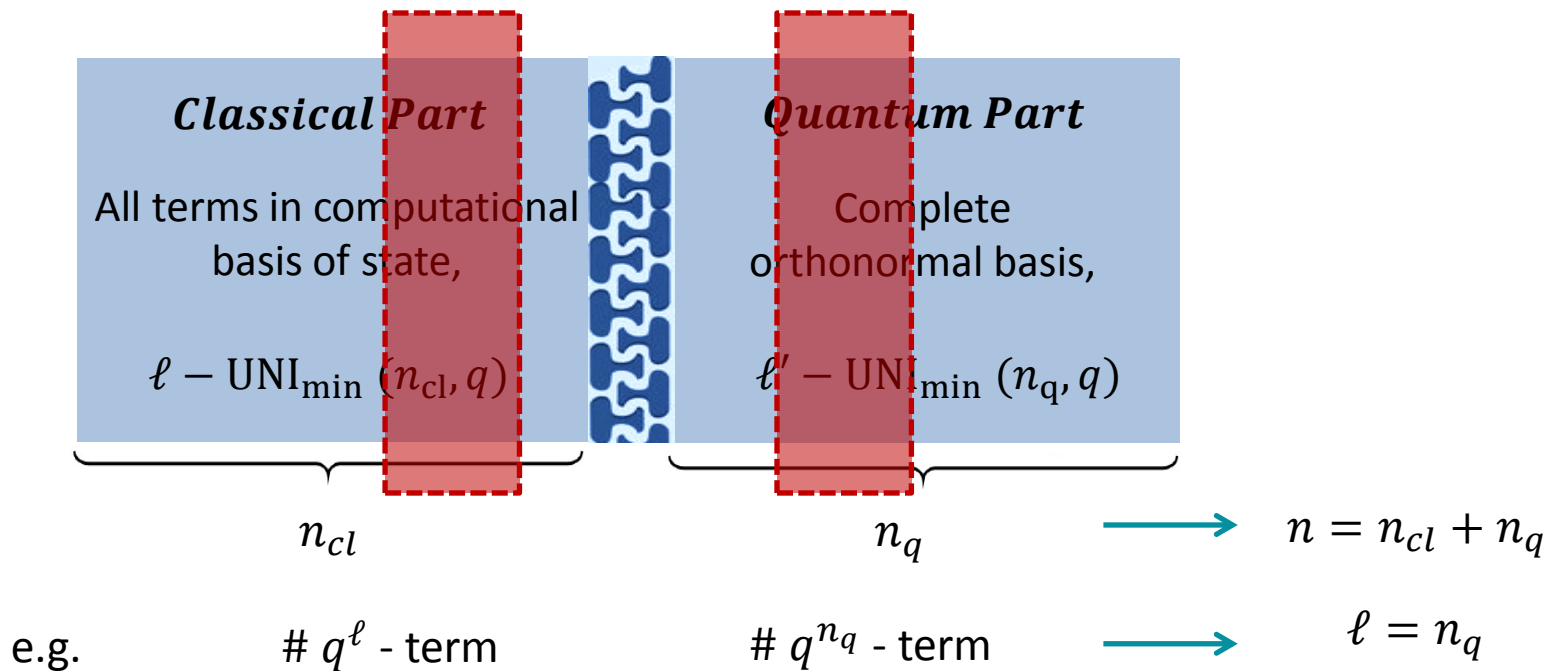


Hence we have non-minimal support  $k$ -uniform state:

$$k = \min\{\ell + 1, \ell' + 1\}$$

# $k$ -uniform states of non-minimal support

- The Q-Clutch connects **all the terms** of a  $\ell$ -uniform state to the quantum state of  $\ell'$ -uniform **basis** and lead to construction of **non-minimal support  $k$ -uniform state**



Hence we have non-minimal support  $k$ -uniform state:

$$k = \min\{\ell + 1, \ell' + 1\}$$

# Examples of $k$ -uniform of non-minimal support

- AME( $n = 5, q = 2$ ):

		<i>Classical Part</i>		<i>Quantum Part</i>	
		1	2	3	4&5
$ \psi\rangle =$		$ 0\rangle$	$ 0\rangle$	$ 0\rangle$	$ \phi^+\rangle$
		$ 0\rangle$	$ 1\rangle$	$ 1\rangle$	$ \psi^+\rangle$
		$ 1\rangle$	$ 0\rangle$	$ 1\rangle$	$ \psi^-\rangle$
		$ 1\rangle$	$ 1\rangle$	$ 0\rangle$	$ \phi^-\rangle$

$$AME(5,2) \ni |\psi\rangle = \sum_{i,j} |i, j, i+j\rangle Z^i \otimes X^j \sum_m |m, m\rangle$$

$\xleftrightarrow{\quad}$   
 $n_{cl} = 3$

$\xleftrightarrow{\quad}$   
 $n_q = 2$

# Examples of $k$ -uniform of non-minimal support

$$AME(7,4) \ni |\psi\rangle = \sum_{i,j,l} |i, j, l, i+j+l, i+xj+(1+x)l\rangle Z^{i+l} \otimes X^{j+xl} \sum_m |m, m\rangle$$

$\xleftarrow{n_{cl} = 5}$ 
 $\xrightarrow{n_q = 2}$

$$GF(2^2) = \{0, 1, x, x+1\}$$

$$\equiv \{0, 1, 2, 3\}$$

generated by  $x^2 = x + 1$

	2	3	4	5	6	7	8
2	Stab CMDS	Stab CMDS	Stab CMDS	Stab CMDS	Stab	Stab CMDS	Stab CMDS
3	Stab	Stab CMDS	Stab CMDS	Stab CMDS	Stab	CMDS	Stab CMDS
4		Stab OA	Stab CMDS	Stab CMDS		CMDS	CMDS
5	Stab OA QOA	Stab OA QOA	Stab OA QOA	Stab CMDS OA QOA	OA QOA	CMDS OA QOA	CMDS OA QOA
6	Stab OA QOA	Stab	Stab CMDS	Stab		CMDS	CMDS
7		Stab		Stab CMDS		CMDS	CMDS
8				Stab CMDS			CMDS



F. Huber, N. Wyderka, Table of absolutely maximally entangled states,  
<http://www.tp.nt.unisiegen.de/+fhuber/ame.html>

# Graphical representation

## Classical Part

All terms

$$\ell - \text{UNI}_{\min}(n_{\text{cl}}, q)$$



## Quantum Part

orthonormal basis

$$\ell' - \text{UNI}_{\min}(n_q, q)$$

$$|\Psi\rangle = |\phi\rangle|\psi_{\vec{v}}\rangle$$

$$|\psi_{\vec{v}}\rangle = M(\vec{v})|\psi\rangle$$

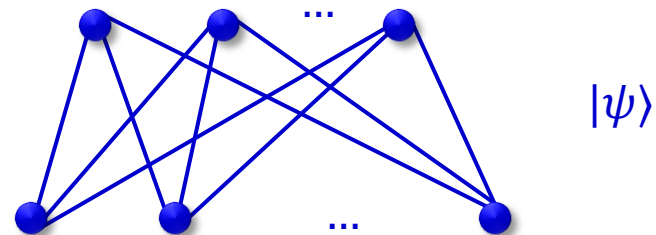
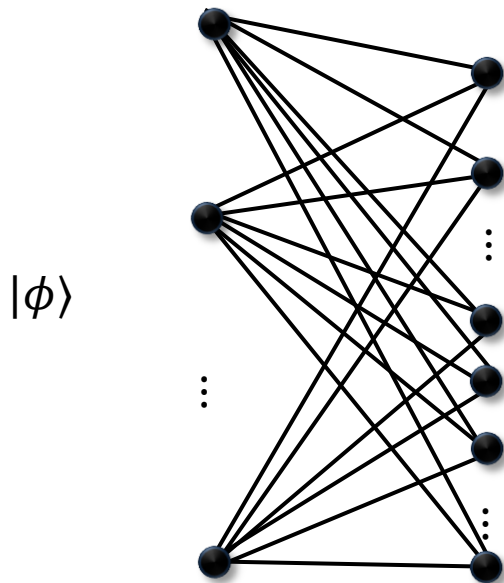
$$M(\vec{v}) := \underbrace{Z^{v_1} \otimes \dots \otimes Z^{v_{\ell'}}}_{\ell'} \otimes \underbrace{X^{a_{v_{\ell'+1}}} \otimes \dots \otimes X^{v_{n_q}}}_{n_q - \ell'}$$

e.g.  $\# q^{\ell}$  - term

$\# q^{n_q}$  - term



$\ell = n_q$



# Graphical representation

## Classical Part

All terms

$$\ell - \text{UNI}_{\min}(n_{\text{cl}}, q)$$



## Quantum Part

orthonormal basis

$$\ell' - \text{UNI}_{\min}(n_q, q)$$

$$|\Psi\rangle = |\phi\rangle|\psi_{\vec{v}}\rangle$$

$$|\psi_{\vec{v}}\rangle = M(\vec{v})|\psi\rangle$$

$$M(\vec{v}) := \underbrace{Z^{v_1} \otimes \dots \otimes Z^{v_{\ell'}}}_{\ell'} \otimes \underbrace{X^{a_{v_{\ell'+1}}} \otimes \dots \otimes X^{v_{n_q}}}_{n_q - \ell'}$$

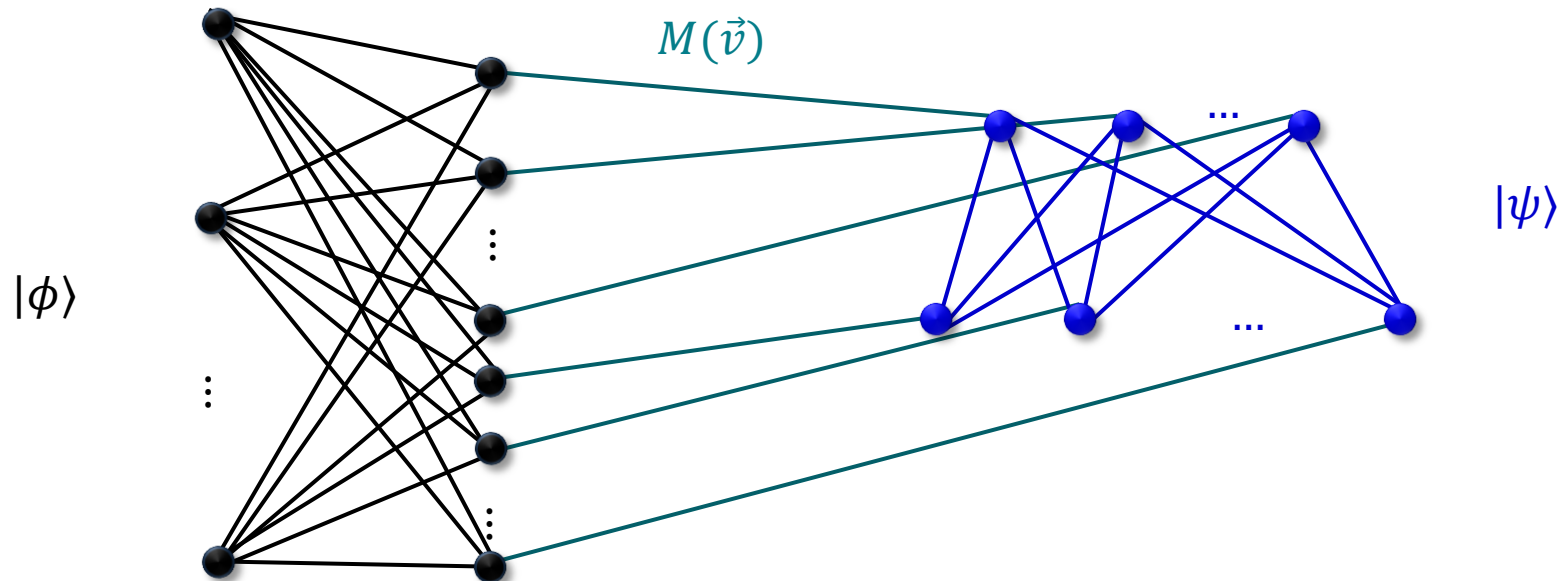
e.g.

#  $q^{\ell}$  - term

#  $q^{n_q}$  - term



$\ell = n_q$



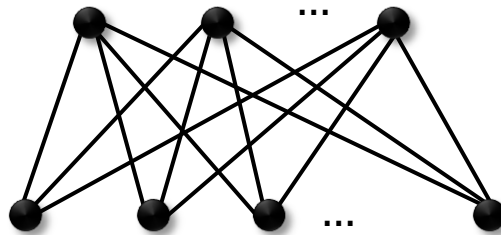
# Summary

Classical error  
correcting codes



$k$ -uniform states of  
minimal support

Graph states:



orthonormal basis:

$$|\psi_{\vec{a}}\rangle := M(\vec{v}) |\psi\rangle$$

## *Classical Part*

All terms

$$\ell - \text{UNI}_{\min}(n_{\text{cl}}, q)$$

## *Quantum Part*

orthonormal basis

$$\ell' - \text{UNI}_{\min}(n_{\text{q}}, q)$$



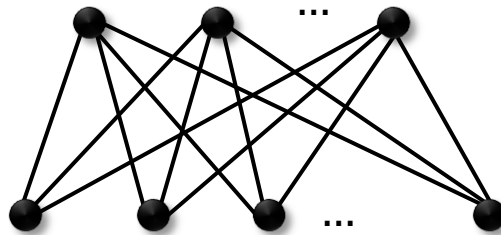
# Summary

Classical error  
correcting codes



$k$ -uniform states of  
minimal support

Graph states:



orthonormal basis:

$$|\psi_{\vec{a}}\rangle := M(\vec{v}) |\psi\rangle$$

***Classical Part***

All terms

$$\ell - \text{UNI}_{\min}(n_{\text{cl}}, q)$$

***Quantum Part***

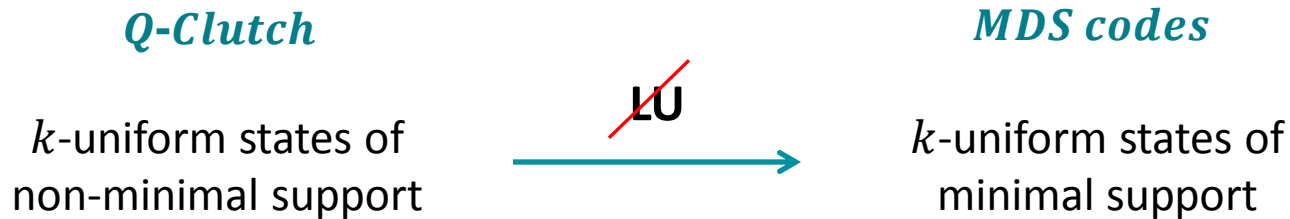
orthonormal basis

$$\ell' - \text{UNI}_{\min}(n_{\text{q}}, q)$$

Thank you for your attention!

# $k$ -uniform of non-minimal support vs $k$ -uniform of minimal support

- We show that the set of the states constructed by Q-Clutch method have **better parameters** compare to the  $k$ -uniform states that are obtained from the MDS codes.
- for given  $n$  and  $q$ :



We are  $\varepsilon$ -close to a full proof. . .

- Which state is better for **teleportation** and ... ?