

Investigation of Kaluza-Klein Compact Star Equation of States

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Outline

1. Motivation for Kaluza-Klein theory & stars
2. A special solution in $1+4D$ space-time
3. Compact Star in Kaluza-Klein World
4. M - R relation via varying R_c for $1+4D$ compact star
5. Comparison with standard hyperon star models



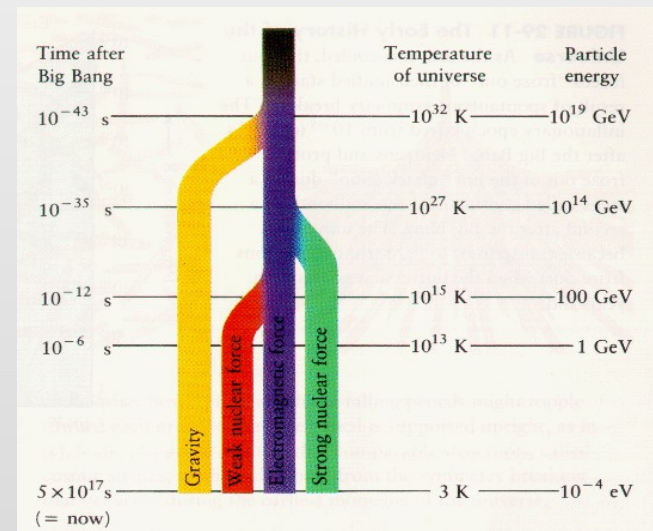
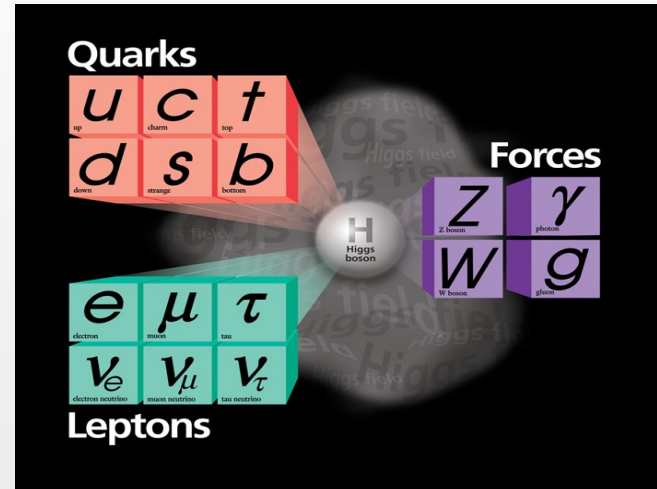
1. Motivation for Kaluza-Klein Stars

→ Standard matter by Standard Model

- Electromagnetic;
- Weak and
- Strong interactions

→ Grand Unified Theory...

- Gravity and QFT are not fitting into the same picture
- GR locally valid; curved space-time
- QFT globally valid; Minkowski



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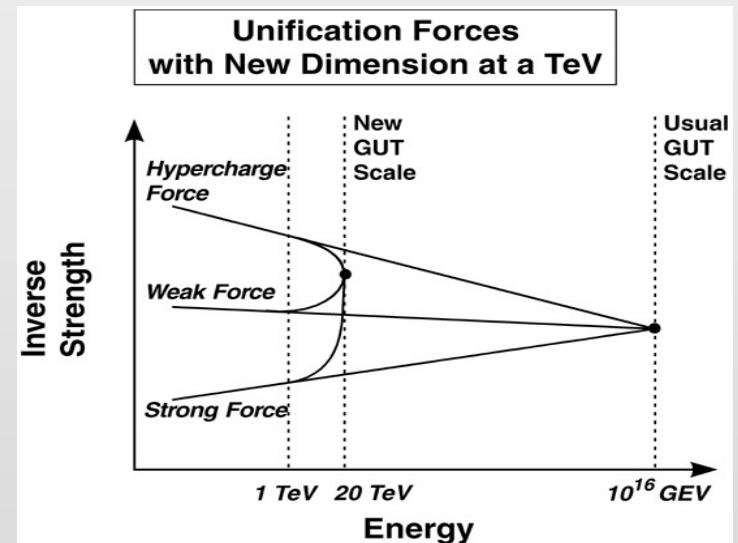
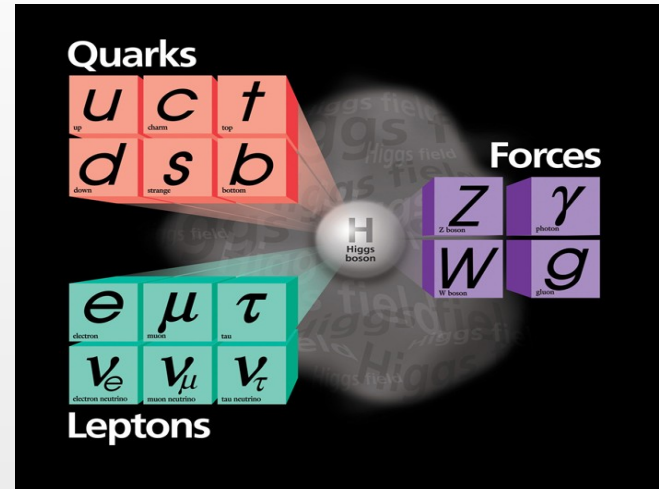
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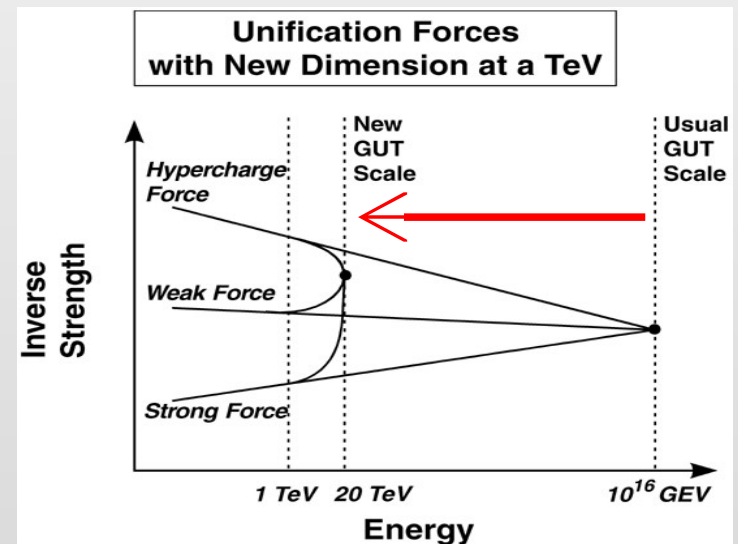
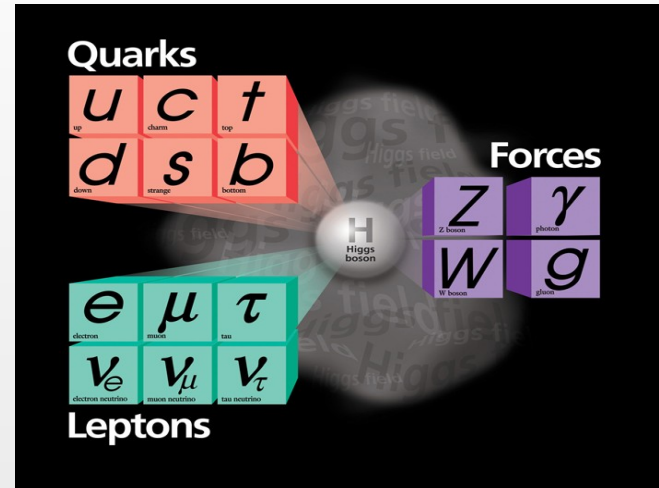
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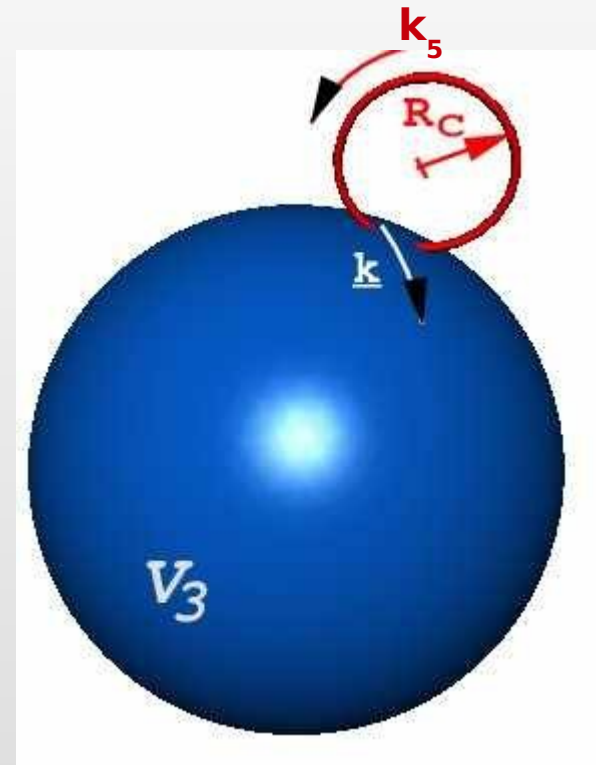
→ **Possible way:**

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1. Motivation for Kaluza-Klein Stars

- Compact Stars: extremely high energy density →
- ❖ Investigation of:
 - Overlap of strong-, electroweak & gravitational forces
 - Compactified extra dimensions
- Extended description of compact stars: introducing a compactified **extra dimension**
- Extra dimensional picture:
 - Excited states of particles: geometrical degrees of freedom
 - e.g. strangeness \sim particle moving in extra dimension



Symmetries in case of a Compact Star

- **Spherical:** invariant under $O(3)$ rotations
- **Static:** elements of metric and $T^{\mu\nu}$ are t independent
- **Ideal relativistic fluid:** only diagonal elements, $T^{\mu\nu}_{;\nu} = 0$
- **Isotropic:** no explicit φ & θ dependence of the components of $g^{\mu\nu}$ and $T^{\mu\nu}$

- **Extra dimension(s)** need extra assumptions...

[G. G. Barnaföldi, P. Lévai, B. Lukács:
„Searching extra
dimensions in compact stars”
Astron. Nachr. **328**, 809 (2007)]

Assumptions on Kaluza-Klein extra dimension(s)

- i. **$1+(3+d_c)$ dimensional space-time:**
dimensions are space-like, except first one: time-like
- ii. **GR is the same** as in $1+3$ D \rightarrow 'Equivalence Principle' is unchanged
- iii. **All causality postulates are the same** as in $1+3$ D
(including lightcone structure)
- iv. **Extra space-like d_c** dimensions are microscopical
- v. **Complete Killing-symmetry** in the d_c -dimensional
microscopical subspace

Focus on the simplest **$d_c=1$** case!

[G. G. Barnaföldi, P. Lévai, B. Lukács: J. Phys.: CS. **218** (2010)]

2. A special solution in 1+4D space-time Metric

- i. **Spherical symmetry** $\rightarrow O(3)$ symmetry
- ii. **Static picture** $\rightarrow g_{\mu 0} = 0$ and $g_{\mu\nu,0} = 0$
- iii. **4D** $g^{\mu\nu}$ is χ^5 independent
- iv. **Killing transformations** $\rightarrow g_{01} = 0$ and $g_{51} = 0$

$$g_{\mu\nu} = \begin{bmatrix} g_{00} & g_{01} & 0 & 0 & g_{05} \\ g_{01} & g_{11} & 0 & 0 & g_{15} \\ 0 & 0 & g_{22} & 0 & 0 \\ 0 & 0 & 0 & g_{22} \sin^2 \vartheta & 0 \\ g_{05} & g_{15} & 0 & 0 & g_{55} \end{bmatrix}$$

Coordinates:

$$t = x^0$$

$$r = x^1$$

$$\vartheta = x^2$$

$$\varphi = x^3$$

$$\chi = x^5$$

$$g_{\mu\nu} = \text{diag} (e^{2\nu}, -e^{2\lambda}, -r^2, -r^2 \sin^2 \vartheta, e^{2\Phi})$$

Radial functions for the metric components: $\nu(r)$, $\lambda(r)$, $\Phi(r)$

Energy-momentum tensor for 1+4D matter

(a) **Energy-momentum tensor** for anisotrop liquid:

$$T_{\mu\nu} := \epsilon u_\mu u_\nu - p (g_{\mu\nu} - u_\mu u_\nu + v_\mu v_\nu) - p_5 v_\mu v_\nu$$

Liquid is isotrop for 3 dimension: $T_1^1 = T_2^2 = T_3^3 = p$

BUT pressure $T_5^5 = p_5$ in the 5th direction is anisotrop,
with energy-density:

$$T_{\mu\nu} = \text{diag}(\epsilon e^{2\nu}, p e^{2\lambda}, p r^2, p r^2 \sin^2 \vartheta, p_5 e^{2\Phi})$$

(b) Let's construct the **$R_{\mu\nu}$ Ricci tensor** and **R Ricci scalar**:

$$R := R_i^i = R_1^1 = R_2^2 = R_3^3 = R_4^4 = R_5^5$$

[B. Lukács, T. Pacher: KFKI-1985-74,
Budapest, Hungary]

Einstein equations in 1+4D space-time

$$R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} = -8\pi G T_{\mu\nu}$$

Components of the **Einstein equation**:

$$\begin{aligned}
 -8\pi G \epsilon &= e^{-2\lambda} \left[\Phi'' + \Phi'^2 - \lambda' \Phi' + \frac{2\Phi'}{r} - \frac{2\lambda'}{r} + \frac{1}{r^2} \right] - \frac{1}{r^2} \\
 -8\pi G p &= e^{-2\lambda} \left[v' \Phi' - \frac{2\Phi'}{r} - \frac{2v'}{r} - \frac{1}{r^2} \right] + \frac{1}{r^2} \\
 -8\pi G p &= e^{-2\lambda} \left[-v'' - v'^2 - v'\lambda' + \Phi'' - \Phi'^2 - v'\Phi' + \lambda'\Phi' - \frac{v'}{r} + \frac{\lambda'}{r} - \frac{\Phi'}{r} \right] \\
 -8\pi G p_5 &= e^{-2\lambda} \left[-v'' - v'^2 + v\lambda - \frac{2v'}{r} + \frac{2\lambda'}{r} - \frac{1}{r^2} \right] + \frac{1}{r^2}
 \end{aligned}$$

- **Extra variables: $p_5, \Phi(r)$**

'Observation' of the 5th Dimension

(a) **Strangeness as a new degree of freedom:**

Connection between $\bar{m} = m_s$ and R_C radius:

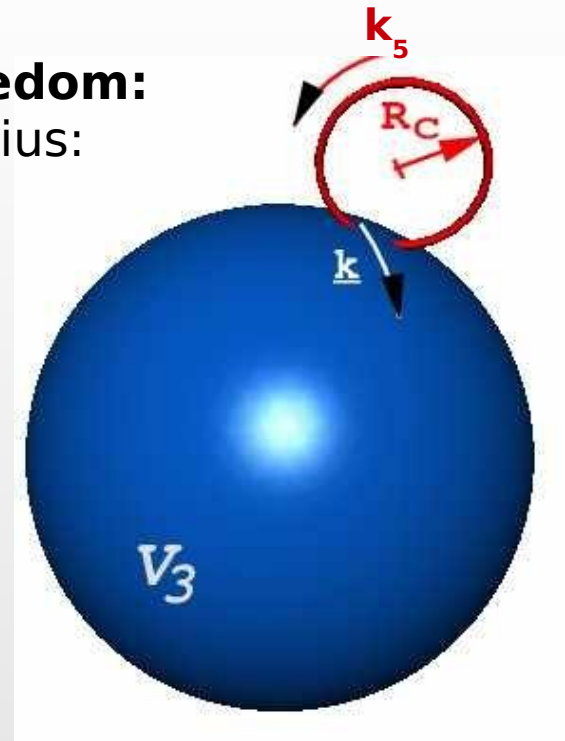
$$E_5 = \sqrt{\underline{k}^2 + \left(\frac{n}{R_C}\right)^2} + m^2 = \sqrt{\underline{k}^2 + \bar{m}^2}$$

$$\bar{m}^2 = \left(\frac{n}{R_C}\right)^2 + m^2$$

m: light (u, d) quark mass

n: excitation number, n=1

\bar{m} heavy (s) quark mass



(b) **Extra 5thD is compactified** in an **S¹ circle** with radius R_C

→ periodical boundary condition

$$\psi(x_5) \approx e^{ik_5 \cdot x_5} \quad \text{and} \quad \psi(x_5 + 2\pi R_C) \sim \psi(x_5) \quad \rightarrow \quad k_5 = \frac{n}{R_C}$$

k_5 : momentum in the 5th direction; x_5 : coordinate in the 5th direction; $n \in \mathbb{Z}^+$

Thermodynamics for 1+4D

(a) Thermodynamical potential for 1+4D Fermion gas

$$\Omega_5 = -2 \frac{V_4}{\beta} \int \frac{d^4 k}{(4\pi)^4} \left[\ln \left(1 + e^{-\beta(\sqrt{k^2 + \bar{m}^2} - \mu)} \right) (+\mu \leftarrow \rightarrow -\mu) \right]$$

$$\bar{m}^2 = (n/R_C)^2 + m^2 \quad \text{excited mass}$$

$$\int dk_5 \rightarrow \frac{1}{R_C} \Sigma_n \quad \text{discretization}$$

$$V_5 = 2\pi R_C V_4 \quad \text{volume}$$

(b) Thermodynamical potential and its quantities

$$\Omega_5 = \sum_n \Omega_4 \left(m^2 + \frac{n^2}{R_C} \right) = \Omega_4(\bar{m})$$

$$p = -\frac{1}{2\pi R_C} \frac{\partial \Omega_5}{\partial V} \quad p_5 = -\frac{1}{2\pi V} \frac{\partial \Omega_5}{\partial R_C} \quad \epsilon = \frac{U}{V_4}$$

TOV equations for 1+4D

(1) ' $d\Phi/dr = 0$ ' special case:

$$\frac{d p(r)}{d r} = - \frac{[p(r) + \epsilon(r)][M(r) + 4\pi r^3 p(r)]}{r[r - 2M(r)]}$$

$$p_5 = 1 - \frac{2M(r)}{r} + \frac{1}{r} \ln \left[1 - \frac{2M(r)}{r} \right] - 2p$$

3. Compact Star in Kaluza-Klein World

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Hyperon Star: n^0, Λ^0, \dots

vs.

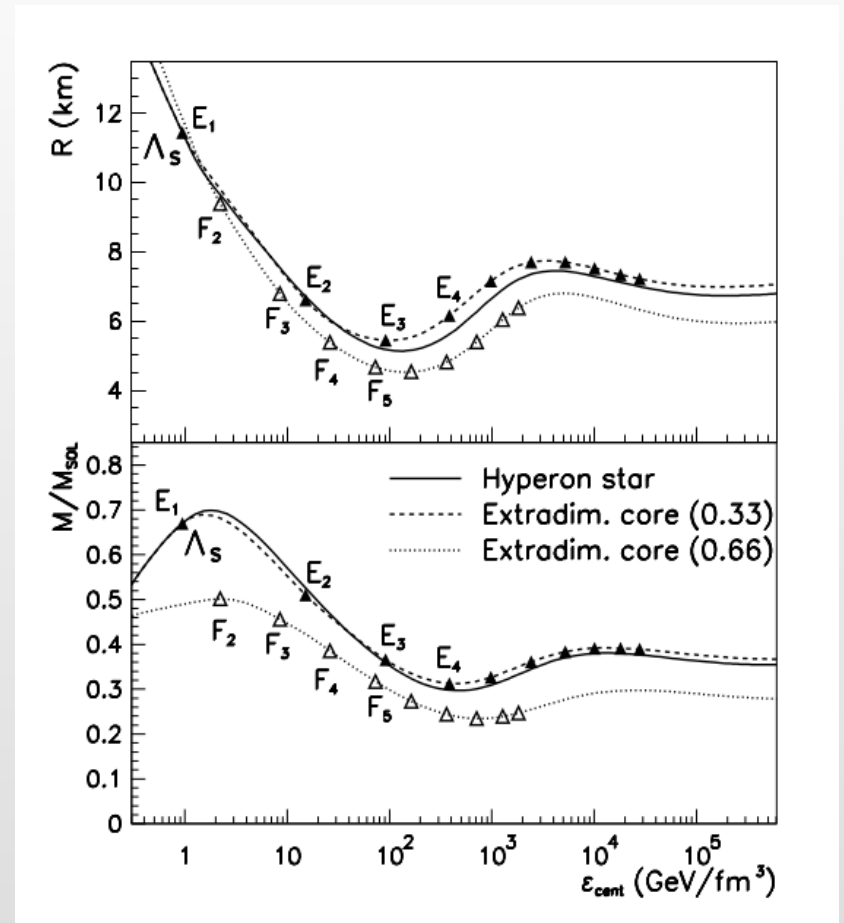
Compact star: with $k_F > \hbar/R_C$

1_c extra dimension

M(R) with n^{th} excited state:

▲ $R_c = 0.33$ fm (E_1, E_2, E_3, \dots)

△ $R_c = 0.66$ fm (F_1, F_2, F_3, \dots)



[G. G. Barnaföldi, P. Lévai, B. Lukács;
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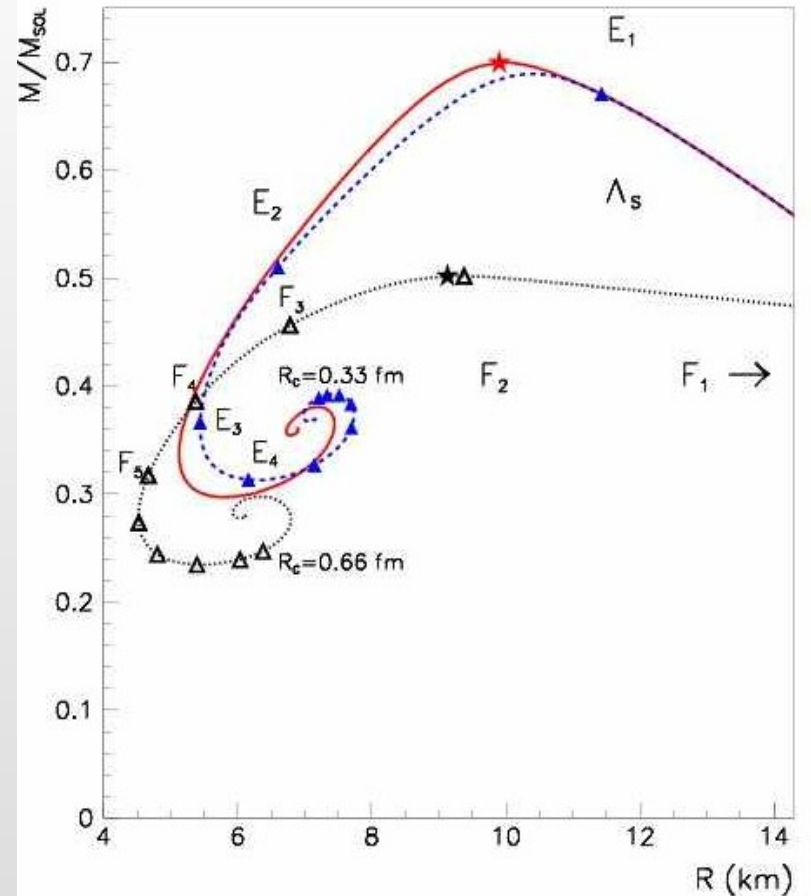
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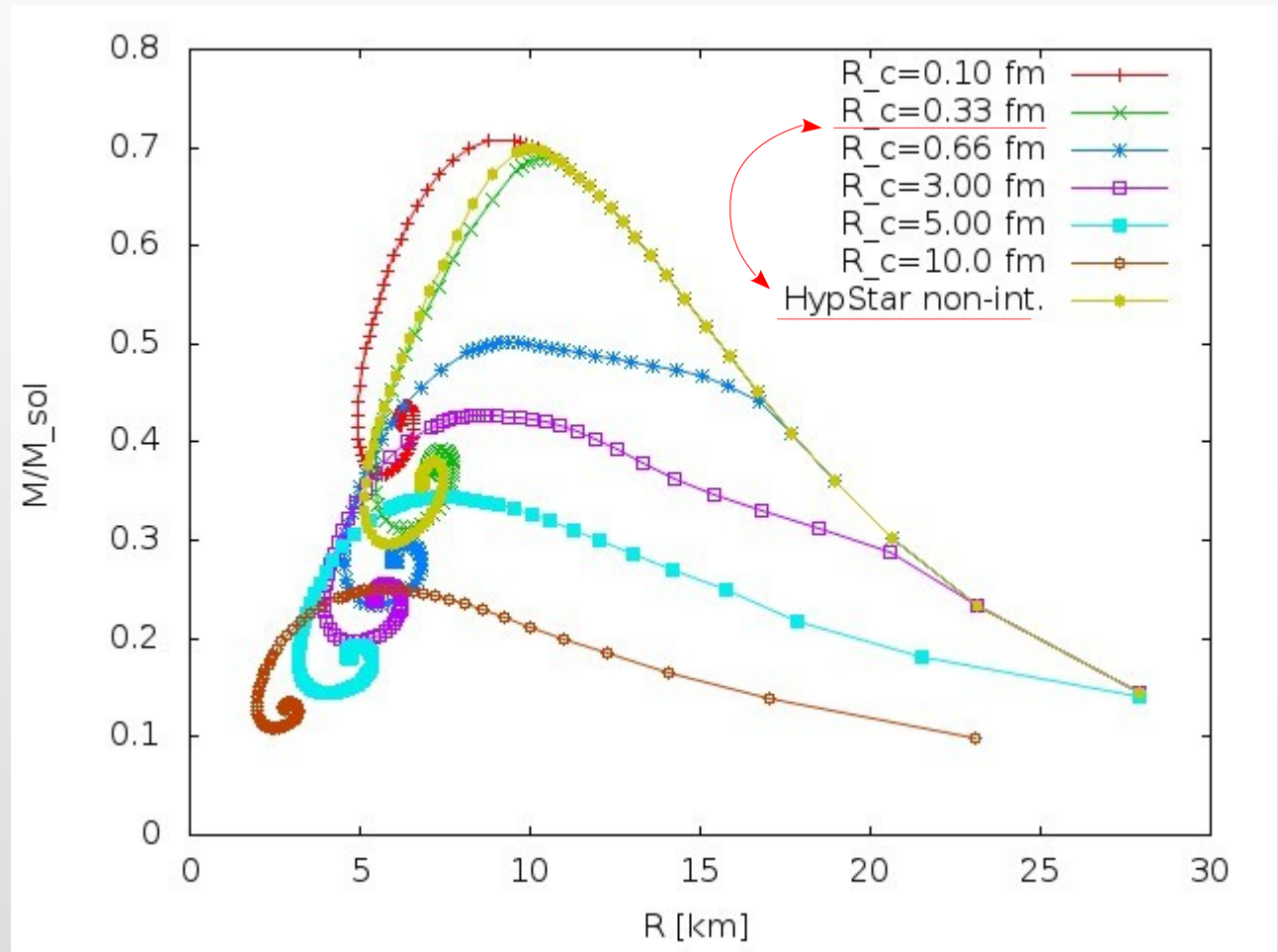


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Smooth variation of the compactified D

Change of the R_c :

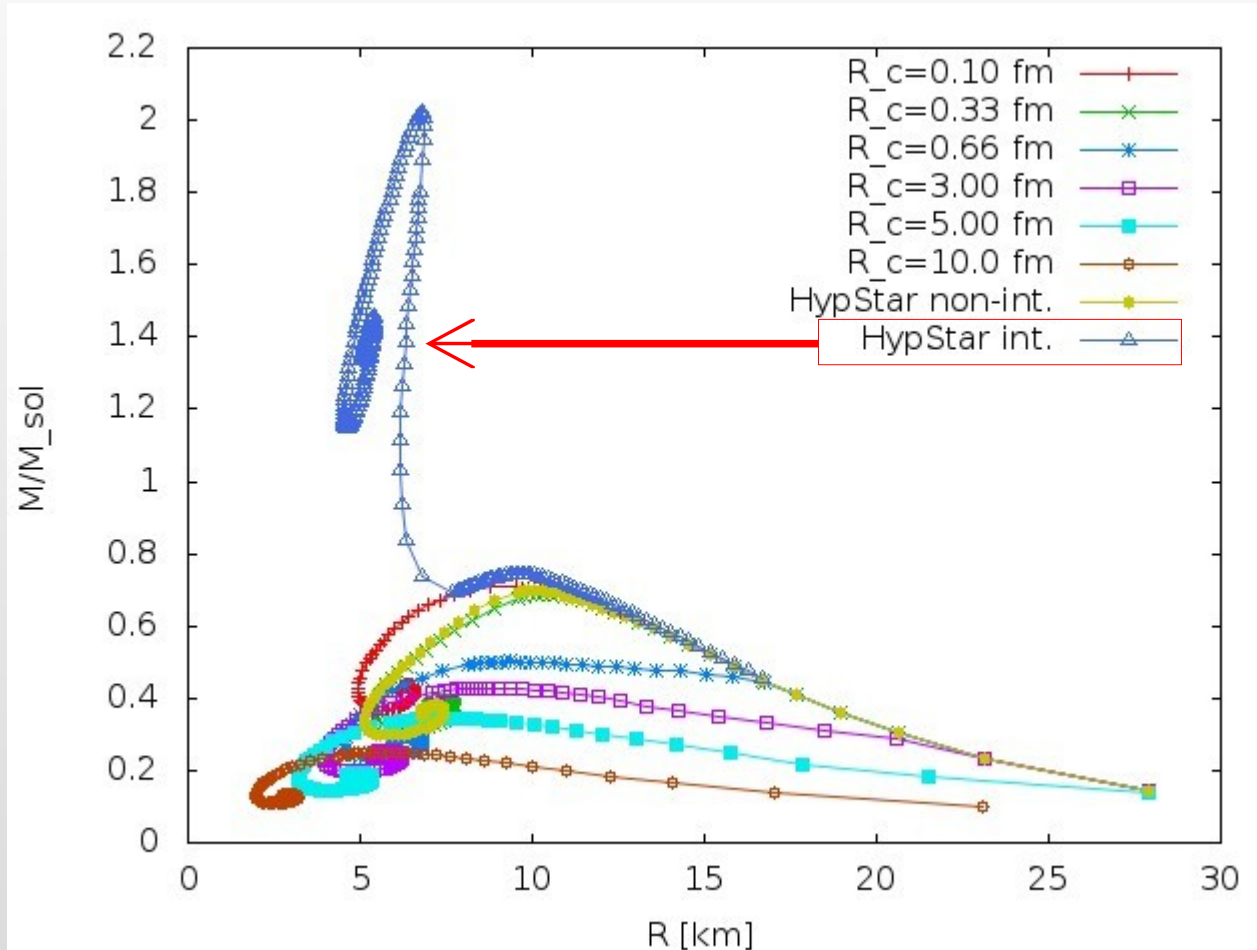
- excitation with:
- larger the $R_c \rightarrow \Delta m_n \sim \frac{n}{R_c}$
smaller the excitations to m_n
- similar solutions (shape, stability, etc.)
- **Hyp Star non-interacting**
 $\sim R_c = 0.33 \text{ fm}$



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- **Hyp Star interacting matter**
 \rightarrow more massive



Summary

→ Compact stars in 1+4D were analyzed:

- Static, spherical Schwarzschild-like space-time
- TOV-like eqs. with specific, but exact (stable) solution
- Solutions overlap with strange star models if R_c is set to the mass of strangeness: $10^{-13} \text{ cm} < R_c < 10^{-9} \text{ cm}$
- Mass limit: determined by R_c , $D \rightarrow$
- larger R_c results \rightarrow decrease in star's mass

→ Extra dimensional fermion stars:

- For solutions massive enough \rightarrow
- **Interaction of matter should be treated!**

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Thank You for Your attention!

Backup slides

Smooth variation of the compactified D

Change of the R_c :

- excitation with: larger the $R_c \rightarrow$ smaller excitations to m_n $\Delta m_n \sim \frac{n}{R_c}$

- Realistic Hyp EoS:**

Hyperon star
interacting matter

- Hyp Star interacting**

- Pure N Star interacting**

