SEGMENTED STRINGS AND HOLOGRAPHY

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¹HUN-REN Wigner Research Centre for Physics ²Budapest University of Technology and Economics 2024 Supported by the ÚNKP-23-3-I-BME-68 New National Excellence Program of the Ministry for Culture and Innovation from the source of the National Research, Development and Innovation Fund.



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Outline

- Introduction
- The AdS₃/CFT₂ correspondence
- The Ryu-Takayanagi formula
- Segmented strings in AdS₃
- Ryu-Takayanagi formula for segmented strings
- Correspondence in even dimensions
- Continuous limit
- Summary and outlook

Introduction



String theory on curved backround \uparrow

Field theory in Minkowski space

Original form: [Maldacena'97]

- IIB superstring theory on $AdS_5 \times S^5$
- $\mathcal{N} = 4$ supersymmetric Yang-Mills theory

Other aspect: [Ryu, Takayanagi'06] [Raamsdonk'10]

• Spacetime is built up of quantum entanglement



Introduction





AdS₃ space

 $\Lambda > 0$: de Sitter space

 $\Lambda < 0$: Anti-de Sitter (AdS) space

Einstein equations:

- In vacuum: $G_{\mu\nu} + \Lambda g_{\mu\nu} = 0$
- Cosmological constant: $\Lambda \leftarrow \Lambda = 0$: Minkowski space

AdS₃ space:

- Embedding space: $\mathbb{R}^{2,2}$
- 2+1 dimensional subset: AdS_3 space
- Metric:

$$ds^{2} = -(dX^{-1})^{2} - (dX^{0})^{2} + (dX^{1})^{2} + (dX^{2})^{2}$$

• Constraint:

 $X \cdot X = -(X^{-1})^2 - (X^0)^2 + (X^1)^2 + (X^2)^2 = -L^2$

Boundary of AdS₃ space:

•
$$\partial_{\infty}AdS_3 = \mathbb{P}\{U \in \mathbb{R}^{2,2} | U \cdot U = 0\}$$

AdS $\partial_{\infty}AdS$



AdS₃ space

Poincaré upper half-space model:

- 2 + 1 dimensional representation of the AdS_3 space
- Coordinates: (z, x^0, x^1)

$$(X^{-1}, X^{\mu}, X^{d}) = \left(\frac{-z^{2} - x^{2} - L^{2}}{2z}, \frac{Lx^{\mu}}{z}, \frac{-z^{2} - x^{2} + L^{2}}{2z}\right)$$

• Metric: $ds^{2} = L^{2} \frac{dz^{2} - (dx^{0})^{2} + (dx^{1})^{2}}{z^{2}}$
AdS
AdS
AdS
AdS
AdS
AdS
AdS

Boundary:

- $z \to 0 \Rightarrow$ Coordinates: (x^0, x^1)
- Metric: $ds^2 \propto -(dx^0)^2 + (dx^1)^2$
- Conformally equivalent to the 1+1 dimensional Minkowski space



The AdS/CFT correspondence

Conformal field theory:

- Quantum field theory
- With conformal invariance
- In d = 2: infinite-dimensional symmetry algebra
- Exactly solvable!

AdS₃/CFT₂ correspondence:

- $\partial_{\infty}AdS_3 \sim 1+1$ dimensional Minkowski space
- Conformal field theory on the boundary

Classical quantities (AdS_3) Field theoretical quantities (CFT_2)



$$Z_{grav}\left[\Phi\Big|_{\partial AdS}=J\right]=Z_{CFT}[J]$$



Entanglement in CFT₂

Von Neumann entropy:

- Statistical ensemble
- Density matrix ρ
- Von Neumann entropy: $S = -tr\{\rho \log \rho\}$
- Measures indeterminacy of the system

Entanglement in CFT₂:

- 1+1 dimensional CFT in vacuum state
- Observer that only has acces to a region A
- Measures different density matrix: $\rho_A = tr_{\bar{A}}\rho$
- Entanglement entropy: $S(A) = -tr\{\rho_A \log \rho_A\}$
- Measures the entanglement between A and \overline{A}

For an interval: [Calabrese, Cardy'18]

- Let *A* be an interval
- Length *R*
- *Note:* Cutoff dependent (δ)

$$S(A) = \frac{c}{3}\log\frac{R}{\delta}$$

HYSICS



c: central charge of CFT





Entanglement in CFT₂

Causal diamonds:

- Causal diamond = causality domain of a subsystem A
- Described by its past and future tips x^{μ} and y^{μ}
- Reduced density matrix: $\rho_A = e^{-H_A}$
- Where *H_A*: modular Hamiltonian

Kinematic space: [Boer'16]

- Space of causal diamonds (or subsystems) \rightarrow Coordinates: (x^{μ}, y^{ν})
- Coset structure: $\frac{SO(2,2)}{SO(1,1) \times SO(1,1)} \rightarrow$ Invariant metric: $\omega_{\mu\nu}$





Spherical minimal surfaces of AdS₃

Minimal surfaces of AdS₃:

- Two null vectors: $U, V \in \mathbb{R}^{2,2}$: $U \cdot U = V \cdot V = 0$
- Minimal surface= $\{X \mid U \cdot X = 0 \cap V \cdot X = 0\}$

In the Poincaré model:

- $U \cdot X = 0$ and $V \cdot X = 0 \rightarrow$ Two cones
- Tips of cones are on the boundary: x_u, x_v
- Minimal surface: One dimensional circular arc
- Image on the boundary: Interval, causal diamond

Area of minimal surfaces:

• Area: $A(U,V) = L \log \frac{4R^2}{\delta^2}$

• Where:
$$R^2 = -\frac{1}{4}(x_u - x_v) \cdot (x_u - x_v)$$

• Cutoff: $z > \delta$

Proportional to the entanglement entropy of the resulting boundary interval

2024.04.12.





$$x_{u}^{\mu} = L \frac{U^{\mu}}{U^{d} - U^{-1}}$$
$$x_{v}^{\nu} = L \frac{V^{\nu}}{V^{d} - V^{-1}}$$

 x_u

 $V \cdot X = 0$

 $U \cdot X = 0$

The Ryu-Takayanagi formula

The Ryu-Takayanagi formula: [Ryu, Takayanagi'06]

- *AdS*₃ space
- Vacuum CFT_2 on the boundary
- Brown-Henneaux formula: $c = \frac{3L}{2G}$ [Brown'86]
- An *A* spacelike *CFT* subsystem, border ∂A :
 - Entanglement entropy: $S(A) = -Tr\{\rho_A \log \rho_A\}$
- $\varepsilon_A AdS$ minimal surface, border ∂A :
 - Surface area: $A_{min}(A)$

$$S(A) = \frac{A_{min}(A)}{4G}$$

Ryu-Takayanagi formula

In our case:

- Image of minimal surface on the boundary: Interval
- Entanglement entropy: $S(U, V) = \frac{L}{4G} \log \frac{4R^2}{\delta^2}$











Classical strings in AdS₃

One dimensional classical strings:

- String: One dimensional object
- Propagation in spacetime: Two dimensional "worldsheet"
- Two parameters: (τ, σ) or (σ^+, σ^-)

• Action:
$$S = -\frac{T}{2} \int d\tau d\sigma \sqrt{-h} h^{ab} \partial_a X \cdot \partial_b X$$
 ~ Surface area

• $\delta S = 0$ { Equation of motion Virasoro constraints

Strings in AdS₃:

- Embedding space: AdS₃
- Equations of motion:

$$\partial_{+}\partial_{-}X - (\partial_{-}X \cdot \partial_{+}X)X = 0$$

$$\partial_{+}X \cdot \partial_{+}X = \partial_{-}X \cdot \partial_{-}X = 0$$

$$X \cdot X = -L^{2}$$



• Normal vector: $N_a = \frac{\epsilon_{abcd} X^b \partial_- X^c \partial_+ X^d}{\partial_- X \cdot \partial_+ X}$



Segmented strings in AdS space

Segmented strings: [Callebaut'15]

- Simplest solution: constant normal vector
- String segment: quadrangle with constant normal vector
- Segmented string: solution built up by segments
- Vertices: $V_i \cdot V_i = -L^2$, i = 1,2,3,4
- Edges: $p_i = \pm (V_i V_{i+1}), i = 1,2,3,4$ $p_i \cdot p_i = 0$

$$X(\sigma^{-},\sigma^{+}) = \frac{L^{2} + \sigma^{+}\sigma^{-}\frac{1}{2}p_{1} \cdot p_{4}}{L^{2} - \sigma^{+}\sigma^{-}\frac{1}{2}p_{1} \cdot p_{4}}V_{1} + L^{2}\frac{\sigma^{-}p_{1} + \sigma^{+}p_{4}}{L^{2} - \sigma^{+}\sigma^{-}\frac{1}{2}p_{1} \cdot p_{4}}$$

Area of a string segment:

• Evaluating the string action with the segmented solution

$$A_{::} = L^2 \log \frac{(p_1 \cdot p_4)(p_2 \cdot p_3)}{(p_1 \cdot p_2)(p_3 \cdot p_4)}$$

• *Note:* Cutoff independent!



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Segmented strings and the boundary

Projection to the boundary:

- $p_i \cdot p_i = 0$!
- p_i define four cones in AdS via $p_i \cdot X = 0$
- At each vertex two edges meat
- Therefore each vertex lies on an intersection of a pair of cones
- Let: $x_i^{\mu} = L \frac{p_i^{\mu}}{p_i^d p_i^-}$
- Then (x_i^{μ}, x_j^{ν}) defines a causal diamond and a subsystem (ij)

Area of the string segment:

$$A_{::} = L^2 \log \frac{R_{14}^2 R_{23}^2}{R_{12}^2 R_{34}^2} = A_{14} + A_{23} - A_{34} - A_{12}$$

• Where:
$$R_{ij}^2 = -\frac{1}{4}(x_i - x_j)^2$$







Ryu-Takayanagi formula for segmented strings

Timelike segmented string:

- Edges: $p_i : p_i \cdot p_i = 0, i = 1, 2, 3, 4$
- For vertices: $p_i \cdot X = 0$
- Area: $A = L^2 \log \frac{R_{14}^2 R_{23}^2}{R_{12}^2 R_{34}^2}$

Spacelike minimal surfaces:

- Cones: $p_i \cdot X = 0$
- Their intersections: (14), (23), (12), (34) minimal surfaces
- Segment vertices lie on them!
- Their areas: $A(ij) = L \log \frac{4R_{ij}^2}{\delta^2}$

CFT vacuum subsystems:

- Images of minimal surfaces: (*ij*) subsystems
- Their entanglement entropies: $S(ij) = \frac{A(ij)}{4G}$

 $A \equiv 4GL(S(14) + S(23) - S(34) - S(12))$







Interpretation and consequences

Interpretation:

- CFT subsystems
- Flow of boundary causal diamonds

Interesting properties: Eg.:

- For a segment: $p_1 + p_2 = p_3 + p_4$
- On the boundary: causal diamonds lie on a common hyperbola
- Accelerating frame

Consequences: Eg.: strong subadditivity:

• Subsystems A, B, C

 $S(AB) + S(BC) - S(B) - S(ABC) \ge 0$

Strong subadditivity!

- Measuring on a larger subsystem reduces uncertainity
- $\exists A, B, C$ subsystems: $A_{\Box} \sim S(AB) + S(BC) S(B) S(ABC)$
- Geometrically: Positivity of segment area $A \ge 0$







Minimal surfaces in higher dimensions

Minimal surfaces of the AdS_{d+1} :

- Two null vectors: $U, V \in \mathbb{R}^{2,d}$: $U \cdot U = V \cdot V = 0$
- Minimal surface= $\{X \mid U \cdot X = 0 \cap V \cdot X = 0\}$

In the Poincaré model:

- $U \cdot X = 0$ and $V \cdot X = 0 \rightarrow$ Two cones
- Tips of the cones are on the boundary: x_u, x_v
- Minimal surface: d-1 dimensional sphere
- Image on the boundary: d-2 sphere, causal diamond

Area of minimal surfaces: [Ryu, Takayanagi'06]

• If
$$d = \text{even:} \quad A = \left(Powers \ of \frac{R}{\delta} \right) + \alpha \log \frac{R^2}{\delta^2}$$

• If
$$d = \text{odd}$$
 $A = \left(Powers \ of \frac{R}{\delta}\right)$



Correspondence in even dimensions

Ryu-Takayanagi proposal in higher dimensions:

- AdS_{d+1} space
- Boundary CFT_d in vacuum state
- A spacelike CFT subsystem with boundary ∂A
- $\varepsilon_A AdS$ minimal surface with boundary ∂A
- Entanglement entropy in general:

$S(A) = \frac{A_{min}(A)}{4G}$

Entropy in higher dimensions:

• $S = S^{\partial} + S^{uni} + S^{other}$

•
$$S \propto \frac{A(\partial A)}{\delta^{d-2}}$$

- If d = even: $S^{uni} \propto \log \frac{R}{\delta}$
- If d = odd: $S^{uni} = const$.







Segmented strings in higher dimensions

Strings in AdS_{d+1}:

- Embedding space: AdS_{d+1}
- Equations of motion:

Segmented strings:

- Vertices: $V_i \cdot V_i = -L^2$, i = 1,2,3,4
- Edges: $p_i = \pm (V_i V_{i+1}), i = 1,2,3,4$
- Same interpolation ansatz

Projection to the boundary:

• Let:
$$x_i^{\mu} = L \frac{p_i^{\mu}}{p_i^d - p_i^-}$$

Area of the string segment:

$$A_{::} = L^2 \log \frac{R_{14}^2 R_{23}^2}{R_{12}^2 R_{34}^2}$$

• Where:
$$R_{ij}^2 = -\frac{1}{4} (x_i - x_j)^2$$

$$\partial_{+}\partial_{-}X - (\partial_{-}X \cdot \partial_{+}X)X = 0$$

$$\partial_{+}X \cdot \partial_{+}X = \partial_{-}X \cdot \partial_{-}X = 0$$

$$X \cdot X = -L^{2}$$

Correspondence in even dimensions:

- d = even
- Edges: $p_i \leftrightarrow \text{Cones}$
- Vertices on minimal surfaces
- Minimal surfaces ↔ Spherical subsystems
- Area of segment \leftrightarrow Entanglement

$$A \sim (S^{uni}(14) + S^{uni}(23) - S^{uni}(34) - S^{uni}(12))$$



Continuous limit

Continuous limit:

- General AdS_{d+1} space
- Directional derivatives: ∂_X, ∂_X \rightarrow Virasoro constraints: $\partial_X, \partial_X = \partial_X \cdot \partial_X = 0$
- Poincaré coordinates: $x^{\mu} = L \frac{\partial_{-} X^{\mu}}{\partial_{-} X^{d} \partial_{-} X^{-1}}$, $y^{\mu} = \frac{\partial_{+} X^{\mu}}{\partial_{+} X^{d} \partial_{+} X^{-1}}$
- On the worldsheet: $\partial_X \cdot X = 0$, $\partial_+ X \cdot X = 0$ \rightarrow Causal diamond with tips x^{μ} , y^{μ}
- String action in causal diamond coordinates:

$$S = \int d\sigma^{-} d\sigma^{+} \sqrt{-h} h^{ab} \omega_{\mu\nu} \partial_{(a} x^{\mu} \partial_{b)} y^{\nu}$$

$$\downarrow$$

$$\omega_{\mu\nu} \equiv \text{Kinematic space metric}$$

$$\downarrow$$

$$\frac{SO(2,2)}{SO(1,1) \times SO(1,1)} \text{ invariant!}$$





Fidelity susceptibility

Infinitesimally close causal diamonds:

- Causal diamonds at (x^{μ}, y^{μ}) and $(x^{\mu} + \delta x^{\mu}, y^{\mu} + \delta y^{\mu})$
- Density matrices: ρ and $\rho + \delta \rho$
- Parallel purifications: ψ and $\psi + \delta \psi$ [Uhlmann'86]

Fidelity susceptibility and complexity:

- Overlap of states: $|\langle \psi | \psi + \delta \psi \rangle| = 1 \omega_{\mu\nu}^{FS} \delta x^{\mu} \delta y^{\nu} + \cdots$
- Where: $\omega_{\mu\nu}^{FS}$: fidelity susceptibility
- Kinematic space metric $\omega_{\mu\nu} \propto \omega_{\mu\nu}^{FS}$
- Where prefactors contain geometric factors and a_d^* trace anomaly

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Bulk string geometry ↔ Boundary quantum geometry











Summary

Summary:

- AdS_3/CFT_2 correspondence
- Ryu-Takayanagi formula → Connection between surfaces of segmented strings, areas of minimal surfaces and entanglement entropies of subsystems in CFT vacuum
- Geometry of classical strings \leftrightarrow Field theoretical entanglement
- Dictionary: Eg.:
 - \circ Positivity of area \leftrightarrow Strong subadditivity
- Duality for segmented strings in AdS_{d+1}/CFT_d , if d = even
- Continious limit \leftrightarrow Quantum geometry

Outlook:

- Entanglement as a glue?
- *CFT* excitations?
- Further entanglement inequalities?





THANK YOU FOR YOUR ATTENTION!



[B. Boldis, P. Lévay, Phys. Rev. D 109, 046002]



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